

Neutrinos: from past surprises to new puzzles?

21 January 2022

U. Torino

Silvia Pascoli



Horizon2020

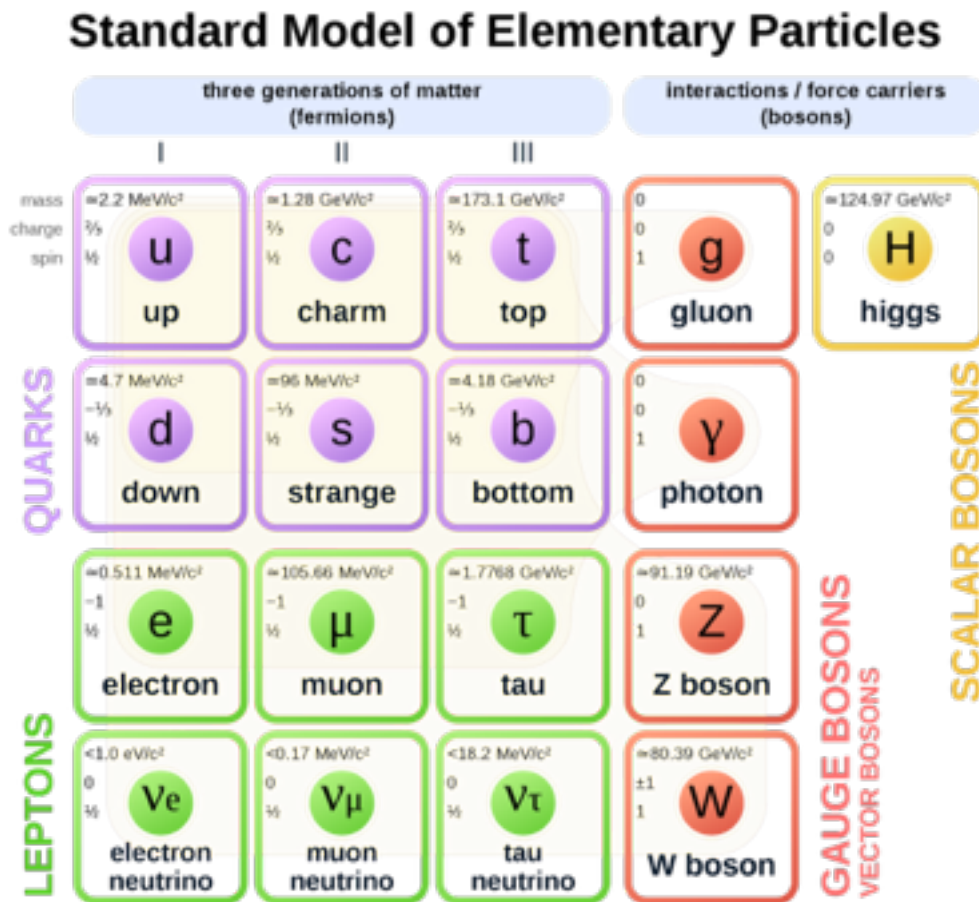


ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA



Neutrinos in the SM

Neutrinos are the lightest and most elusive of all the known **elementary fermions**.



- Neutrinos come in 3 flavours (types), corresponding to each of the charged leptons.

- Neutrinos in the SM are of only one type (left-handed Weyl spinors) as the right-handed component is not included.

SM is a gauge theory: $SU(3)_c \times SU(2)_L \times U(1)_Y$

Particles	$SU(3)$	$SU(2)_L$	$U(1)_Y$
Leptons			
$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L, \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L, \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$	1	2	$-1/2$
e_R, μ_R, τ_R	1	1	-1
Quarks			
$\begin{pmatrix} u \\ d \end{pmatrix}_L, \begin{pmatrix} c \\ s \end{pmatrix}_L, \begin{pmatrix} t \\ b \end{pmatrix}_L$	3	2	$1/6$
u_R, c_R, t_R	3	1	$2/3$
d_R, s_R, b_R	3	1	$-1/3$

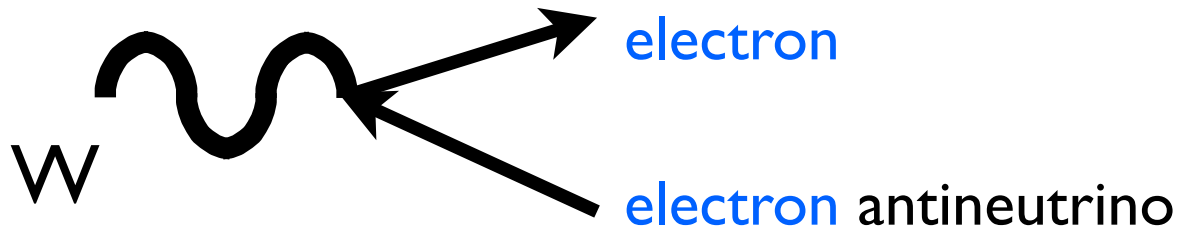
Neutrinos belong to $SU(2)$ doublets and are neutral.

- They carry lepton number, $U(1)_{\text{lepton}}$

$$\ell \xrightarrow{U(1)_{\text{lepton}}} e^{i\alpha} \ell \qquad \nu_L \xrightarrow{U(1)_{\text{lepton}}} e^{i\alpha} \nu_L$$

They have charge current (CC) and neutral current (NC) interactions

$$\mathcal{L}_{\text{SM}} = -\frac{g}{\sqrt{2}} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_{\alpha L} \gamma^\mu \ell_{\alpha L} W_\mu - \frac{g}{2 \cos \theta_W} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_{\alpha L} \gamma^\mu \nu_{\alpha L} Z_\mu + \text{h.c.}$$



Number of active neutrinos

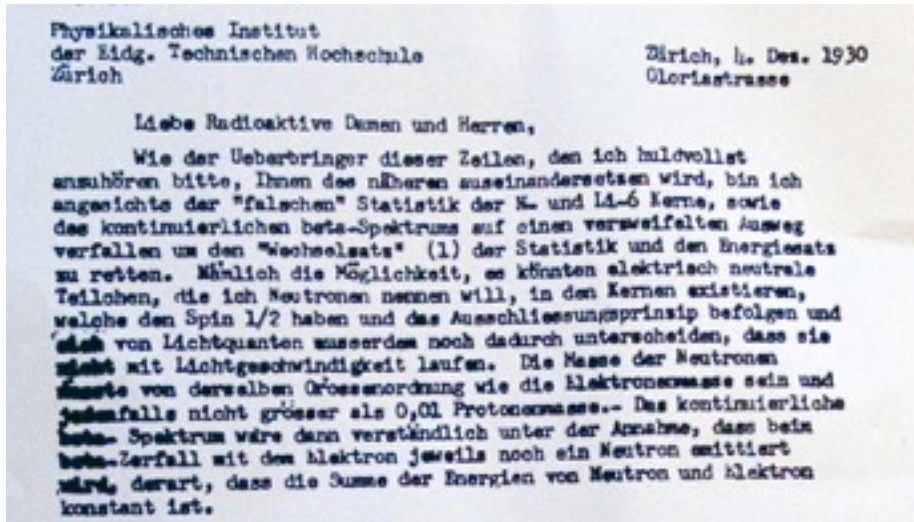
The invisible width of the Z (measured precisely at LEP) restricts the number of active (=interacting) neutrinos to

$$N_\nu = \frac{\Gamma_{inv}}{\Gamma_{\bar{\nu}\nu}} = 2.984 \pm 0.008$$

Note: Additional neutrinos can be present but they cannot partake of the SM interactions and are called sterile neutrinos.

*A REALLY brief
history of our
knowledge of
neutrinos*

The proposal of the “neutrino” was put forward by W. Pauli in 1930. [Pauli Letter Collection, CERN]



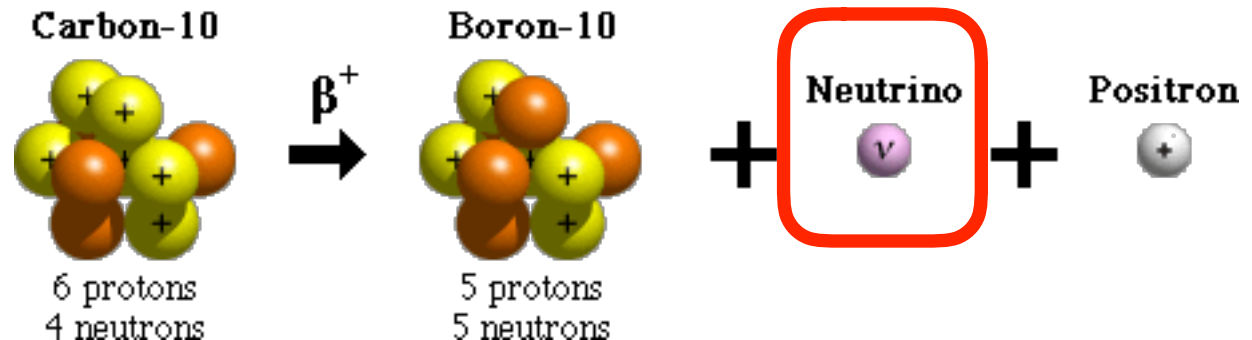
Dear radioactive ladies and gentlemen,

...I have hit upon a **desperate** remedy to save the ... energy theorem. Namely the possibility that there could exist in the nuclei **electrically neutral particles** that I wish to call neutrons, which have **spin $1/2$** ...

The **mass of the neutron must be ... not larger than 0.01 proton mass**. ...in β decay a neutron is emitted together with the electron, in such a way that the sum of the energies of neutron and electron is constant.



W. Pauli

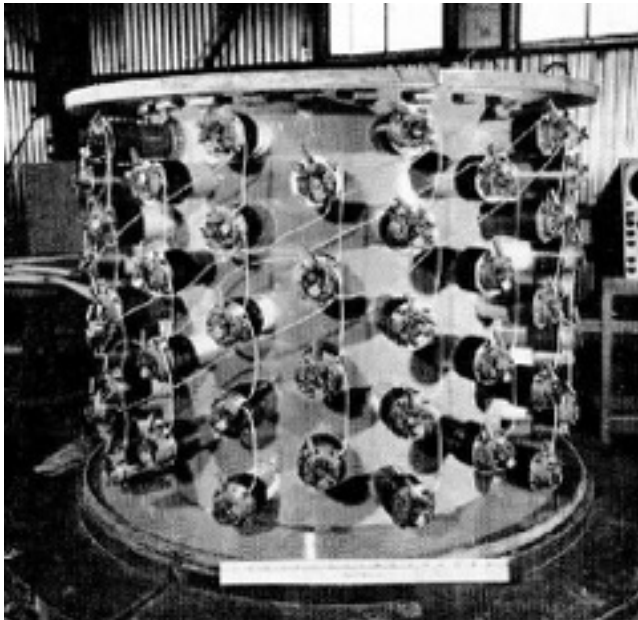


- Fermi, following E. Amaldi, used the name “**neutrino**” (little neutron) and later proposed the Fermi theory of beta decay.

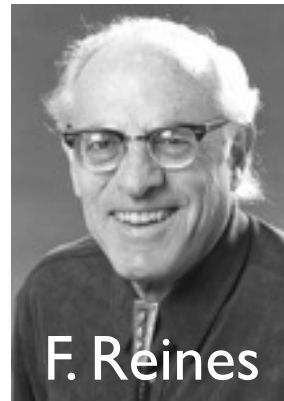


E. Fermi

- Reines and Cowan discovered neutrinos in 1956 using inverse beta decay.

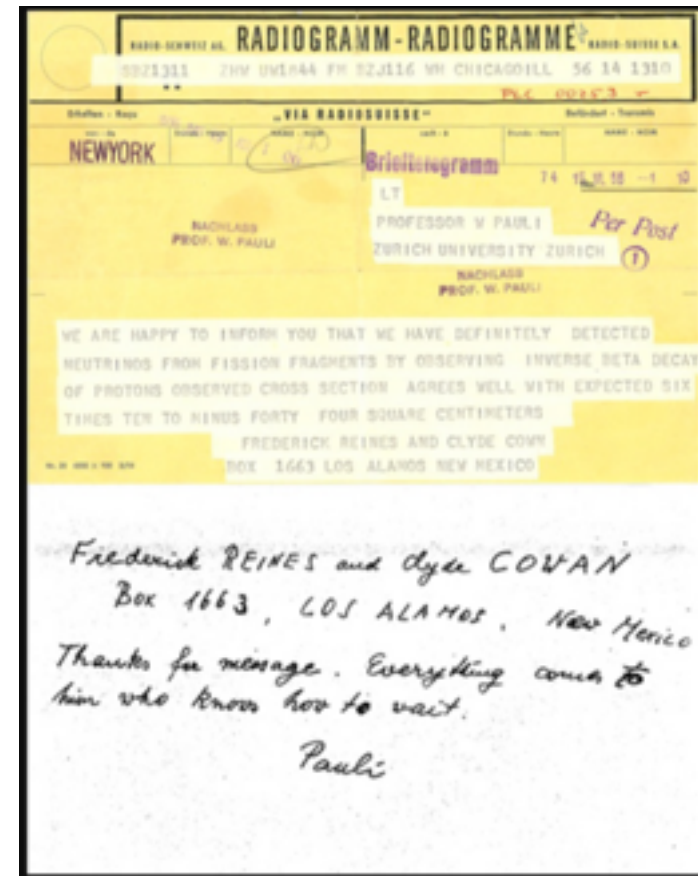


Savannah River experiment



F. Reines

The Nobel
Prize in
Physics
1995



- Fermi, following E. Amaldi, used the name “**neutrino**” (little neutron) and later proposed the Fermi theory of beta decay.



E. Fermi

- Reines and Cowan discovered neutrinos in 1956 using the Savannah River experiment.
- beta decay.

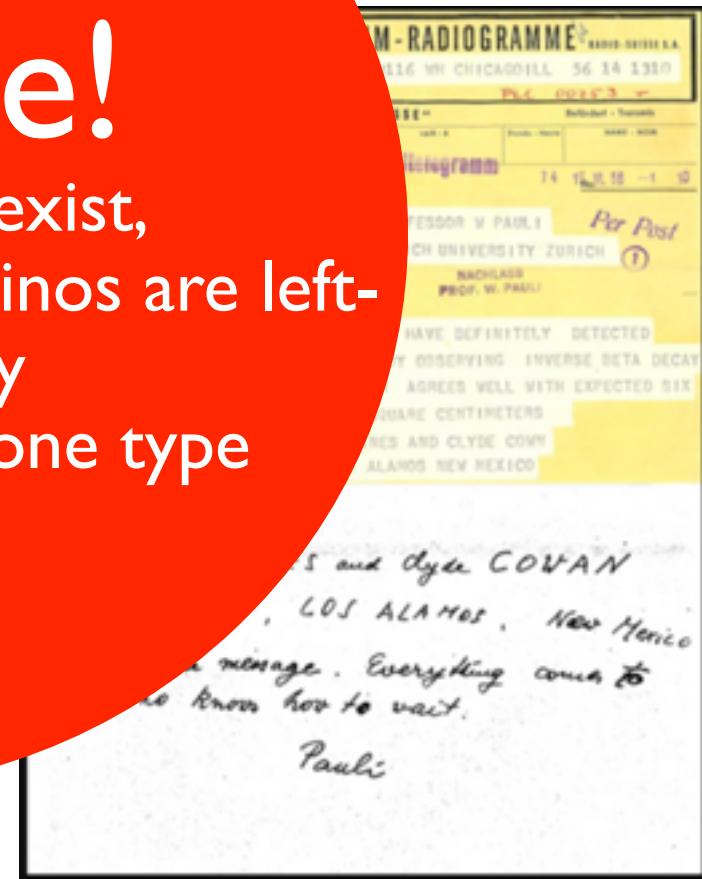


Savannah River experiment

Surprise!

- Neutrinos exist,
- P is violated and neutrinos are left-handed only
- there is more than one type

1995



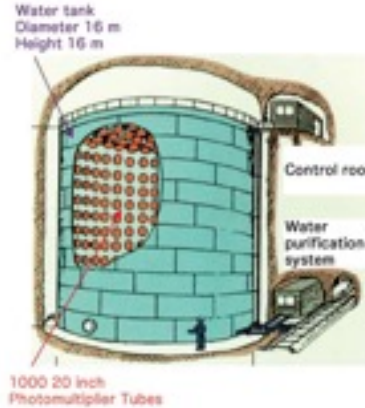
Searches for **astrophysical neutrinos** (from the Sun, SN and atmosphere) started, reporting anomalies.

Homestake



R. Davis Jr.

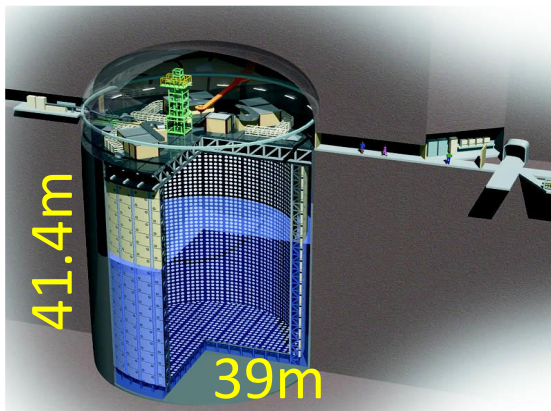
Kamiokande



M. Koshiba

Nobel
prize
in
2002

This led to the discovery of neutrino oscillations.



Super-Kamiokande



T. Kajita



SNO

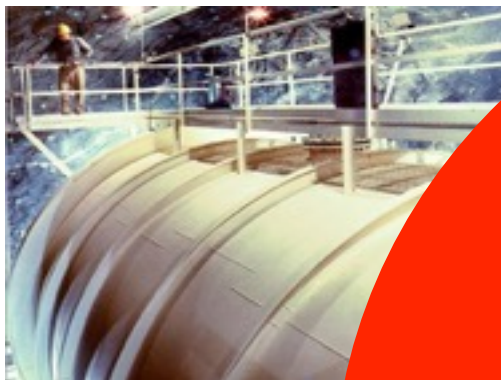


A. McDonald

Nobel Prize in
Physics 2015

Searches for **astrophysical neutrinos** (from the Sun, SN and atmosphere) started, reporting anomalies.

Homestake



Kamio

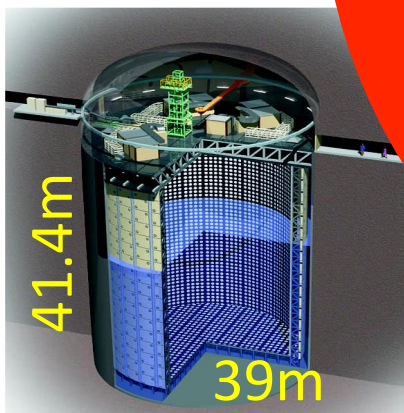


Nobel
prize
in
2002

This

Surprise!
Neutrinos oscillate

ions.



Super-Kamiokande



SNO



Nobel Prize in
Physics 2015

*Neutrino
oscillations:
a quantum
mechanical
phenomenon*

The first idea of neutrino oscillations was put forward by B. Pontecorvo in 1957.



Бруно Понтекорво

Neutrinos can change flavour while traveling. This is an eminently quantummechanical effect, similar to other observed ones, such as spin precession. It has an oscillatory behaviour.



Neutrino mixing

Mixing between two basis (flavour and mass) is described by the *Pontecorvo-Maki-Nakagawa-Sakata* matrix:

$$\nu_{\alpha} = \sum_i U_{\alpha i} \nu_i$$

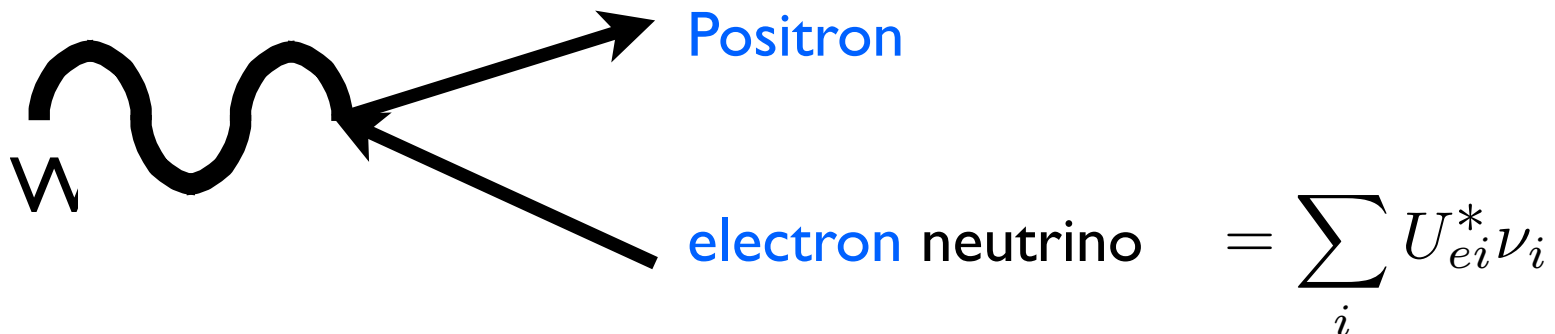
↑
←

Flavour field
Mass field

which enters in the CC interactions

$$\mathcal{L}_{CC} = -\frac{g}{\sqrt{2}} \sum_{k\alpha} (U_{\alpha k}^* \bar{\nu}_{kL} \gamma^{\rho} l_{\alpha L} W_{\rho} + \text{h.c.})$$

This implies that in an interaction with an electron, the corresponding (anti-)neutrino will be produced, as a superposition of different mass eigenstates.



- **2-neutrino mixing** matrix depends on 1 angle only. The phases get absorbed in a redefinition of the leptonic fields (a part from 1 Majorana phase).

$$\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$$

- **3-neutrino mixing** matrix has 3 angles and 1(+2) CPV phases (neutrino - antineutrino $U \rightarrow U^*$).

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_{21}/2} & 0 \\ 0 & 0 & e^{i\alpha_{31}/2} \end{pmatrix}$$

Let's assume that a **muon neutrino** is produced

$$|\nu, t = 0\rangle = |\nu_\mu\rangle = \sum_i U_{\mu i}^* |\nu_i\rangle$$

The time-evolution is described by the Schroedinger equation with free Hamiltonian:

$$\mathcal{H} = \begin{pmatrix} \blacksquare & \blacksquare \\ \blacksquare & \blacksquare \end{pmatrix} \xrightarrow{\nu_e, \nu_\mu} i \frac{d}{dt} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

flavour basis mass basis

The solution is

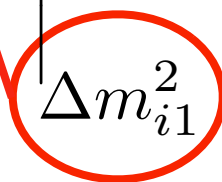
$$|\nu, t\rangle = \sum_i U_{\mu i}^* e^{-iE_i t} |\nu_i\rangle$$

At **detection** one projects over the flavour state. The **probability of oscillation** is

$$P(\nu_\mu \rightarrow \nu_e) = |\langle \nu_\mu | \nu_e \rangle|^2 = \left| \sum_{ij} U_{ei}^* U_{\mu j} e^{-iE_i t} \langle \nu_j | \nu_i \rangle \right|^2$$

very relativistic

neutrinos: $E_i \simeq p + \frac{m_i^2}{2p}$

$$= \left| \sum_i U_{ei}^* U_{\mu i} e^{-i \frac{m_i^2 - m_1^2}{2E} t} \right|^2$$


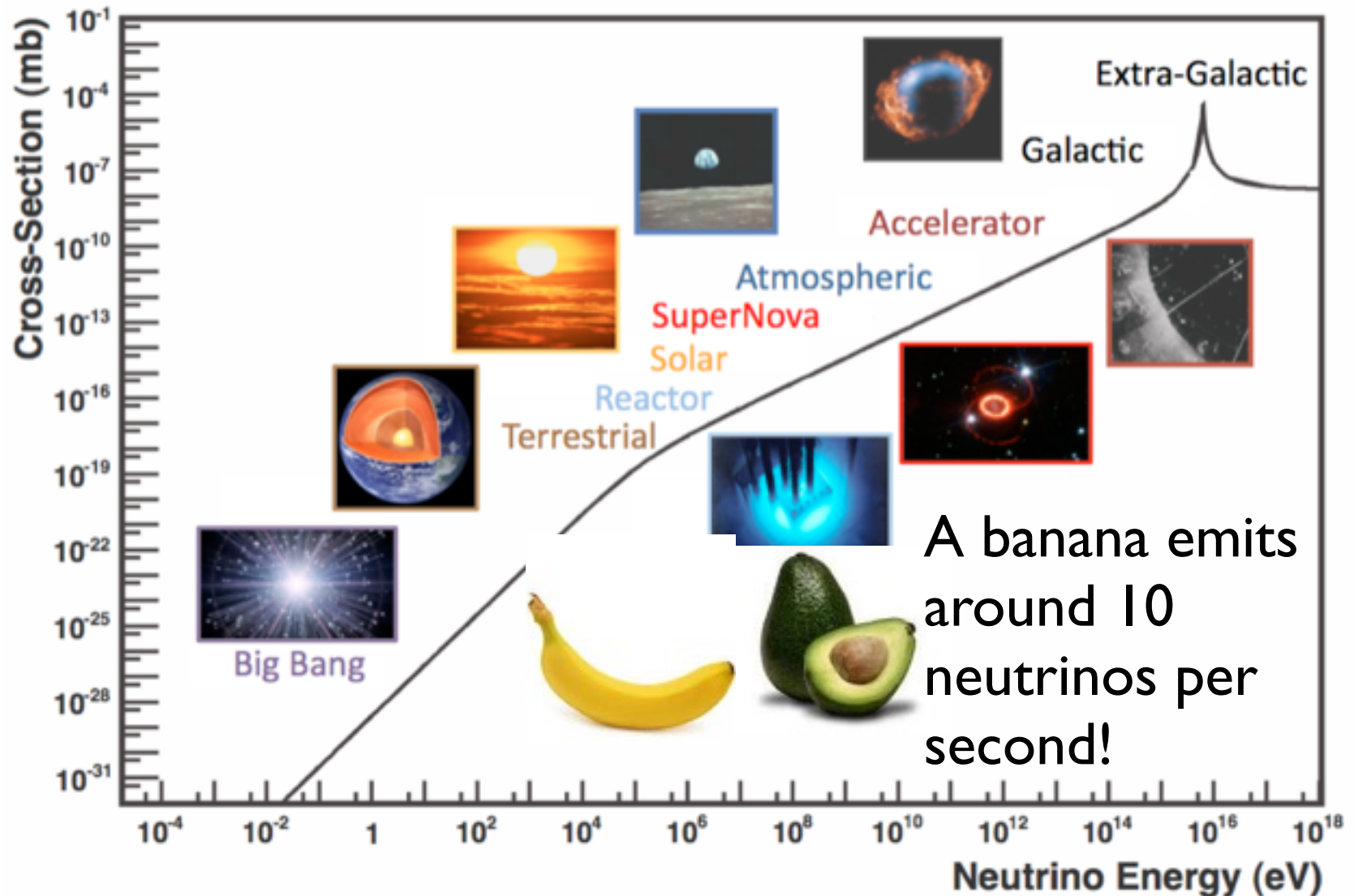
$$P(\nu_\alpha \rightarrow \nu_\beta) = \left| \sum_i U_{\alpha i} U_{\beta i}^* e^{-i \frac{\Delta m_{i1}^2}{2E} L} \right|^2$$

- **neutrinos have mass** (as the different components of the initial state need to propagate with different phases)
- **neutrinos mix** (If they do not mix the flavour eigenstates are also eigenstates of the propagation Hamiltonian.)



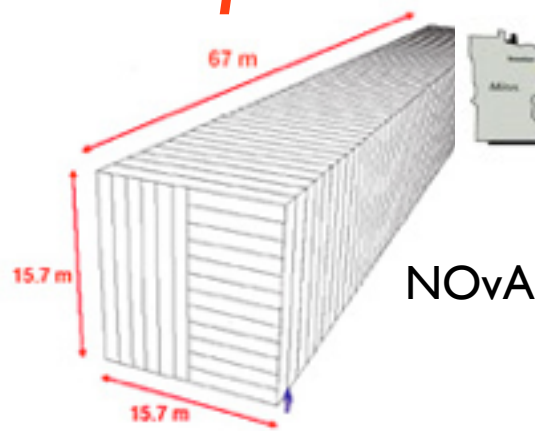
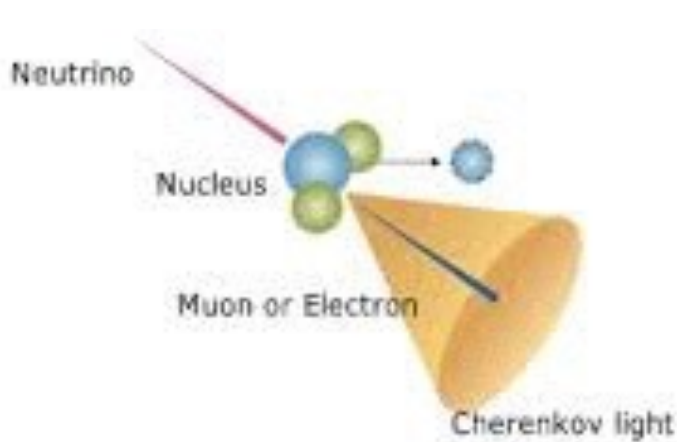
Current knowledge of neutrino properties

Neutrino sources

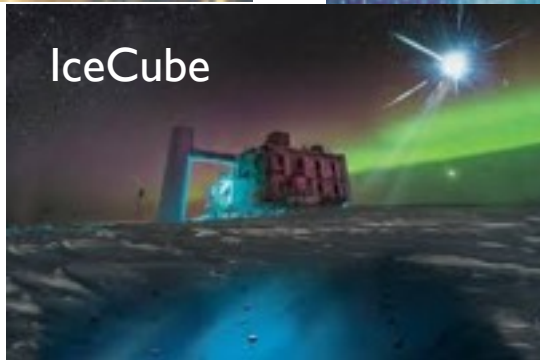
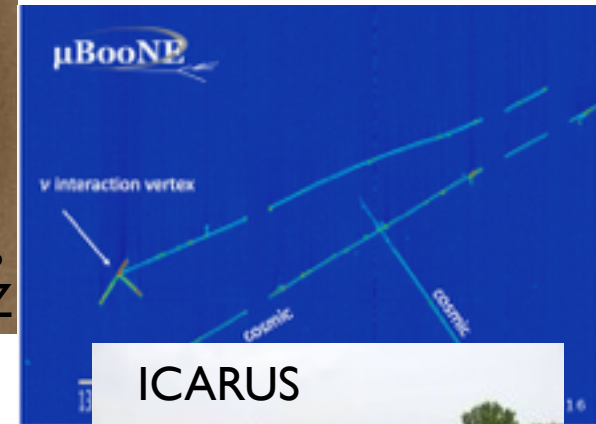
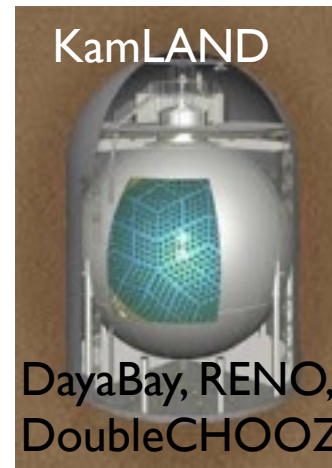


J. Formaggio and S. Zeller, I 305.75 I 3

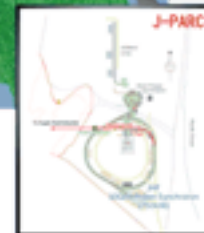
Current/past experiments



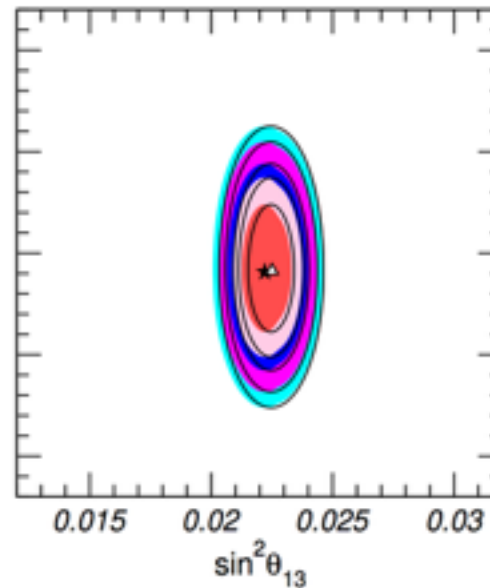
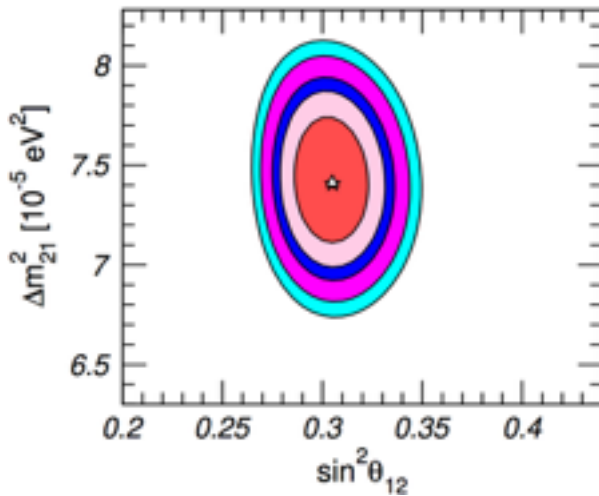
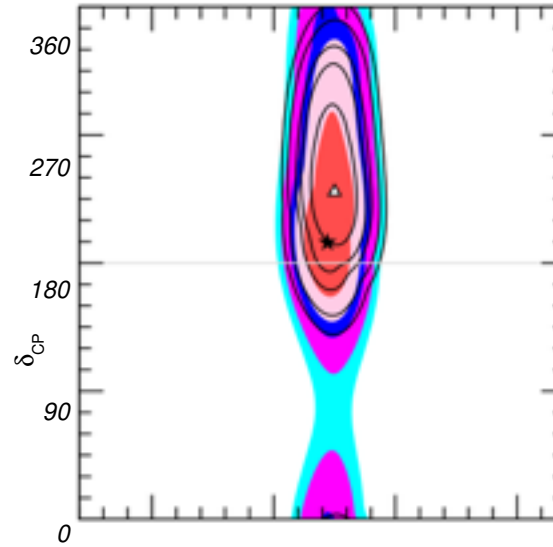
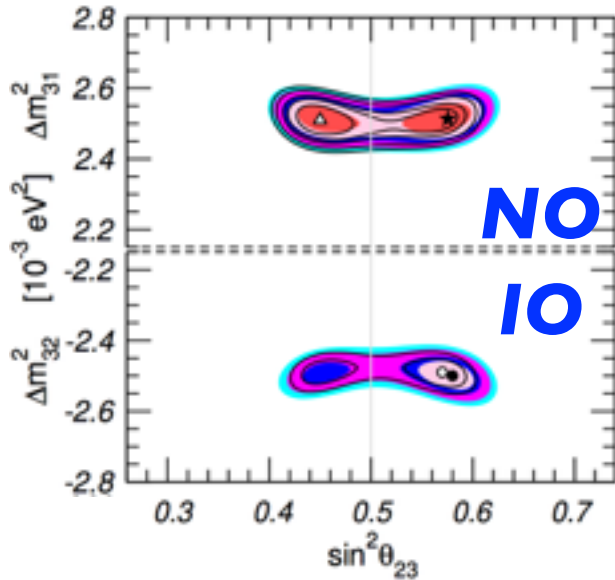
Credit: Super-Kamiokande



T2K



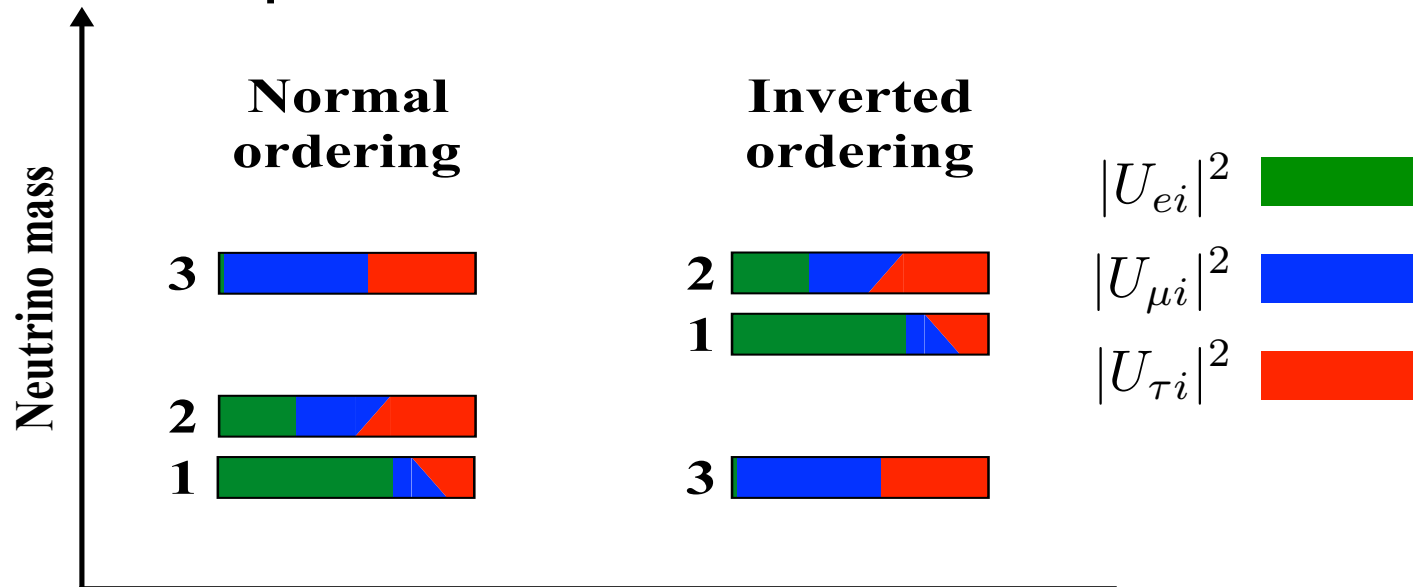
NuFIT 5.1 (2021)



Current status:

- 2 mass squared differences
- 3 sizable mixing angles,
- mild hints of CPV
- mild indications in favour of NO

$\Delta m_{21}^2 \ll \Delta m_{31}^2$ implies at least 3 massive neutrinos.



Fractional flavour content of massive neutrinos

$$m_1 = m_{\min}$$

$$m_2 = \sqrt{m_{\min} + \Delta m_{21}^2}$$

$$m_3 = \sqrt{m_{\min} + \Delta m_{31}^2}$$

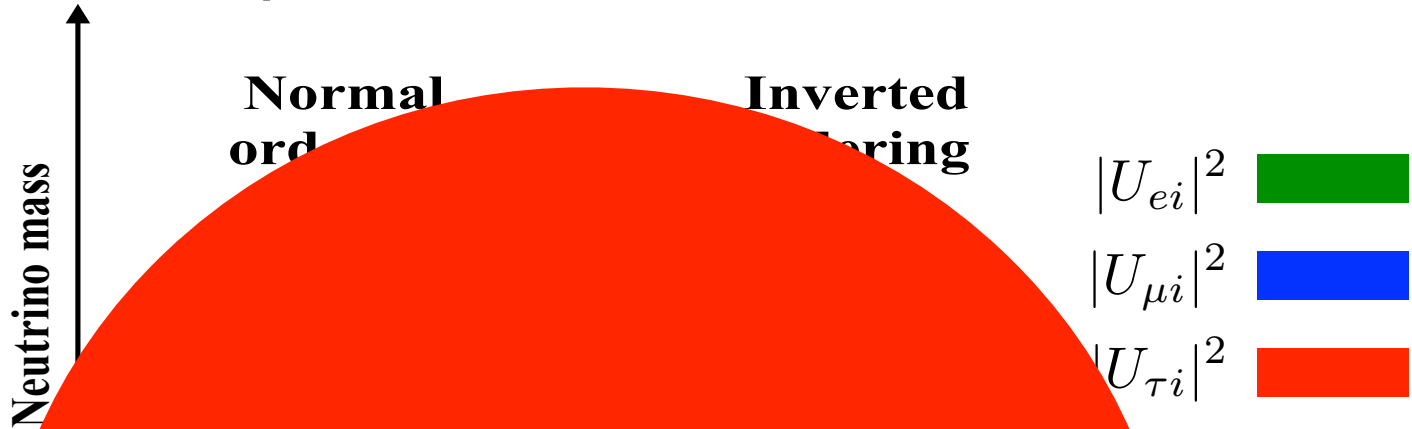
$$m_3 = m_{\min}$$

$$m_1 = \sqrt{m_{\min} + |\Delta m_{32}^2| - \Delta m_{21}^2}$$

$$m_2 = \sqrt{m_{\min} + |\Delta m_{32}^2|}$$

- Neutrino masses: ordering and scale unknown. **They are much smaller and less hierarchical than charged fermions.**
- **Possibly large CPV: fundamental question**
- **Large mixing angles** (differently from quark sector).

$\Delta m_{21}^2 \ll \Delta m_{31}^2$ implies at least 3 massive neutrinos.



Surprise!
 Very small and not-very-hierarchical masses
 Large mixing angles

- Neutrino masses are much smaller and more degenerate than those of charged fermions. They are much smaller and more degenerate than those of charged fermions.
- Possibly large CPV: fundamental question
- Large mixing angles (differently from quark sector).



Questions for the future

What do we still need to know?

- What is the **nature** of neutrinos? Dirac vs Majorana?
- What are the values of the **masses**? Absolute scale (KATRIN, ...?) and the ordering.
- Is there **CP-violation**? Its discovery in the next generation of LBL depends on the value of delta.
- What are the **precise values** of mixing angles? Do they suggest an underlying pattern?
- Is the **standard picture** correct? Are there NSI? Sterile neutrinos? Other effects?

2020

2025

2030

2035

LBL osc.

T2K
NOvA

LBNF-DUNE
T2HK (T2HKK)

ESSnuSB?,
nufactory?

SBL osc.

SBL reactor,...
MicroBooNE

SBN

LBNF-DUNE ND
T2HK ND
???

Other osc.

SK, Borexino,
LBL detectors

JUNO

DUNE
HK

Theia???

Direct mass

KATRIN

Project 8

DBD0n u

KamLAND-Zen
GERDA
CUORE

LEGEND-200

NEXT-100, nEXO...

LEGEND-1000
CUPID
NEXT-HD, PANDAX...

Next-
next
gen?

UHE

IceCube

IceCubeGen2
ORCA, KM3Net

Neutrino nature and neutrinoless DBD

Neutrinos can be **Majorana** or **Dirac** particles. In the SM only neutrinos can be Majorana as they are neutral.

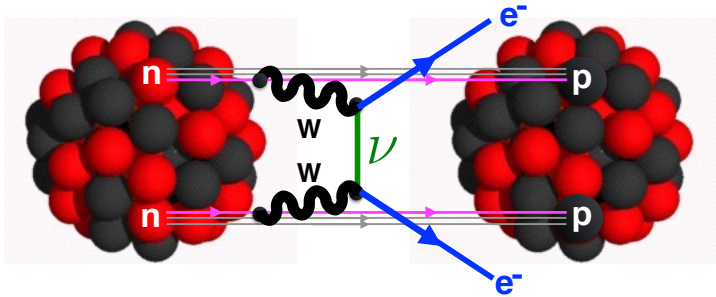
Majorana condition

$$\nu = C\bar{\nu}^T$$

The nature of neutrinos is linked to the conservation of **Lepton number (L)**.

- This is crucial information to unveil the **Physics BSM: with or without L-conservation?** Lepton number violation is a necessary condition for **Leptogenesis**.
- Tests of LNV:
 - At low energy, neutrinoless double beta decay,
 - LNV tau and meson decays, collider searches.

Neutrinoless double beta decay $((A,Z) \rightarrow (A, Z+2) + 2 e^-)$ experiments can proceed via the exchange of massive Majorana neutrinos.



The half-life time depends on neutrino properties

$$(T_{0\nu}^{1/2})^{-1} \propto |M_{NME}|^2 |m_{\beta\beta}|^2$$

SP, CERN Courier, Jul 2016

- The effective Majorana mass parameter:

$$|m_{\beta\beta}| \equiv |m_1|U_{e1}|^2 + m_2|U_{e2}|^2 e^{i\alpha_{21}} + m_3|U_{e3}|^2 e^{i\alpha_{31}}|$$

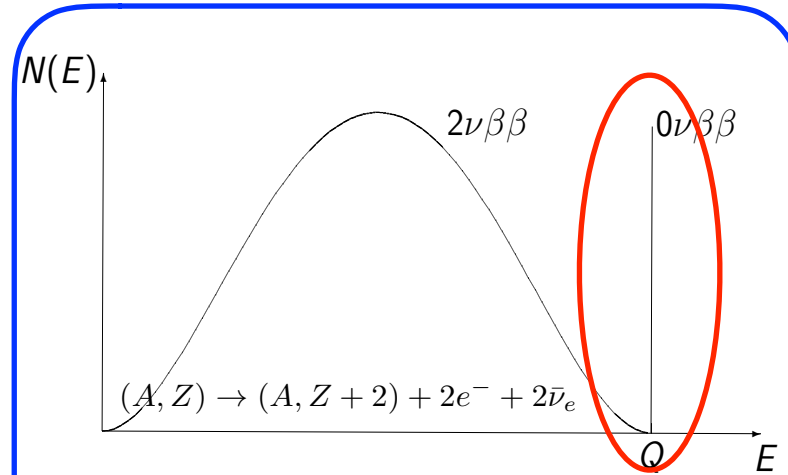
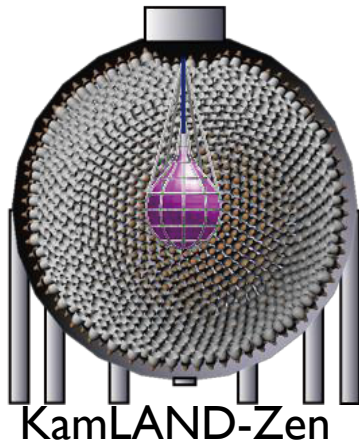
Mixing angles (known)

CPV phases (unknown)

Neutrinoless double beta decay is a very rare process: $T_{1/2} > 10^{26}$ yrs.

$$T_{0\nu} \propto \sqrt{\frac{M t}{B \Delta E}}$$

ton-scale
 $< 1\%$ at Q_{bb}
 < 1 cts/yr/ton/ROI

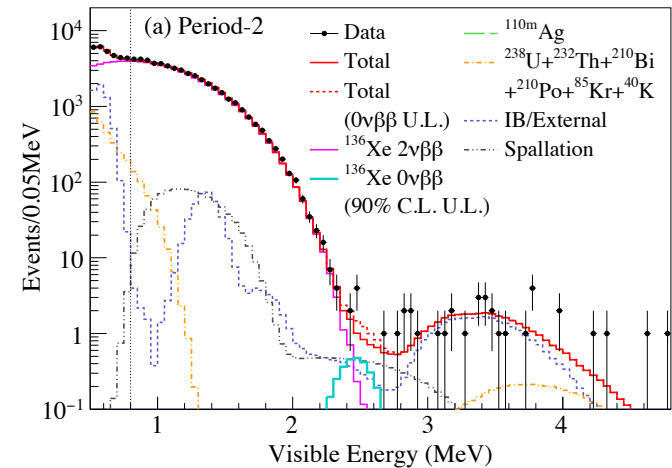


News

Roman ingots to shield particle detector

Lead from ancient shipwreck will line Italian neutrino experiment.

Nature



KamLAND-Zen, PRL 117 (2016)

KamLAND-Zen Loaded LSc with 380 kg ^{136}Xe ,
 $T_{1/2} > 1.07 \times 10^{26}$ yrs (90% C.L.), $m_{bb} < 61-165$ meV

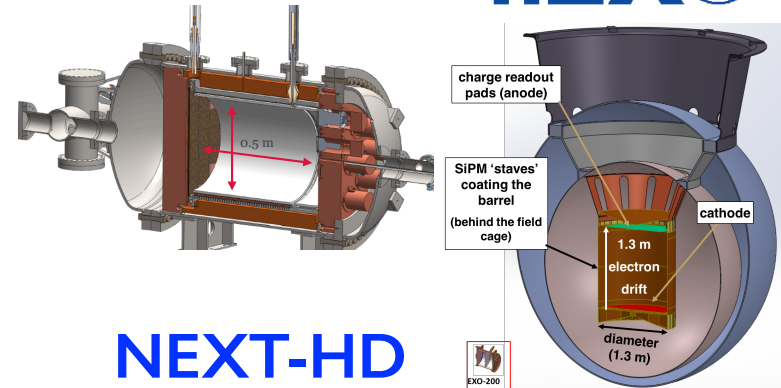
EXO-200 ~75 kg LXe TPC, $T_{1/2} > 3.5 \times 10^{25}$ yrs

GERDA 31 kg (enriched) ^{76}Ge , $T_{1/2} > 1.8 \times 10^{26}$ yrs

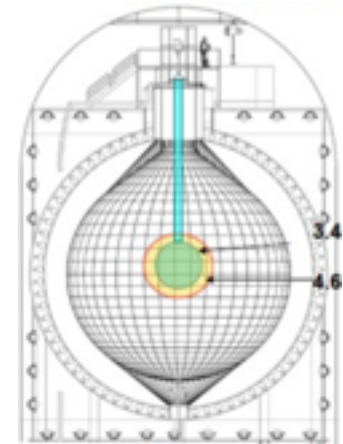
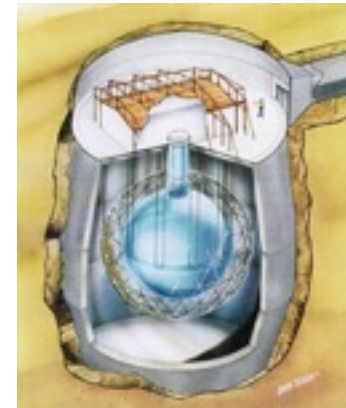
MAJORANA 26.0 kg yrs, $T_{1/2} > 0.27 \times 10^{26}$ yrs

CUORE ^{130}Te , ~206 kg, $T_{1/2} > 2.2 \times 10^{25}$ yrs

nEXO



NEXT-HD



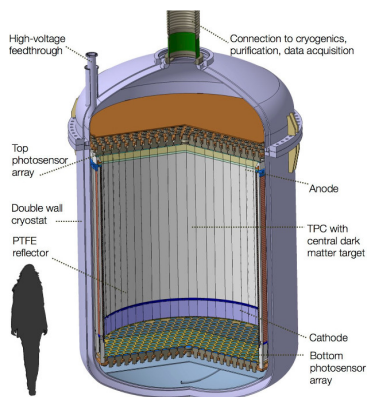
SNQ+

KamLAND2-Zen

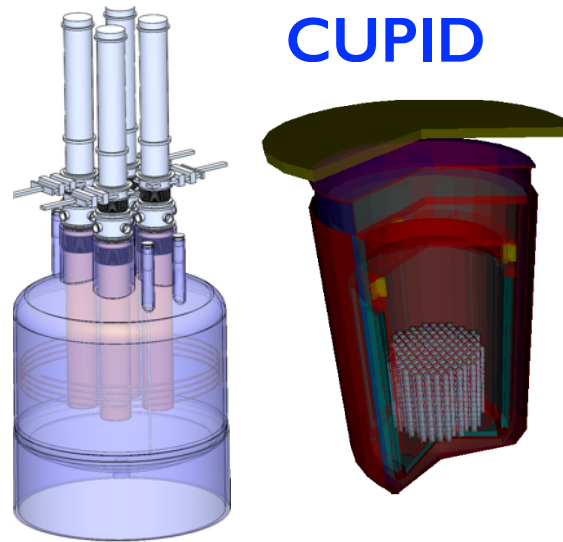
HPXe: PANDAX-III

Bolometers: AMoRE

DARWIN



CUPID

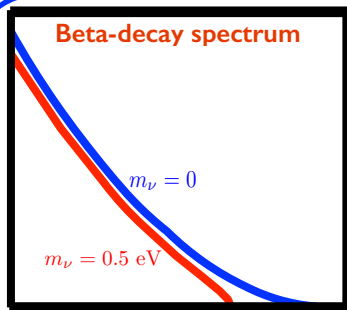


LEGEND Large Enriched Germanium Experiment for Neutrinoless $\beta\beta$ Decay

The ultimate goal of next generation is $m_{bb} \sim 15-20$ meV.

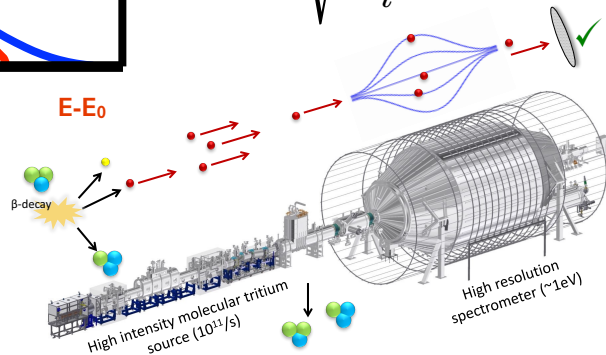
Measuring neutrino masses

- Absolute mass scale.

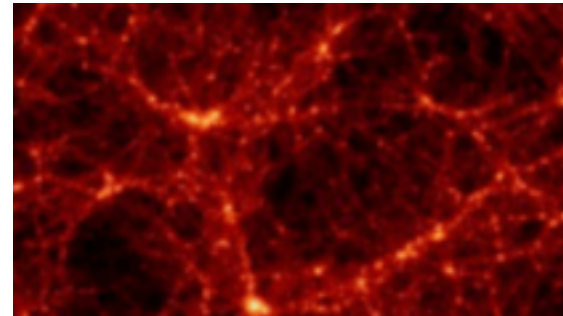


Beta decay

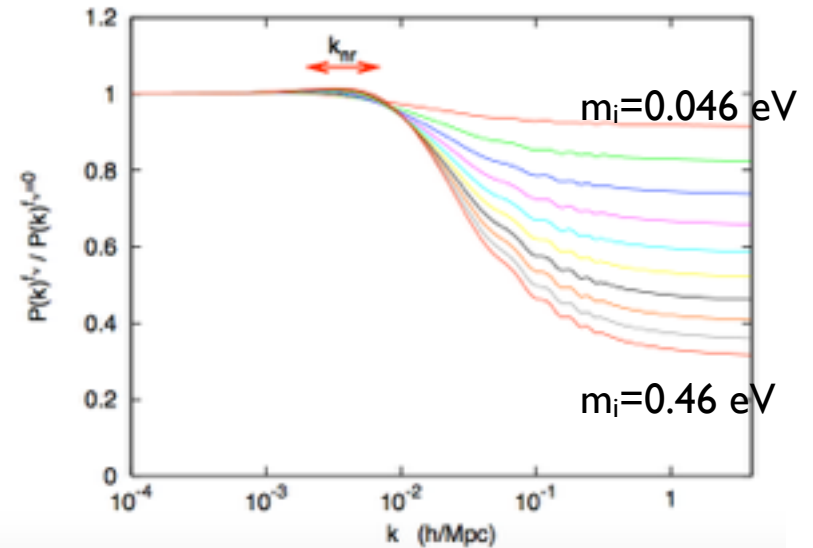
$$m_\beta \equiv \sqrt{\sum_i U_{ei}^2 m_i^2}$$



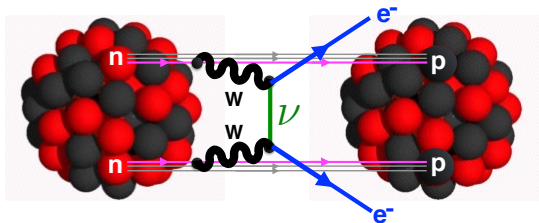
Cosmology



$$\sum_i m_i$$



Neutrinoless dbeta decay

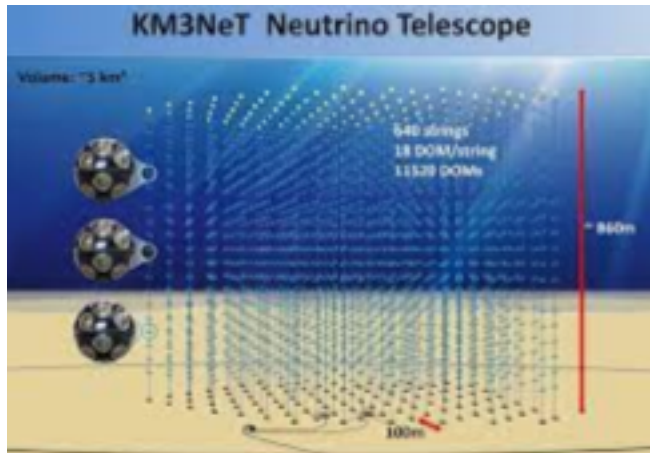


$$m_{\beta\beta} = f(m_i, \alpha_{21}, \alpha_{31}, \delta)$$

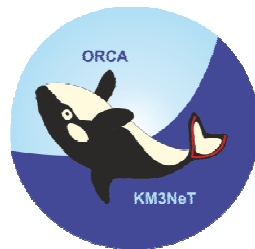
J. Lesgourgues and S. Pastor, Phys. Rep. 429

- **Mass ordering** via **neutrino oscillation in matter or in vacuum** (JUNO). Discovery expected within 10 years thanks to relatively large θ_{13} .

Atm neutrinos



Exploit matter effects in large detectors.



Long baseline neutrino oscillation experiments

JUNO

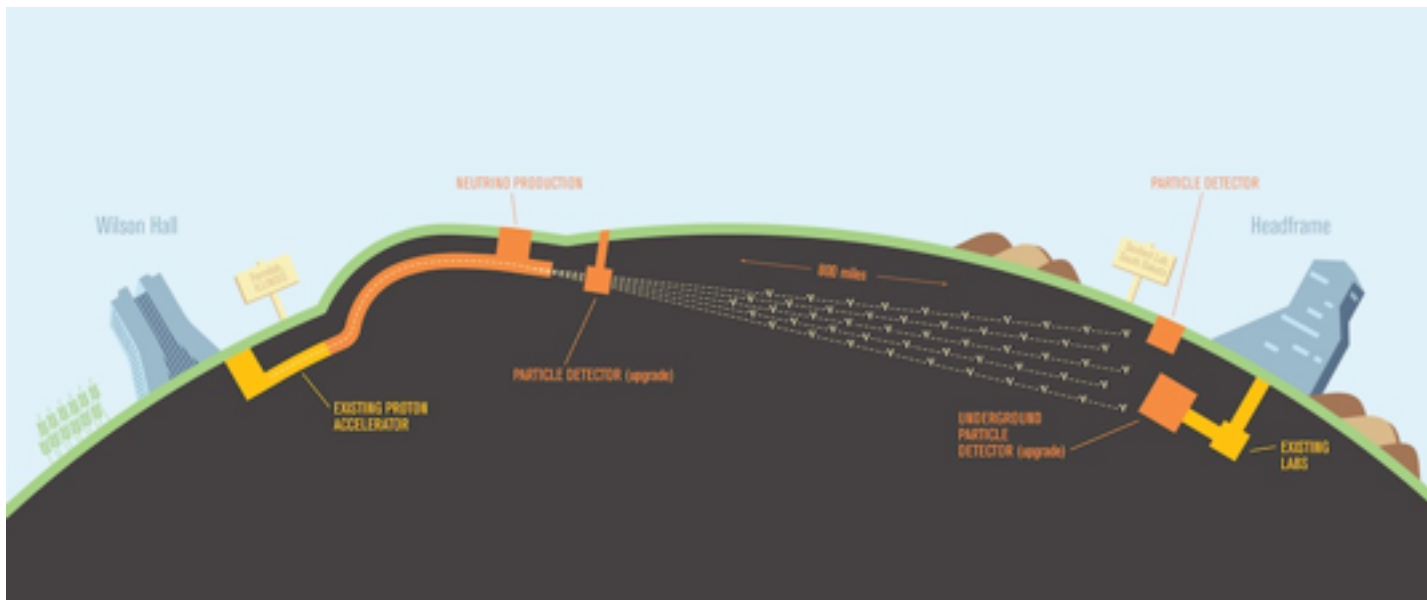


JUNO Coll., 2104.02565

JUNO uses a 20kton LSc detector to detect reactor nus at a baseline of ~ 60 km. Excellent energy resolution is needed. Due to start in 2023.

Long baseline oscillations: mass ordering and CPV

Long baseline neutrino oscillation experiments (T2K, NOvA, DUNE, T2HK) study the subdominant channels muon to electron (anti)neutrino.

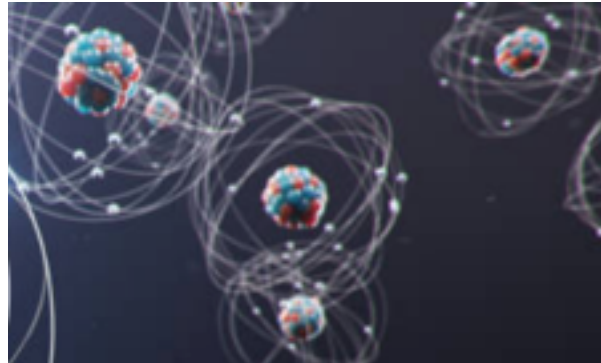


Credit:
Symmetry
magazine

Muon neutrinos are produced in pion decays at accelerator complexes (JPARC, Fermilab) and then travel 100s Km to highly capable very large neutrino detectors.

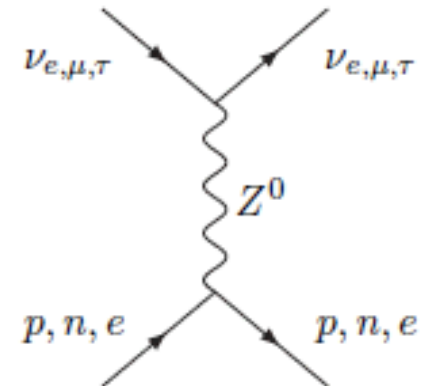
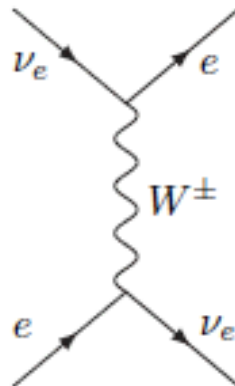
- When neutrinos travel through a medium, e.g. Earth, they interact with the background of e, p and n.

ν_α

Matter (e, p, n)

via SM processes:

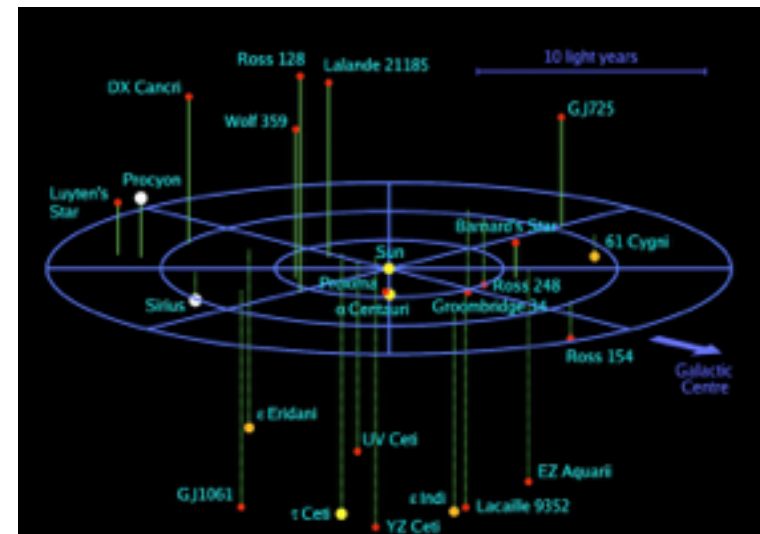


Inelastic scattering in which the neutrino changes direction/energy or converts into a charged lepton is very very very rare.



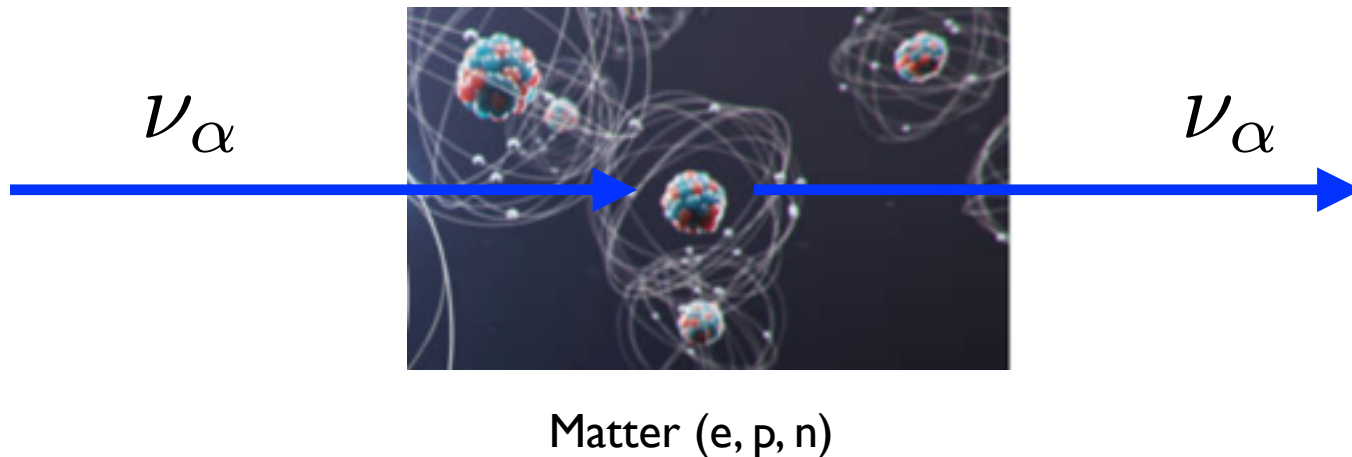
Matter (e, p, n)

At the neutrino energy of beta decays, the average distance before an interaction (mean free path) in water (1 g/cm^3) is $d = 200$ lightyears!



ESO

Forward elastic scattering gives neutrinos a flavour-diagonal effective mass and affects oscillations.



- The background is CP and CPT violating, e.g. the Earth contains only particles and not antiparticles, and the resulting oscillations are CP and CPT violating (therefore, neutrinos and antineutrinos are affected differently).

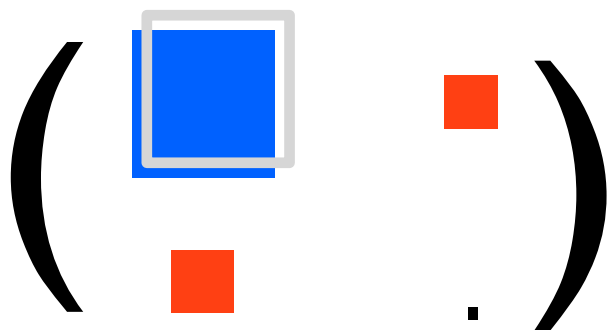
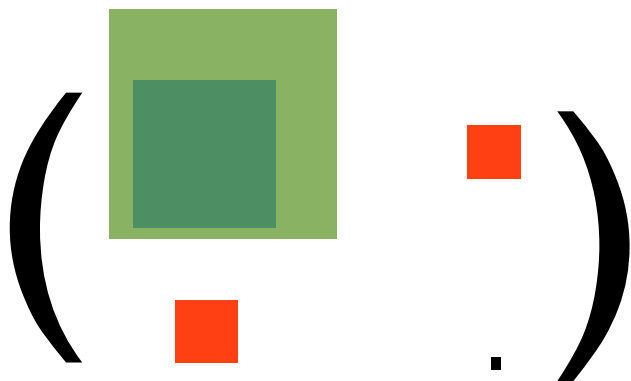
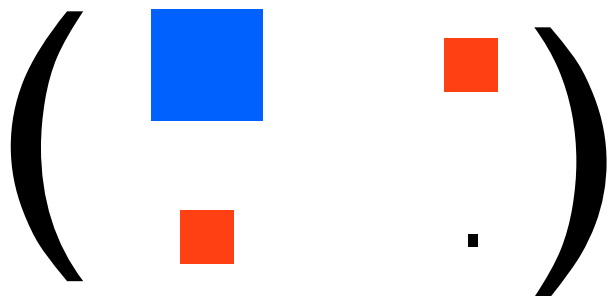
Effective Hamiltonian

$$\begin{pmatrix} \blacksquare & \blacksquare \\ \blacksquare & \cdot \end{pmatrix}$$

Mixing angle

$$\tan 2\theta \sim \frac{\text{vacuum} \blacksquare}{\blacksquare}$$

Effective Hamiltonian



Mixing angle

vacuum

$$\tan 2\theta \sim \frac{2 \text{ (red square)}}{\text{ (blue square)}}$$

matter suppression (Sun, SN)

$$\tan 2\theta^M \sim \frac{2 \text{ (red square)}}{\text{ (blue square)} + \text{ (green square)}} \ll \tan 2\theta$$

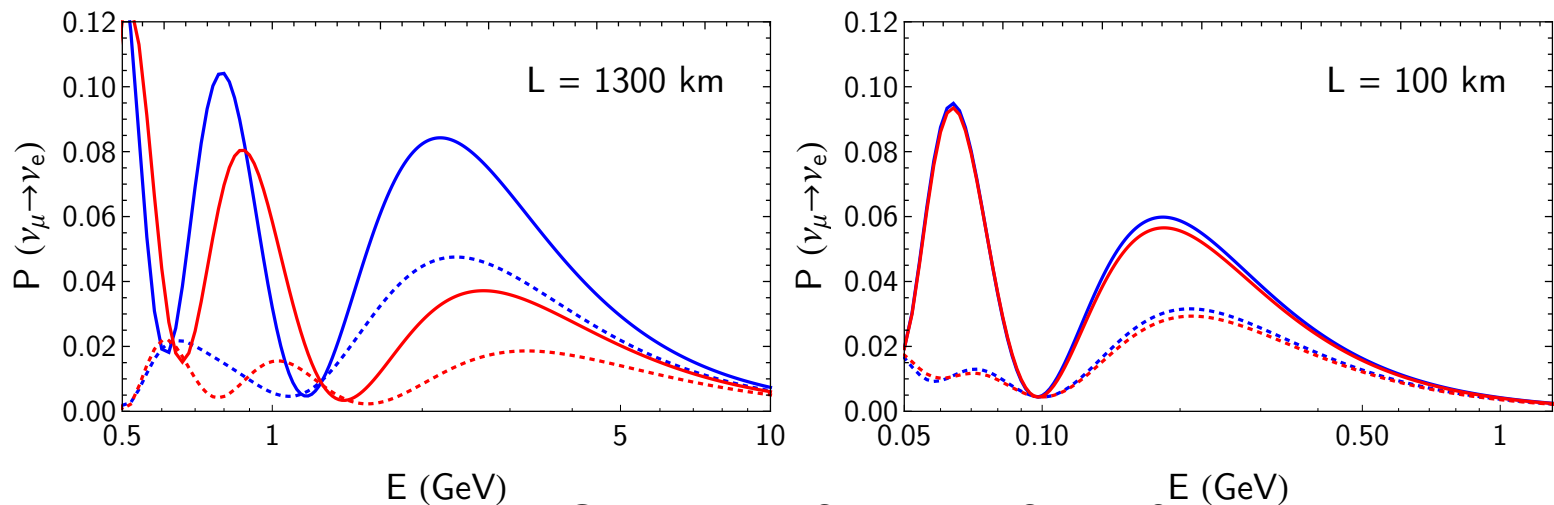
MSW resonance (Sun, SN)

$$\tan 2\theta^M \sim \frac{2 \text{ (red square)}}{\text{ (blue square)} - \text{ (white square)}} \sim \infty$$

- Matter effects are described by a potential V in the effective Hamiltonian which determines the time evolution.

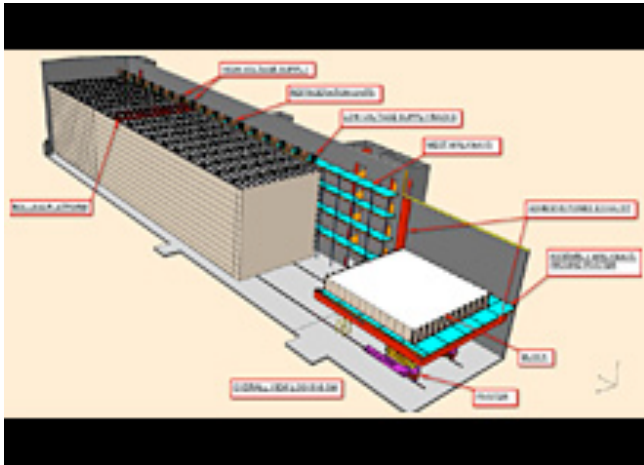
$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} -\frac{\Delta m^2}{4E} \cos(2\theta) + \sqrt{2} G_F N_e & \frac{\Delta m^2}{4E} \sin(2\theta) \\ \frac{\Delta m^2}{4E} \sin(2\theta) & \frac{\Delta m^2}{4E} \cos(2\theta) \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

- The oscillations can be enhanced or suppressed (depending on the sign of $\Delta m^2 =$ mass ordering).



Present/Future LBL exp

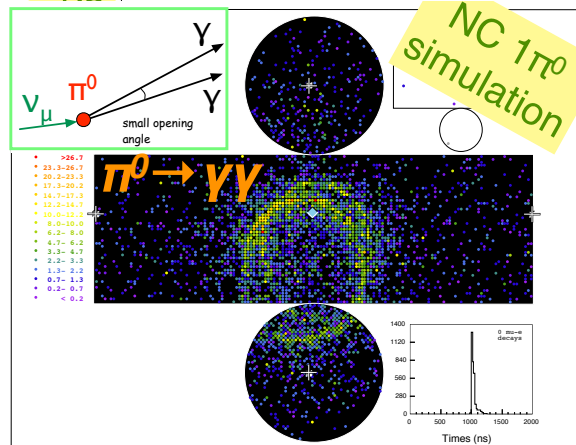
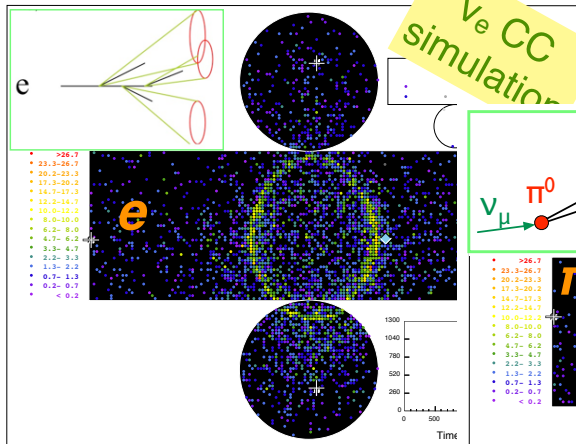
DUNE: 1300 km on-axis
(20)-40 kton LAr detector



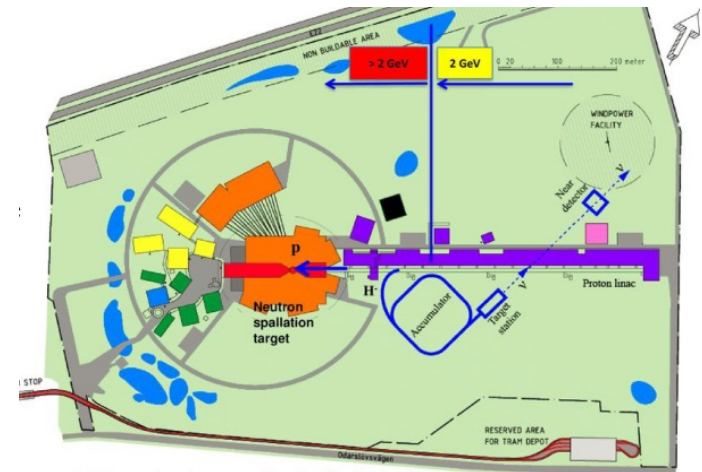
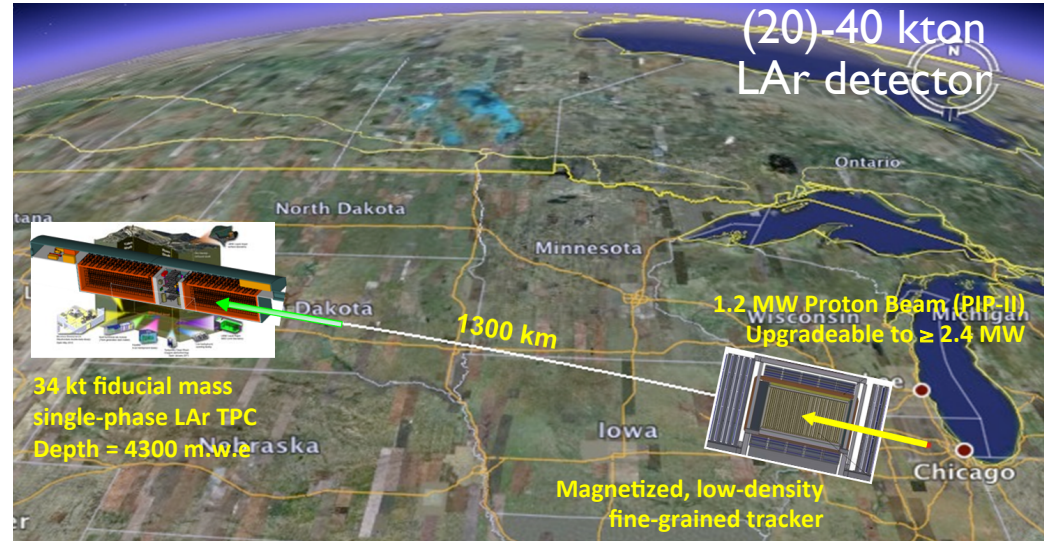
NOVA: 810 km off-axis
~14 kton plastic scintillator detector

T2K: 295 km off-axis
~22.5 kton WC detector

T2HK: 295 km off-axis
~1 Mton WC detector



M. Shiozawa, for T2HK coll., NuPhys 2014

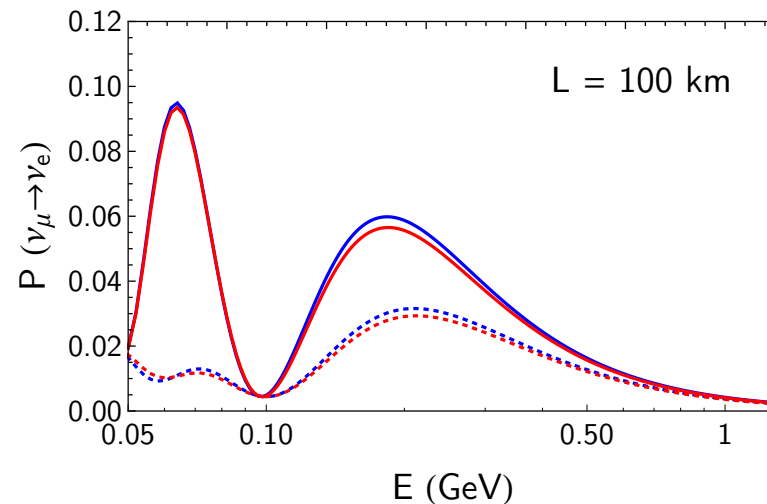


~1 BEuros for the neutrino facility including detector

ESSnuSB: 300-500 km
~0.5 Mton WC detector
second oscillation maximum

CP-violation in LBL experiments

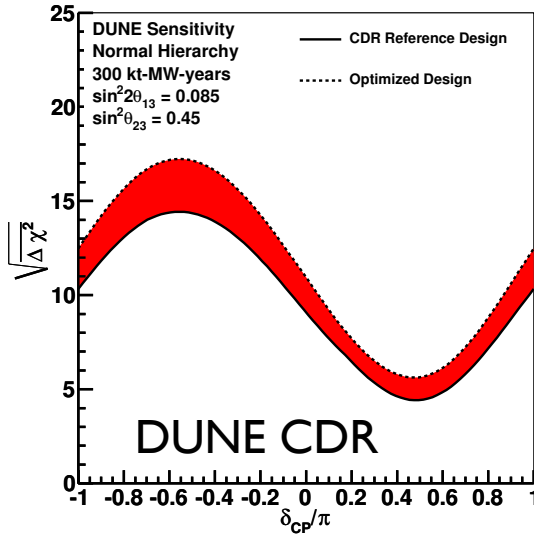
In presence of CP violation (complex mixing matrix U), neutrinos and antineutrinos behave differently.



P. Coloma and SP, World Scientific

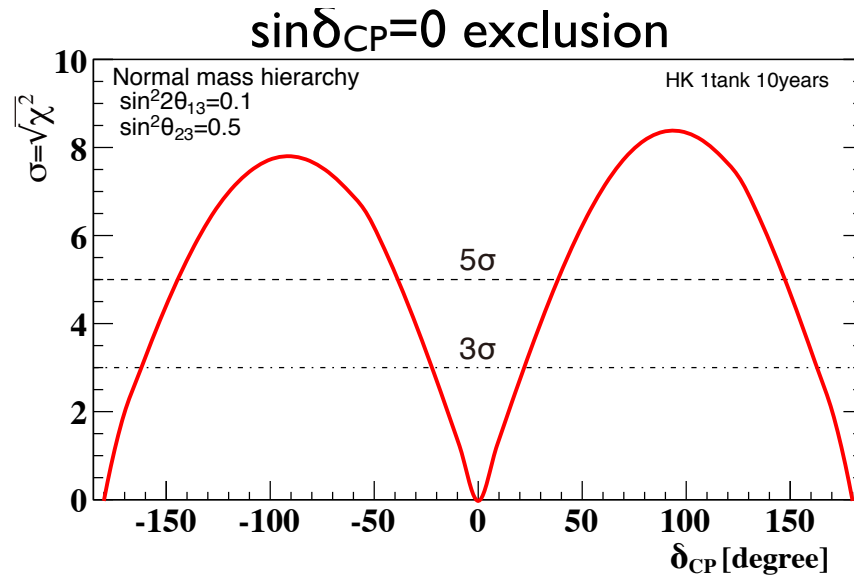
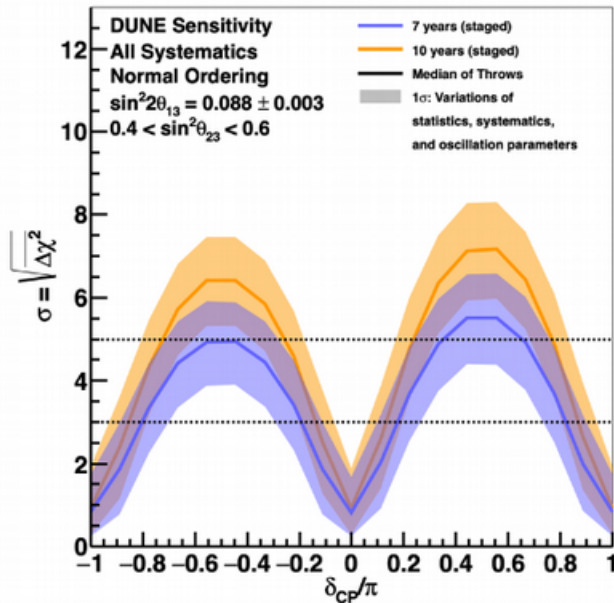
- By comparing neutrino and antineutrino oscillations one can deduce the presence of leptonic CPV.
- One needs to disentangle true CPV from matter effects.
- CPV effects are more pronounced at **low energy**.

Mass ordering and CPV sensitivity



- The mass ordering will be discovered by the end of the decade by DUNE, with strong indications by ~ 2025 by JUNO, KM3Net, Atm neutrino exp.
- If CPV is \sim maximal, we will know it in the next ~ 10 -15 years.

True Normal Ordering 04



*Beyond 3-neutrino
mixing?*

Anomalies in the neutrino sector

There are hints beyond standard 3 neutrino mixing.

$$\pi^+ \rightarrow \mu^+ + \nu_\mu ,$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

$\bar{\nu}_e$

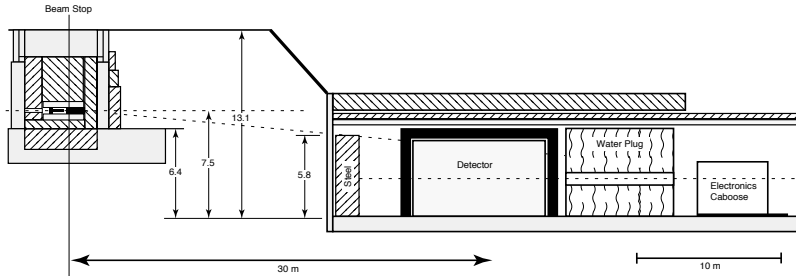
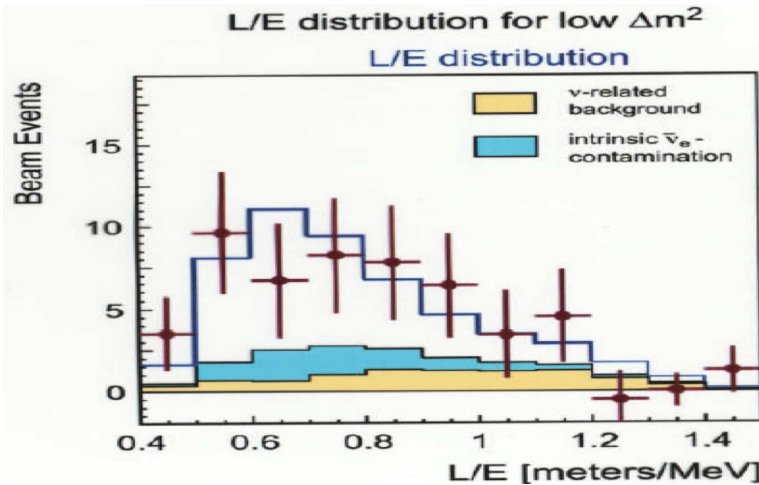
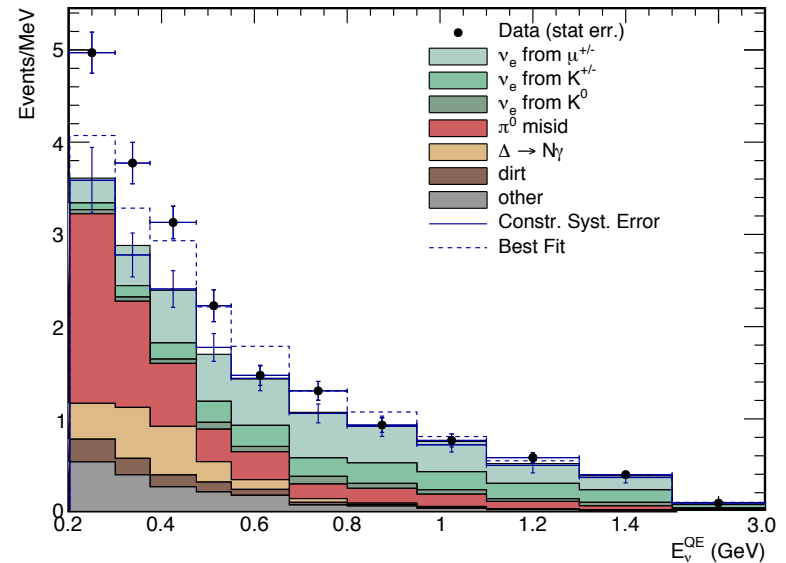
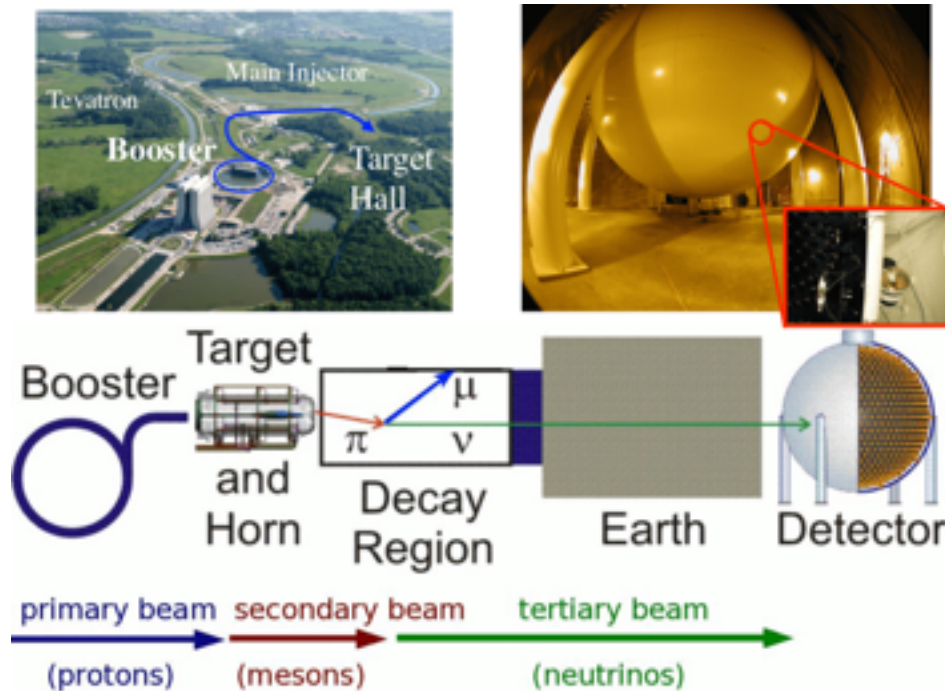


FIG. 1. Detector enclosure and target area configuration, elevation view

LSND reported the appearance of electron anti-neutrinos (inverse beta decay) at short distance (~ 30 m) from muon decays (DAR). A 3.8 sigma effect, not confirmed by KARMEN.

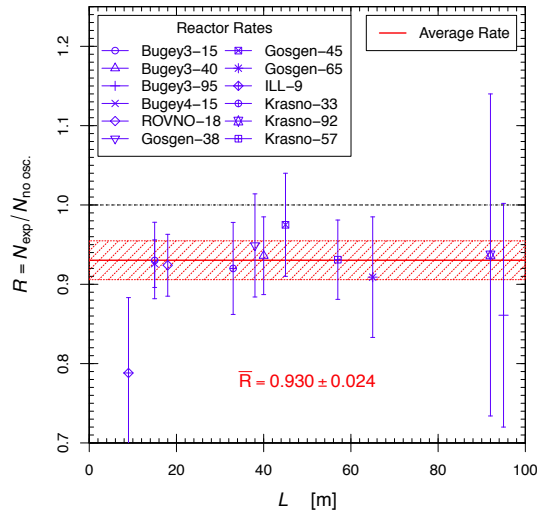


MiniBooNE was designed to test the LSND results.
 $\langle E \rangle \sim 700$ MeV and $L \sim 500$ m.



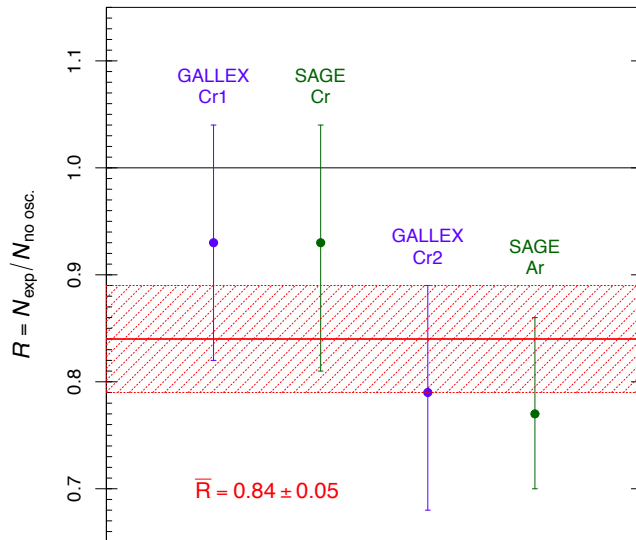
MiniBooNE Coll., PRL 121 (2018)

MiniBooNE reports a **low-E excess** which has increased in significance over time ($3.6\sigma \rightarrow 4.7\sigma \rightarrow 4.8\sigma$) with more data and sophisticated analysis.



Mention et al., 2011. See also Muller et al., PRC 832011; Huber et al., PRC84 2011. And Sinev, I103.2452; Ciuffoli et al., JHEP 12 2012; Zhang et al., PRD87 2013; Ivanov et al., I306.1995.

Reactor anomaly: A recomputation of the reactor fluxes seems to indicate neutrino disappearance, compatible with oscillations into sterile neutrinos with large masses. Recent results have reduced the significance of these results.



Gallium anomaly: The measurement of the fluxes of electron neutrinos in Gallium Radioactive sources reports a value 2.9 sigma away from what expected.

Frekers et al., PLB 706 2011. See also SAGE 2006, 2009; Laveder et al., NPPS 2007, MPLA 2007, PRD 2008, PRC 2011, PRD 2012.

Sterile neutrinos as an explanation?

Sterile neutrinos: hypothetical neutral fermionic singlets of the Standard Model.

Generically they mix with the light neutrinos:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{pmatrix} = U_{4 \times 4} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

Flavour state **Massive state**

Nearly-sterile neutrino, commonly called sterile neutrino

$$\mathcal{L} = \dots + \bar{\ell}_L U_{\ell 4} \gamma_\mu \nu_{4,L} W^\mu + \text{NC} + \text{h.c.}$$

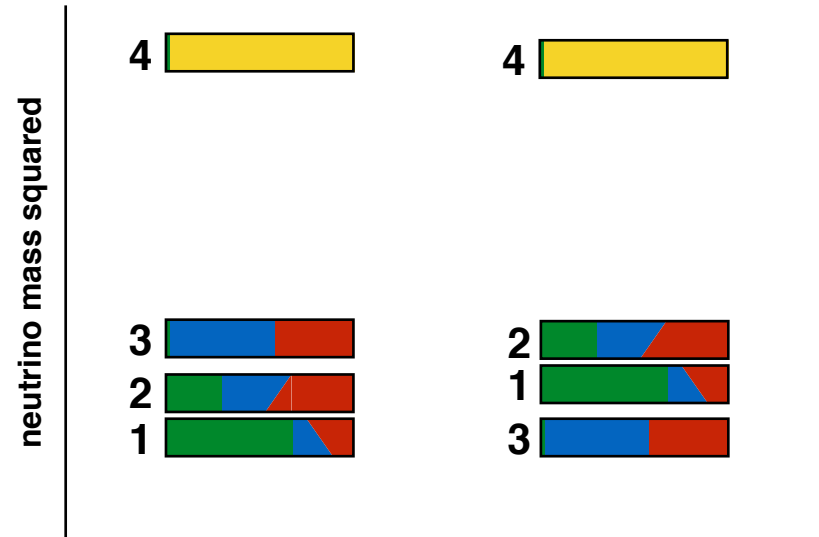
Adding sterile neutrinos to the Standard Model is the simplest possible extension BSM.

- Theory remains anomaly free.
- Can give origin to neutrino masses and explain their smallness (at least in some cases).
- GUT theories embedding L-R symmetries, e.g. $SU(4)$, $SO(10)$,... predict their existence.
- A part from GUT theories, there is no strong motivation for choosing one mass scale instead of another (except for a naturalness principle: setting their mass to zero restores the lepton number symmetry).

Light (nearly-)sterile neutrinos

$$\Delta m_s^2 \ll \Delta m_A^2 \ll \Delta m_{41}^2$$

for 4 massive nus.



Fractional flavour content of massive neutrinos

Appearance oscillation probability at short baselines:

$$P(\nu_\alpha \rightarrow \nu_\beta) = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2 \sin^2 \frac{\Delta m_{41}^2 L}{4E}$$

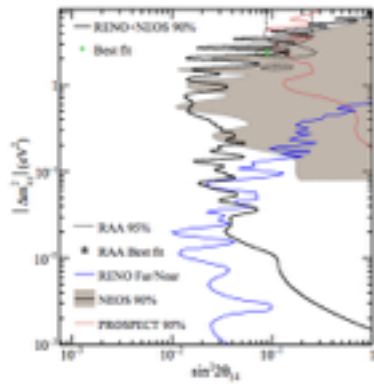
Oscillation **disappearance probability**:

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2) \sin^2 \frac{\Delta m_{41}^2 L}{4E}$$

Light sterile neutrino oscillations can explain the appearance data.

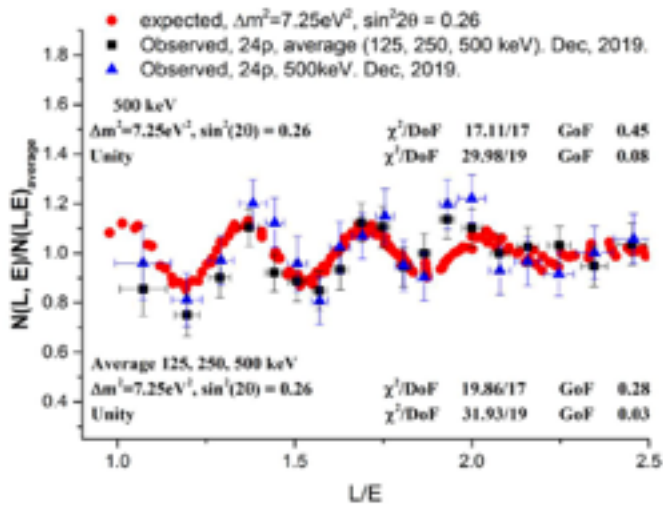
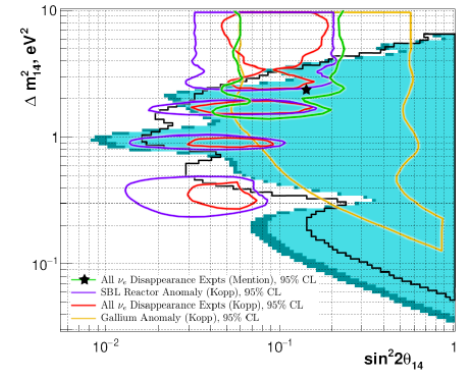
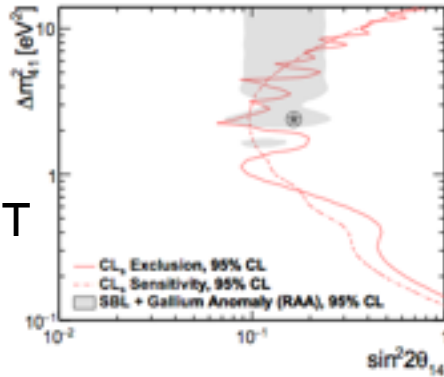
Electron neutrino disappearance

Experiments with very short baselines have been designed to test these anomalies in controlled conditions.



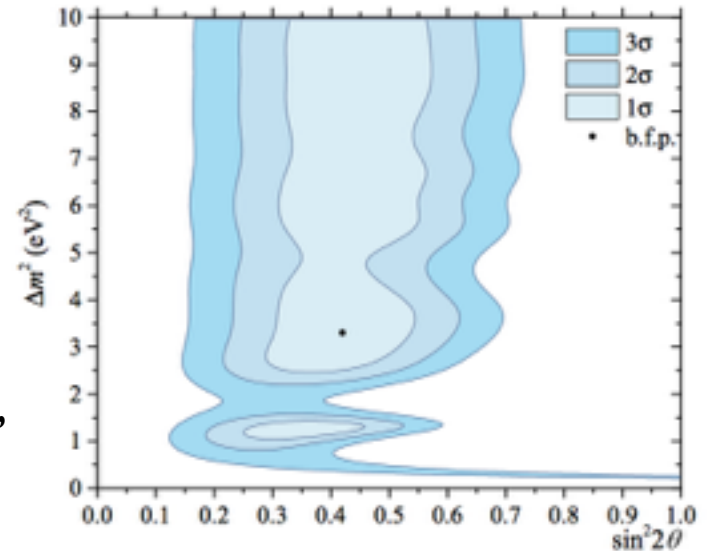
R E N O
+NEOS

PROSPECT

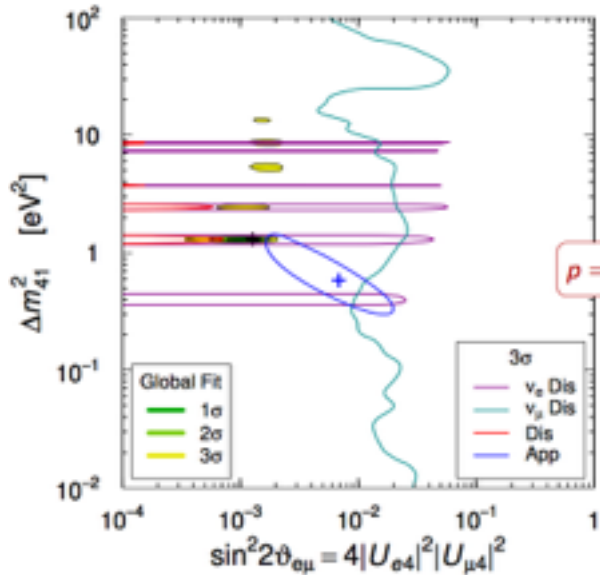


Neutrino4

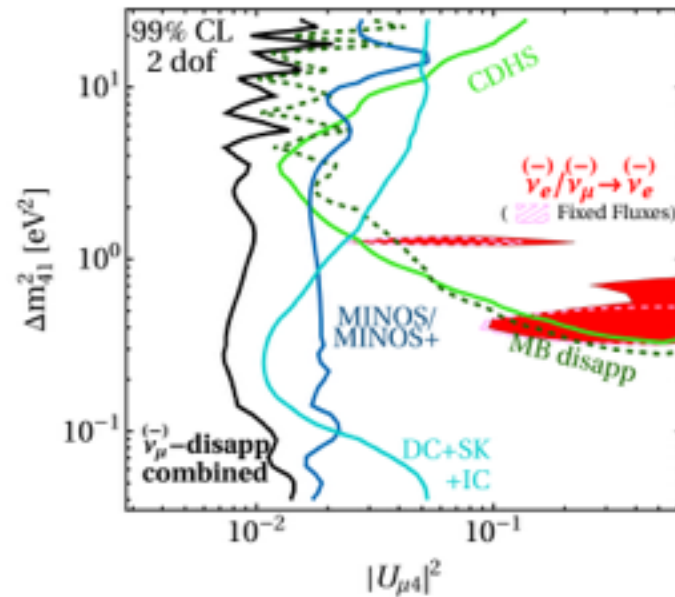
B E S T ,
2109.11482



But there is a **tension** with muon neutrino disappearance data.



Gariazzo et al., in preparation, TAUP 2021



A. Dentler et al., 1803.10661

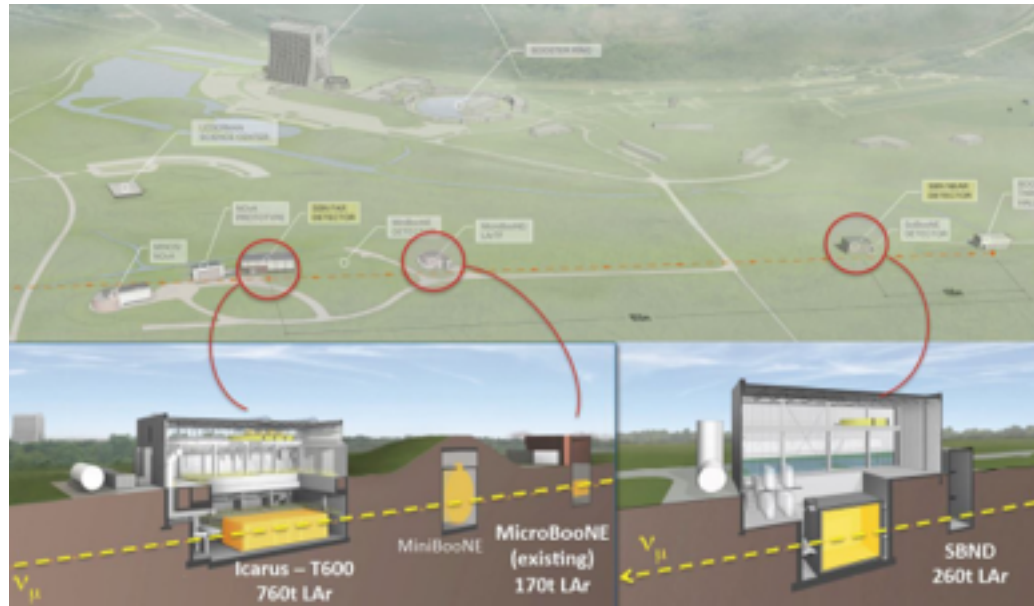
$$P(\nu_\mu \rightarrow \nu_e) = 4|U_{e4}|^2|U_{\mu4}|^2 \sin^2 \frac{\Delta m_{41}^2 L}{4E}$$

$\sin^2 2\theta$

My take: The situation is still rather uncertain. Probably the appearance results are not due to oscillations. Major tension with cosmology.

MicroBooNE and SBN at Fermilab

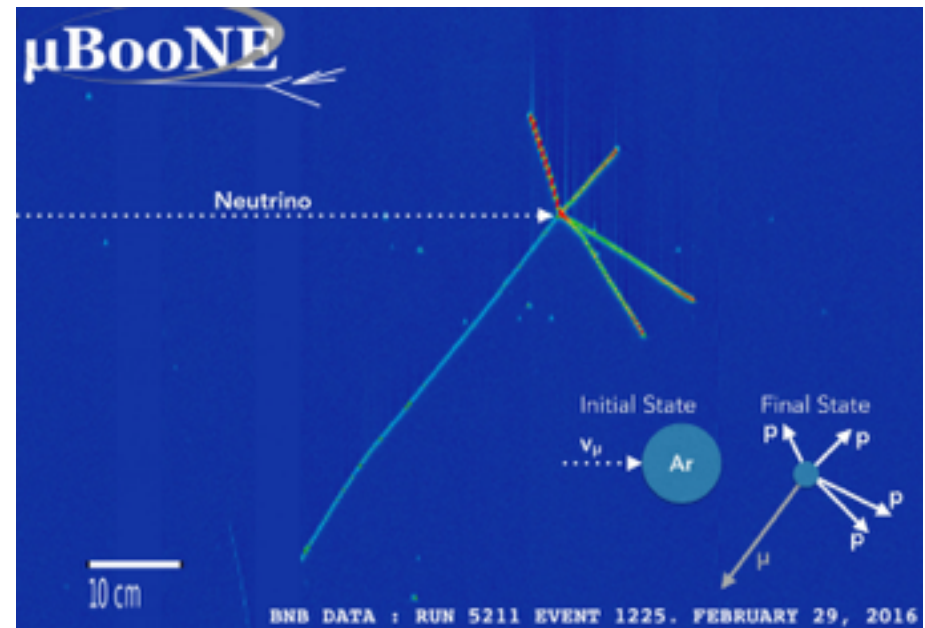
They use accelerator neutrino experiments with $L \sim 100\text{-}600\text{m}$ and $E \sim 700\text{-}800\text{ MeV}$.



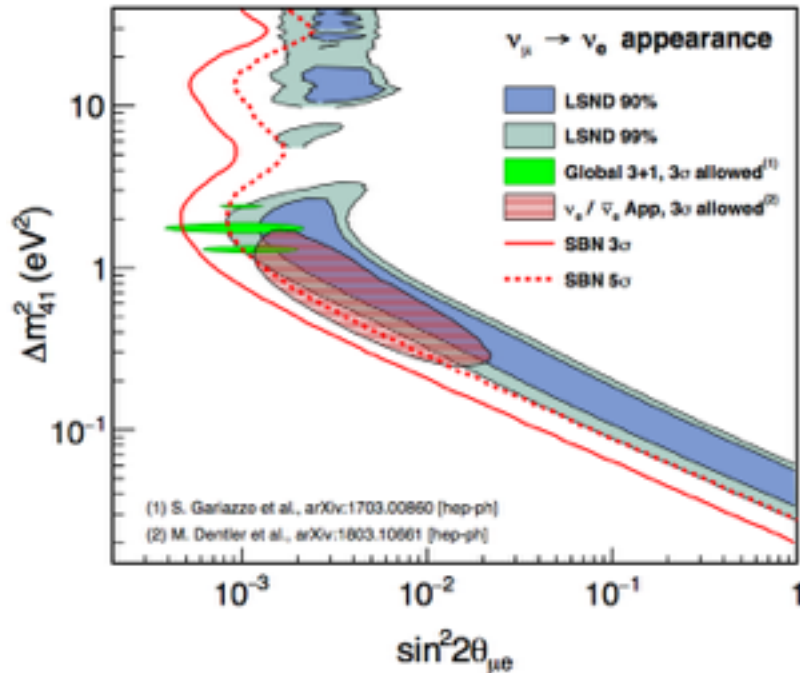
MicroBooNE detector

MicroBooNE event:
muon neutrino
scattering in LAr

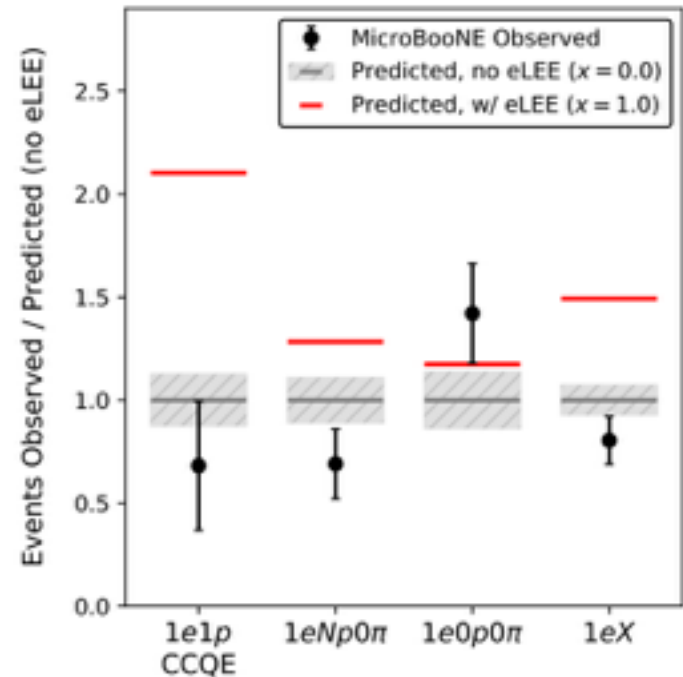
<https://www.bo.infn.it/gruppo2/sbn-it/>



Accelerator neutrino experiments should provide the definitive answer and can check both the appearance and disappearance channels.



MicroBooNE Coll., 1903.04608



MicroBooNE Coll., 2110.14054

MicroBooNE first results **disfavour** a neutrino oscillation explanation, as well as some not understood background (Delta single photons).

BSM explanations for MB LEE

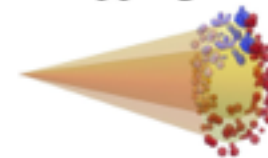
Due to the WC nature of MB, single electrons can be mimicked by photons and by electron-positron pairs (if overlapping or asymmetric).

Electrons? Or Photons?Or Neither?

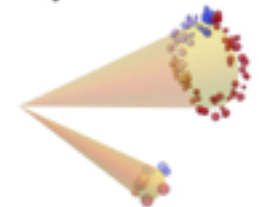
Rich phenomenology developing in recent years around the possibility of the MiniBooNE excess being due to e^+e^- pairs from decays of new exotic particles.

- Decays of **new dark gauge bosons** (Z')
 - E. Bertuzzo, S. Jana, P.A.N. Machado, R.Zukanovich Funchal [Phys.Rev.Lett. 121 24, 241801\(2018\)](#)
 - P. Ballett, S. Pascoli, M. RL [Phys. Rev. D 99, 071701 \(2019\)](#)
 - A. Abdollahi, M. Hostert, S.Pascoli [Phys.Lett.B 820 136531\(2021\)](#)
- General **Extended higgs sectors + Decay**
 - B. Dutta, S. Ghosh, T. Li Phys. [Rev. D 102, 055017 \(2020\)](#)
 - W. Abdallah, R. Gandhi, S. Roy [Phys. Rev. D 104, 055028 \(2021\)](#)
- Decays of **leptophilic axion-like** particles
 - C. V. Chang, C. Chen, S. Ho, S. Tseng [Phys. Rev. D 104, 015030 \(2021\)](#)

Overlapping e^+e^-

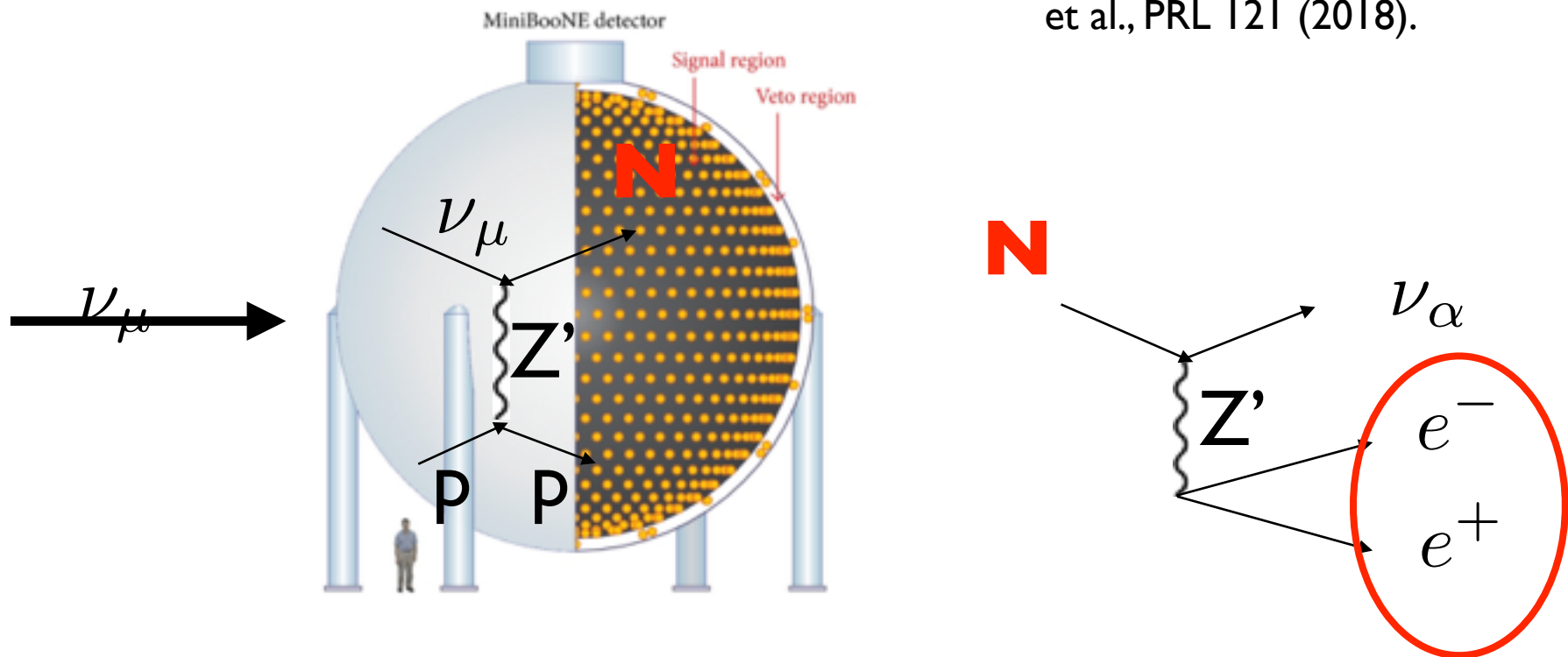


Highly Asymmetric e^+e^-



A viable explanation of the MiniBooNE low-E excess is provided by the **up-scattering of an HNL N** in the detector and **its decay into $ee \nu$** .

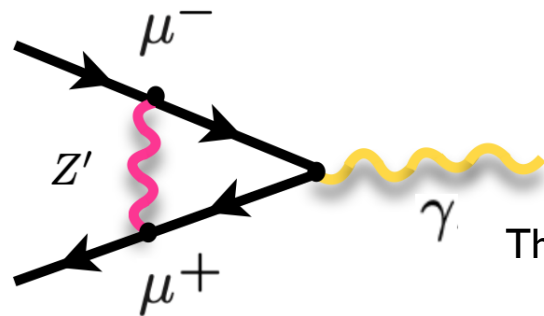
It builds on a decay explanation of MiniBooNE by S. Gineko, PRL 103 (2009). A similar analysis appeared at the same time but with light Z' by E. Bertuzzo et al., PRL 121 (2018).



This type of models require a **low energy extension of the SM** (dark sector) with new gauge or scalar interactions (for the scattering) and new fermions (which need to decay fast).

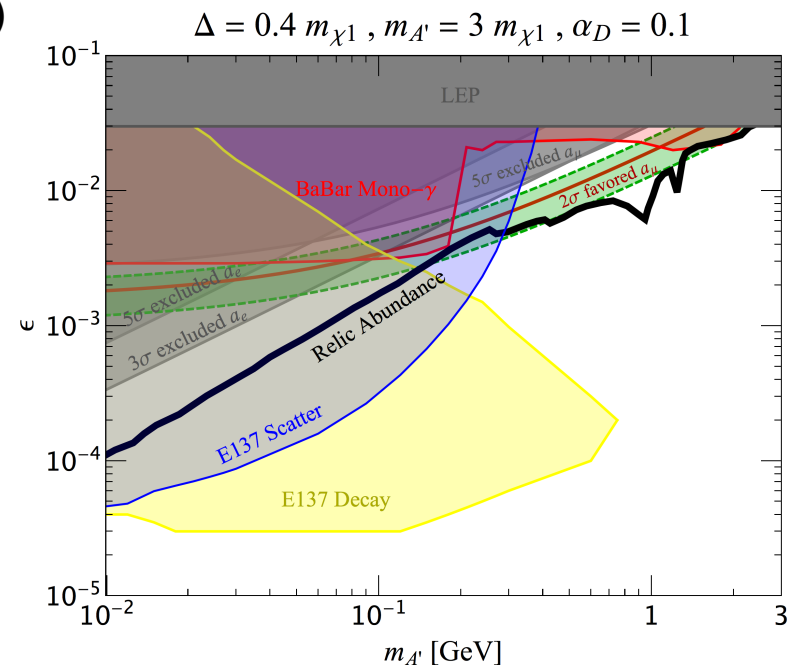
These models have a distinct phenomenology and could also explain other anomalies, specifically the muon ($g-2$).

P. Fayet, PRD75 (2007), M. Pospelov, PRD80 (2009)



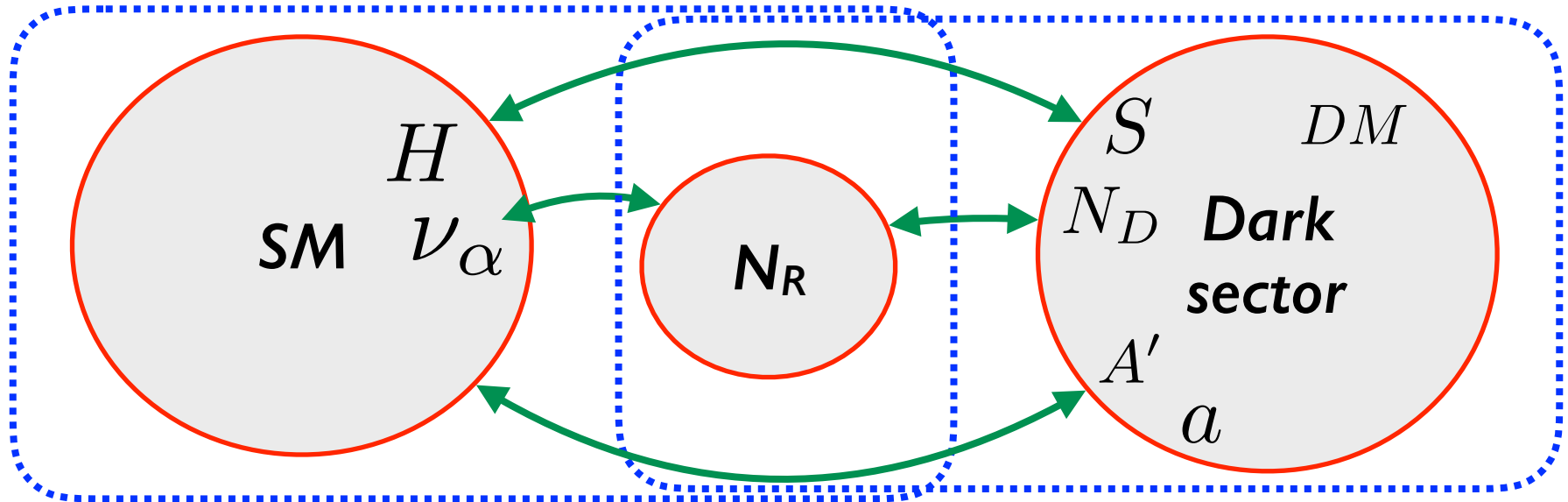
Thanks to A. Abdullahi

as far as Z' decays mainly semi-visibly ($Z' \rightarrow N N$ and N decay fast).



Neutrinos as a window to Dark sectors???

The dark or hidden sector indicate extensions of the SM that are **below the electroweak scale**.



Dark sectors can account for **neutrino masses, the baryon asymmetry, dark matter**.

This would be a major departure from “traditional” BSM thinking and open a very exciting experimental landscape.

Neutrinos as a window to Dark sectors???

The dark or hidden sectors are the extensions of the SM that are relevant at the dark scale.



Dark sectors are not necessarily connected to the baryon

This would be a major departure from “traditional” BSM thinking and open a very exciting experimental landscape.

Conclusions

Neutrinos are the most elusive and mysterious of the known particles. Neutrino masses only particle physics evidence BSM.

Current status: precise knowledge of most of neutrino properties. Key questions open (nature, CPV, neutrino masses) due to be answered in the next decade. Thriving experimental programme.

Surprises in store? MiniBooNE LEE remains a puzzle. New MicroBooNE results point away from sterile neutrinos. Neutrino4 and BEST anomalies?

Are neutrinos pointing towards a new understanding of particles and interactions: **dark sector?**

Neutrinoscope



NeutrinoScope 4+
Bring neutrinos alive with AR!
Cambridge Consultants
★★★★★ 5.0, 7 Ratings
Free

Neutrinoscope is a free App for iPhone and iPad developed by Cambridge Consultants and Durham University. It allows to visualise the neutrinos as they are around us.

