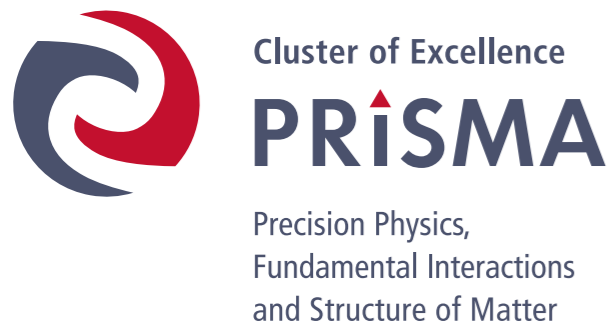


# From the motion of planets to elementary particles

Johannes M. Henn

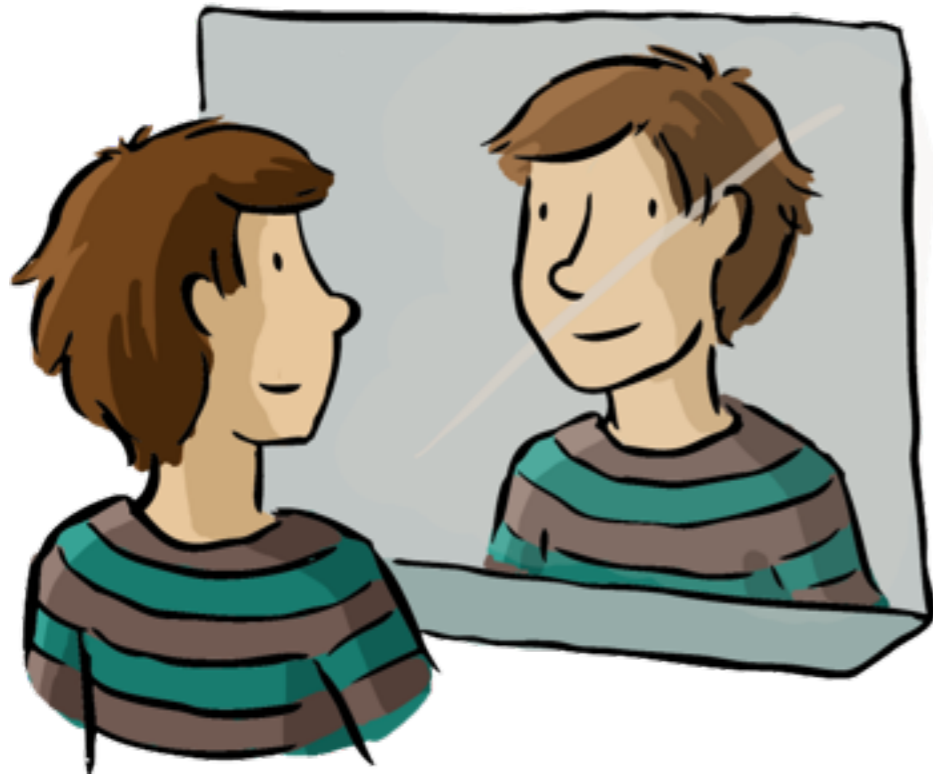
JGU Mainz & Max-Planck Institut für Physik, Munich

Talk at University of Turin  
May 4, 2018



# Outline

Symmetries in  
physics

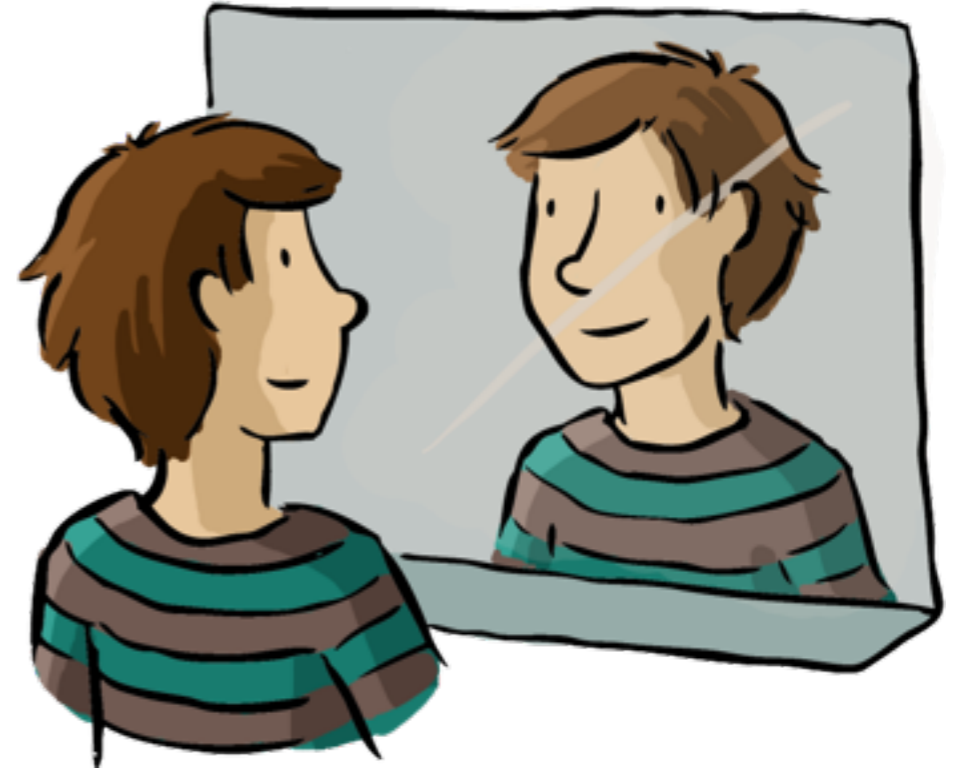


Applications to  
elementary particle  
interactions



picture: Quanta Magazine

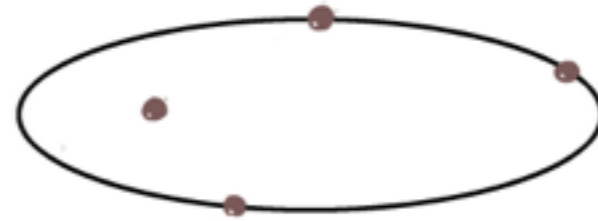
# Symmetries in physics



- guiding principle for finding exact description of Nature
- help to exactly solve idealized models
- obvious versus hidden symmetries

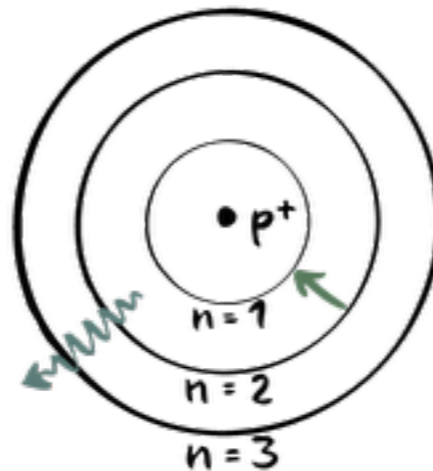
# Symmetry in important physical systems

Kepler problem



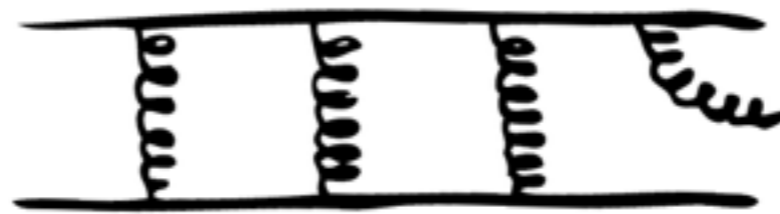
classical  
mechanics

Hydrogen atom



quantum  
mechanics

Interactions of  
elementary particles

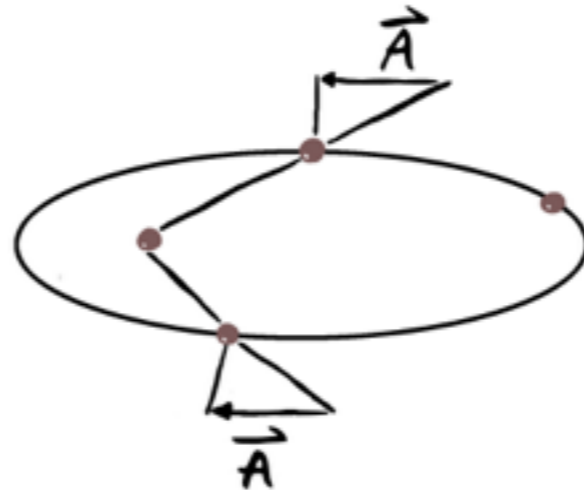


quantum  
field theory

Governed by the same hidden symmetry!

# Regularity of orbits from symmetry

$$V \sim -\frac{\lambda}{r}$$



stable  
orbits

$$V \sim -\frac{\lambda}{r^{0.9}}$$



orbits  
precess

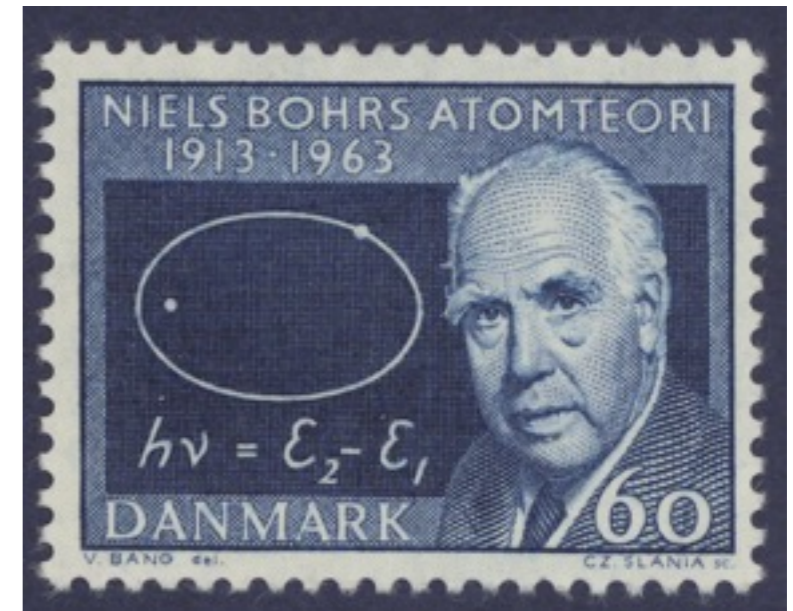
regularity of orbits explained by  
conservation of Laplace-Runge-Lenz vector

$$\vec{A} = \frac{1}{2} \left( \vec{p} \times \vec{L} - \vec{L} \times \vec{p} \right) - \mu \frac{\lambda}{4\pi} \frac{\vec{x}}{|\vec{x}|}$$

# Hydrogen atom

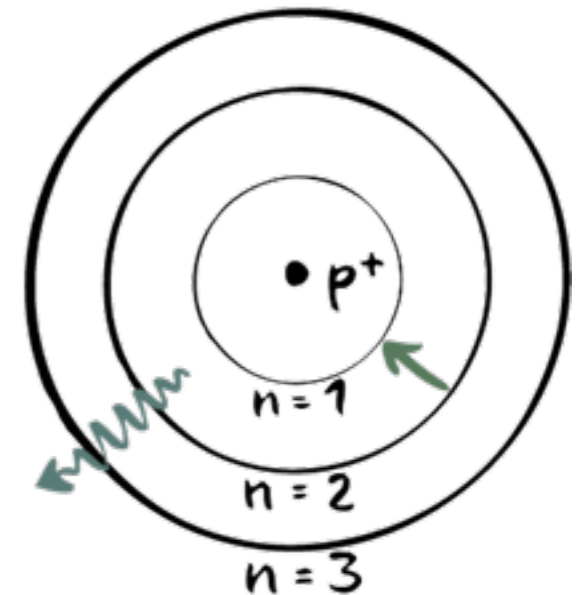
- described by quantum mechanics

- Hamiltonian 
$$H = \frac{1}{2m} p^2 - \frac{k}{r}$$



- spectrum with degeneracy  $n^2$

$$E_n = -\frac{mk^2}{2\hbar^2} \frac{1}{n^2} \quad n = 1, 2, \dots$$



- formula explained by symmetry

# Spectrum determined by symmetry

- Hamiltonian  $H = \frac{1}{2m}p^2 - \frac{k}{r}$

- hidden symmetry:

Laplace-Runge-Lenz-Pauli operator

$$\vec{A} = \frac{1}{2}(\vec{p} \times \vec{L} - \vec{L} \times \vec{p}) - mk \frac{\vec{r}}{r}$$

- conserved quantity in quantum mechanics

$$[H, L_i] = 0 \quad [H, A_i] = 0$$

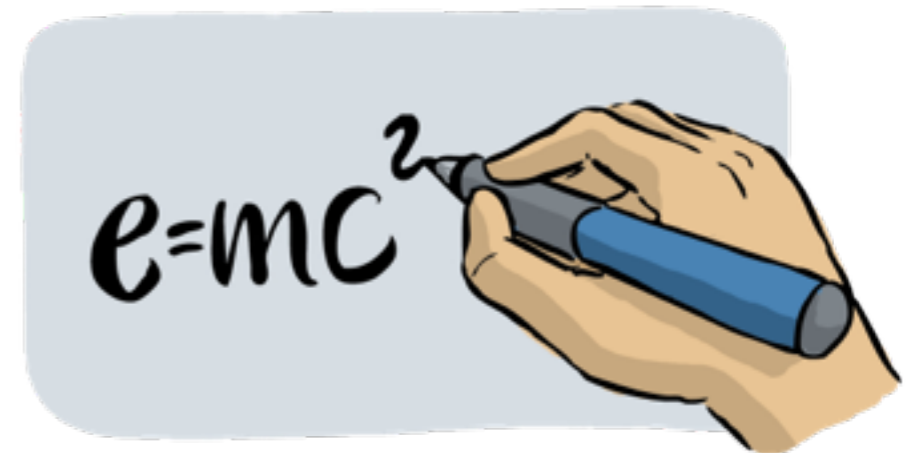
$$[A_i, A_j] = -i\hbar\epsilon_{ijk}L_k \frac{2}{m}H$$

- operator algebra allows to find spectrum



# Hidden symmetry in key physical systems

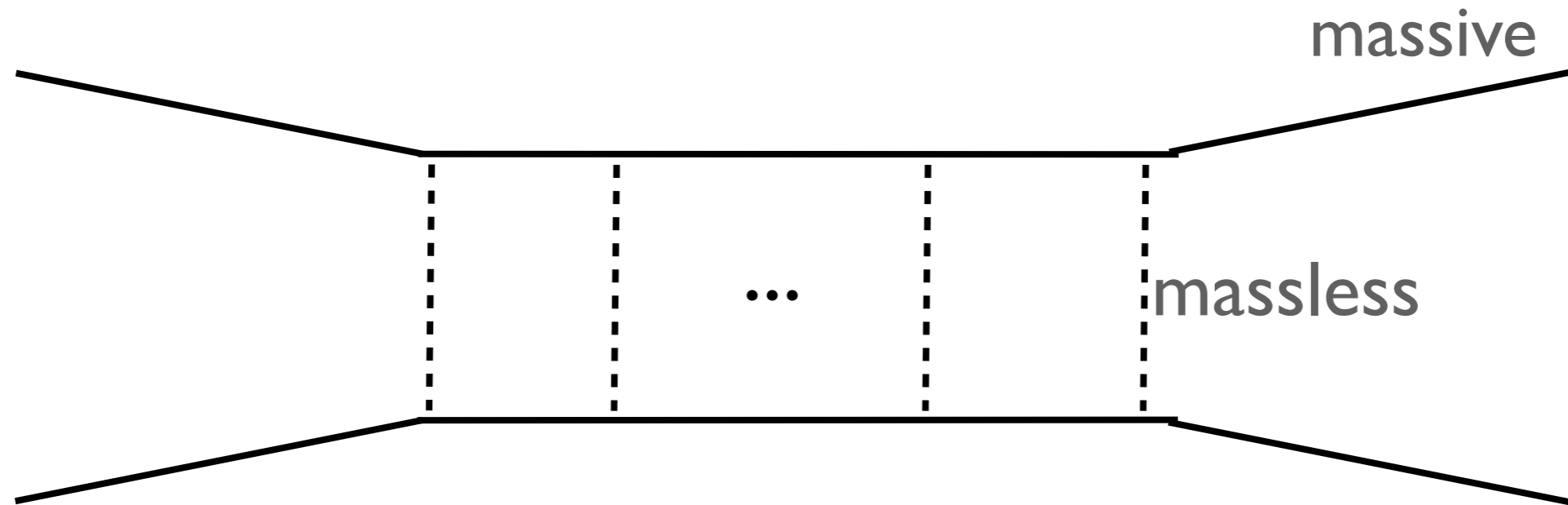
- **Kepler problem** and **hydrogen atom** are important classical and quantum mechanics problems that can be exactly solved
- **have the same hidden Laplace-Runge-Lenz symmetry**
- at higher energies, quantum field theory (QFT) needed
- is there a QFT with the same symmetry?





# towards a relativistic QFT

- Wick-Cutkosky model



- ladder approximation to  $ep \rightarrow ep$ , ignoring spin
- In the non-relativistic limit, this reduces to the hydrogen Hamiltonian

# symmetry of Wick-Cutkosky model

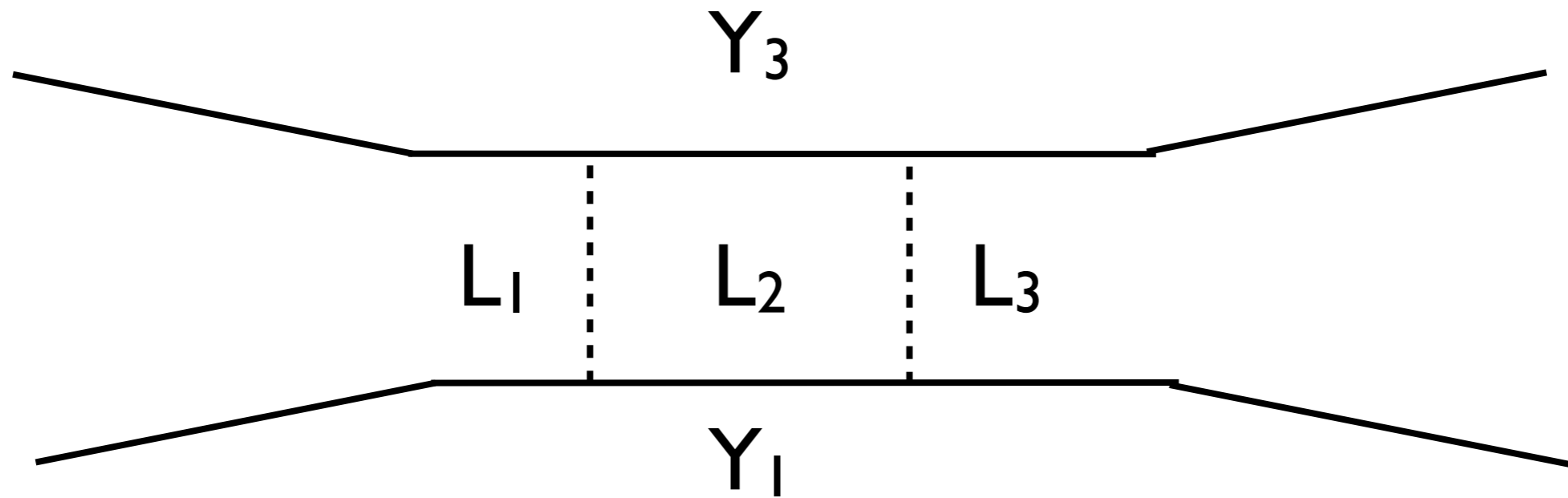
- model possesses an exact  $O(4)$  symmetry, even away from the non-relativistic limit
- consider one rung

$$\dots \int \frac{d^4 \ell_2}{(\ell_2 - \ell_1)^2 [(\ell_2 - p_1)^2 + m^2] [(\ell_2 + p_2)^2 + m^2] (\ell_2 - \ell_3)^2} \dots$$

- symmetry obvious in Dirac's **embedding formalism**

$$L_i^a \equiv \begin{pmatrix} \ell_i^\mu \\ L_i^+ \\ L_i^- \end{pmatrix} = \begin{pmatrix} \ell_i^\mu \\ \ell_i^2 \\ 1 \end{pmatrix} \quad L_i \cdot L_j = (\ell_i - \ell_j)^2 \quad L_i^2 = 0$$

similarly for external momenta



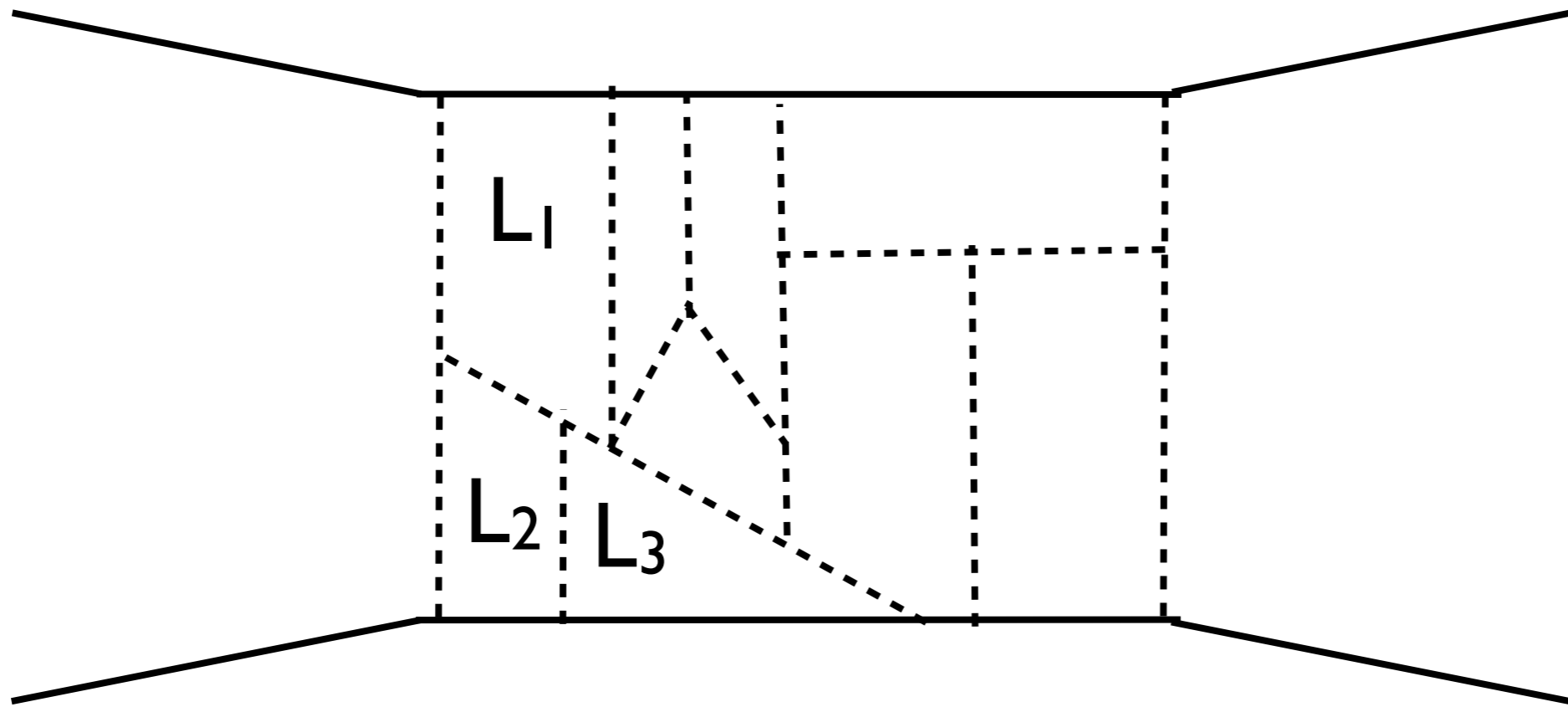
- rung in embedding formalism

$$\dots \int "d^4 L_2" \frac{1}{(L_1 \cdot L_2)(L_2 \cdot Y_1)(L_2 \cdot Y_3)(L_2 \cdot L_3)} \dots$$

- manifest  $SO(6)$  symmetry
- the two vectors  $Y_1, Y_3$  reduce it to  $SO(4)$
- contains the usual  $SO(3)$  as a subgroup
- the remaining 3 generators are the **Runge-Lenz vector!**

# Beyond the ladder approximation

- ladder approximation is arbitrary
- misses multi-particle effects, problems with unitarity
- **Is there a consistent QFT with the LRL symmetry?**
- the simplest way to imagine this requires a planar limit:



- Feynman rules would have to respect the  $SO(6)$  symmetry

# Standard model of elementary particles

THE STANDARD MODEL

	Fermions			Bosons	
Quarks	$u$ up	$c$ charm	$t$ top	$\gamma$ photon	Force carriers
	$d$ down	$s$ strange	$b$ bottom	$Z$ Z boson	
Leptons	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	$W$ W boson	
	$e$ electron	$\mu$ muon	$\tau$ tau	$g$ gluon	
	$H$ Higgs boson*				

\*Yet to be confirmed

Source: AAAS

- Higgs boson: predicted by theorists in the 60's

- as of July 4th, 2012 : discovery by CMS and Atlas experiments

- core part (gluons): non-Abelian gauge theory

$$\mathcal{L} = \frac{1}{4} \text{Tr} \int F_{\mu\nu} F^{\mu\nu}, \quad F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu + ig[A^\mu, A^\nu]$$

$$A^\mu = \sum_{a=1}^{N^2-1} A_a^\mu t_{ij}^a$$

gauge group  $SU(N_c)$ ,  $N_c=3$

- large  $N_c$  limit selects planar Feynman diagrams

# maximally supersymmetric Yang-Mills theory

Particle content similar to QCD:

## QCD

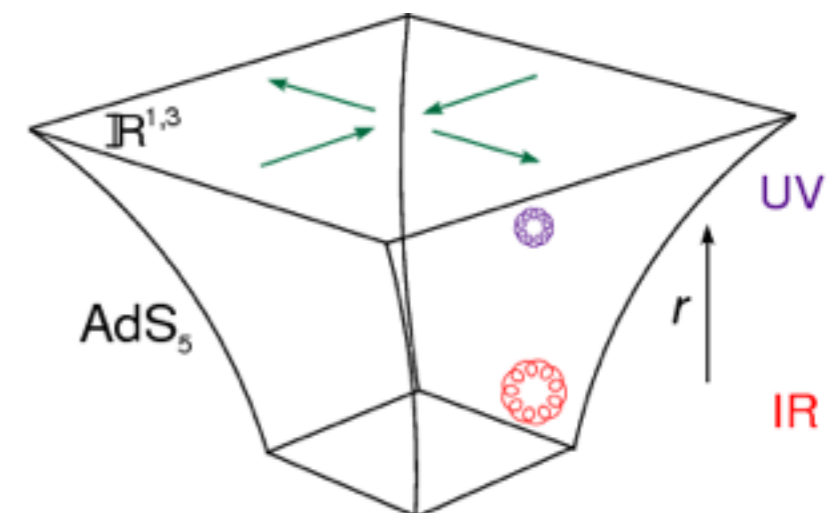
- SU(3) Yang-Mills theory (gluons)
- fermions in fundamental representation

## N=4 supersymmetric Yang-Mills theory

- SU(Nc) Yang-Mills theory
- 4 fermions, adjoint repr.
- 6 scalars

Bonus features:

- supersymmetry; vanishing beta function
- conjectured holographic AdS description



picture from 0803.2475 [hep-th] (L. Dixon)

# Zvi Bern & collaborators studied scattering amplitudes in this theory



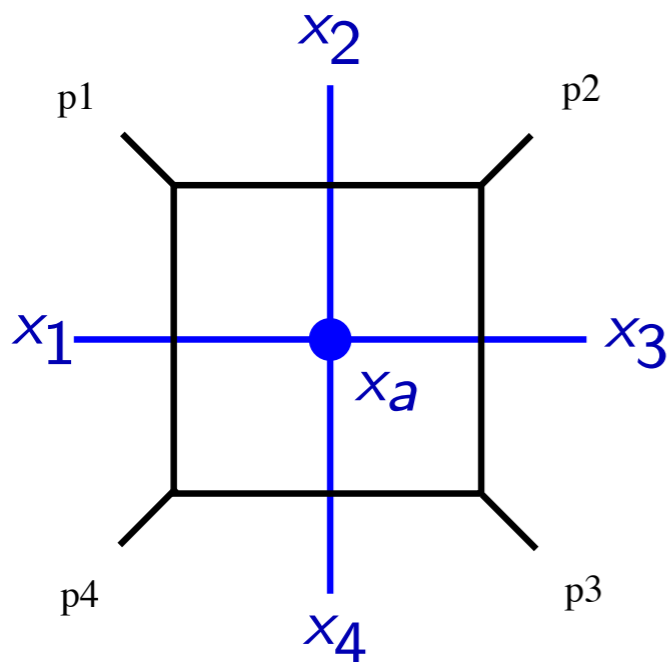
- they used modern (generalized unitarity) methods
- millions of Feynman diagrams sum up to a few 'effective integrals'
- why is this the case?

# Hidden symmetry N=4 sYM

planar N=4 sYM has dual conformal symmetry

[Drummond, JMH, Smirnov, Sokatchev 2006; Alday, Maldacena 2007; Drummond, JMH, Korchemsky, Sokatchev 2007]

e.g. 1-loop four-particle amplitude:



$$p_i^\mu = x_i^\mu - x_{i+1}^\mu$$

$$= x_{13}^2 x_{24}^2 \int \frac{d^D x_a}{x_{1a}^2 x_{2a}^2 x_{3a}^2 x_{4a}^2}$$

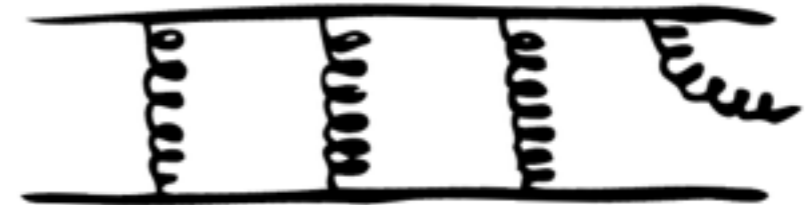
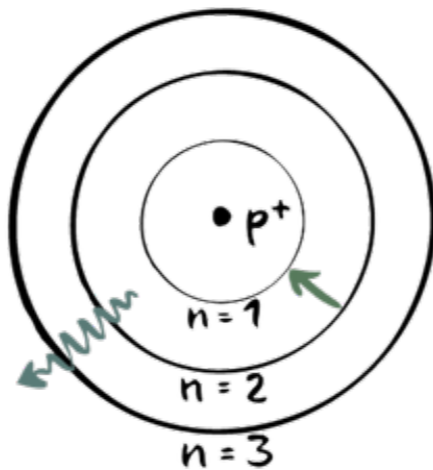
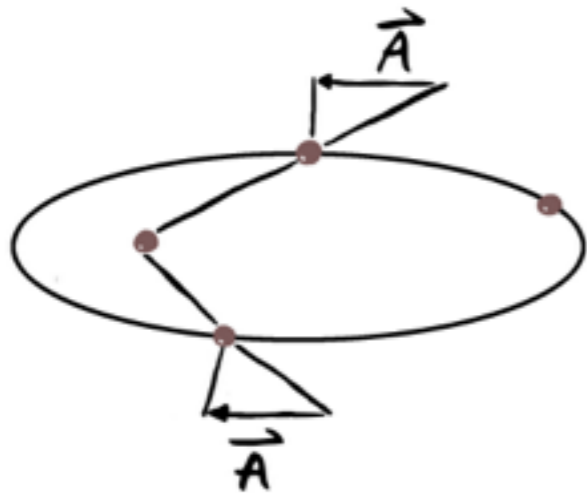
invariant under  $SO(4,2)$  in dual space  $x^\mu \rightarrow x^\mu / x^2$

$$= (Y_1 \cdot Y_3)(Y_2 \cdot Y_4) \int \frac{“d^D L”}{(Y_1 \cdot L)(Y_2 \cdot L)(Y_3 \cdot L)(Y_4 \cdot L)}$$



# summary Laplace-Runge-Lenz symmetry

- LRL symmetry governs several problems



$N=4$  super Yang-Mills theory is the  
'hydrogen atom of the 21st century'

- symmetry explains simplicity
- helpful for finding exact answers

# Applications to elementary particle interactions



picture: Quanta Magazine

# Multi-particle collisions as the next frontier



picture: Quanta Magazine

- at high energies, many particles produced
- challenge to evaluate the virtual corrections
- long experimenter's wishlist for theorists, e.g.

$$pp \rightarrow 3 \text{ jets} \quad pp \rightarrow H + 2 \text{ jets} \quad pp \rightarrow V + 2 \text{ jets}$$

- challenge: 5-particle processes at 2 loops

# 'Ideal' and 'real' scattering amplitudes

Is there some simpler version of QCD that allows to understand key properties of scattering amplitudes?



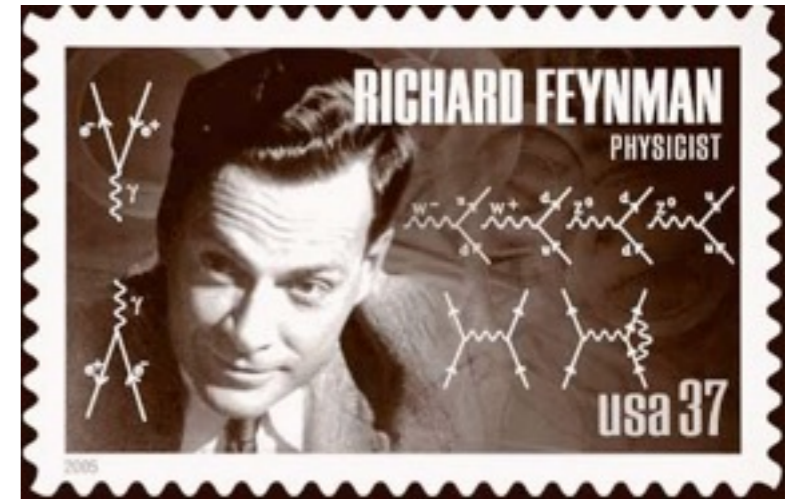
We need to obtain numerical results for cross sections at the LHC.

This talk: **tools for 'real' QCD coming from 'ideal' amplitudes**

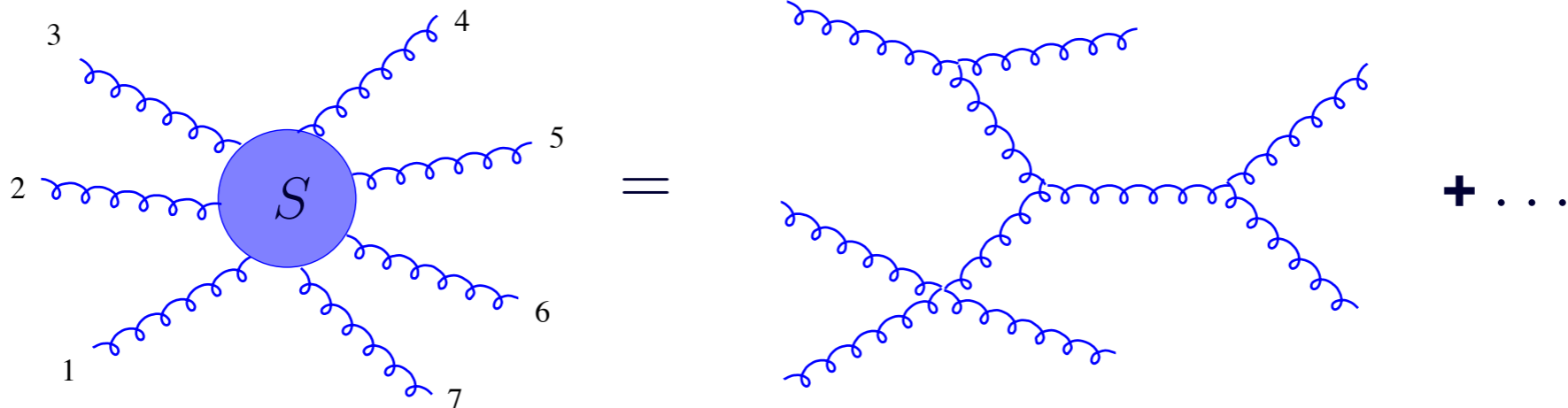
# Scattering amplitudes

Computational recipe:

- (1) draw all Feynman diagrams
- (2) compute them!



Often difficult in practice! E.g. tree-level gluon scattering:



number of external gluons	4	5	6	7	8	9	10
number of diagrams	4	25	220	2485	34300	559405	10525900

Final results much simpler than intermediate steps! **Why?**

# Simplicity of amplitudes from symmetry

Tree-level gluon amplitudes are 'secretly' supersymmetric!

They have the full symmetry of N=4 sYM

- conformal supersymmetry  $J^a = \sum_{i=1}^n J_i^a$
- hidden dual conformal symmetry

combine to  
Yangian symmetry

$$J^a = \sum_{i=1}^n J_i^a \quad Q^c = f^{abc} \sum_{i < j}^n J_i^a J_j^b$$

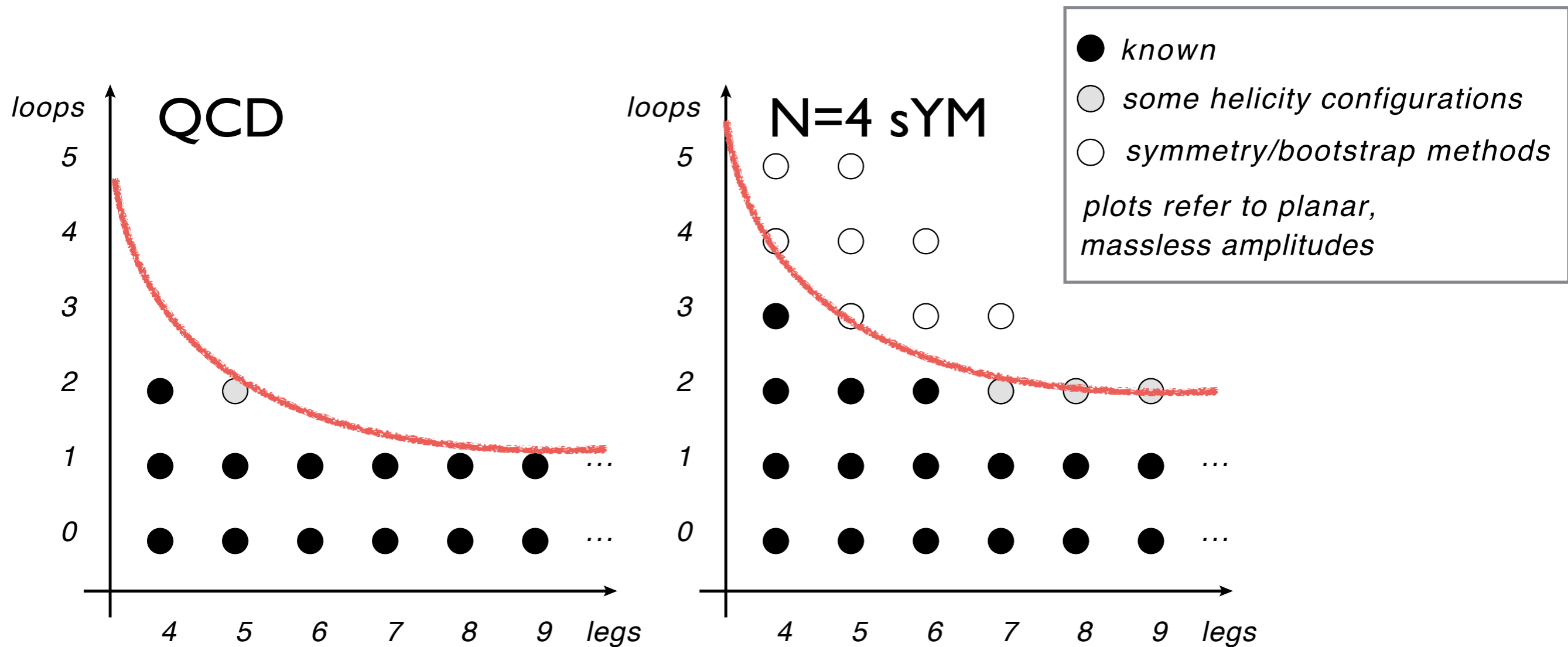
$$J^a \mathcal{A}_n = 0$$

$$Q^a \mathcal{A}_n = 0$$

explains simplicity!

symmetry & collinear behavior fixes tree-level amplitudes!

# State of the art loop amplitudes



- **frontier of knowledge** pushed forward continuously
- N=4 sYM a good prediction what we can hope to achieve next in QCD

# Bootstrapping scattering amplitudes



Can we fix amplitudes from general properties?

- symmetries
- analytic properties
- physical limits



# Bootstrap (pre)history

- 1960's: determine S-matrix from analytic properties



- 1994: 'One loop n point gauge theory amplitudes, unitarity and collinear limits'

[Bern, Dixon, Dunbar, Kosower]

- 2011: bootstrap in planar maximally supersymmetric Yang-Mills theory

[Dixon, Drummond, JMH]

many further developments [Almelid, Bartels, Bargheer, Caron-Huot, Del Duca, Dixon, Druc, Drummond, Duhr, Dulat, Gardi, Harrington, JMH, von Hippel, Marzucca, McLeod, Paulos, Pennington, Parker, Papathanasiou, Scherlis, Schomerus, Sprenger, Spradlin, Trnka, Verbeek, Volovich]

- 2017: first application to multi-loop QCD integrals, non-planar

[Chicherin, JMH, Mitev]

# Bootstrap approach

$\vec{x}$  kinematic dependence

$D = 4 - 2\epsilon$  dimension

$$\mathcal{A}(\vec{x}, \epsilon) = \sum_{i,j,k} c_{ijk} \frac{1}{\epsilon^i} r_j(\vec{x}) f_k(\vec{x}) + \mathcal{O}(\epsilon)$$

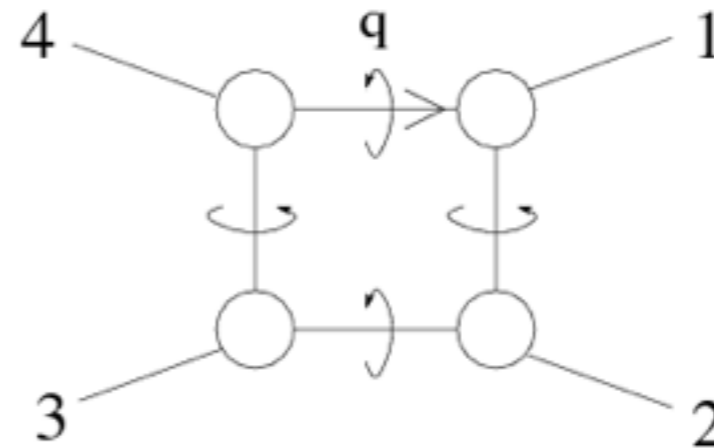
- Laurent expansion in  $\epsilon$
- rational/algebraic normalization factors
- special functions
- unknowns: *finite number* of coefficients

# Constraints on rational factors

- controlled by leading singularities

[Cachazo '08; Arkani-Hamed, Bourjaily, Cachazo, Trnka '10]

- idea: information contained in loop integrand
- perform integral over closed cycles

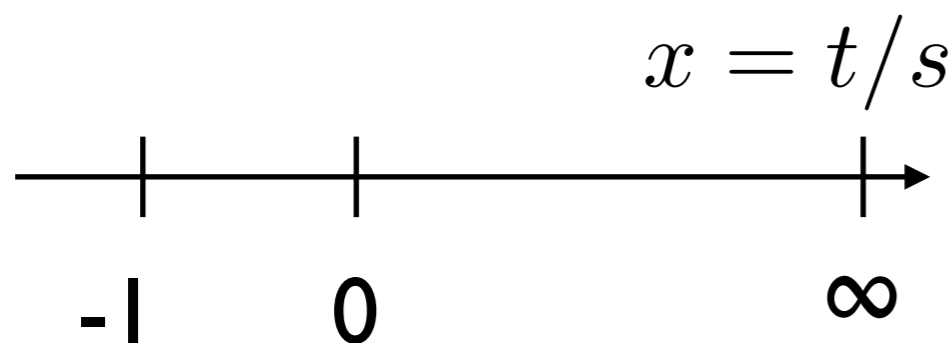


- residue computation much simpler compared to space-time integration

# Finding the space of **special functions**

- improved understanding of iterated integrals
- ‘symbol’ technology [Goncharov, Spradlin, Vergu, Volovich, 2010]
- canonical differential equations defining special functions [JM, 2013]
- singularities of functions from Landau equations

Example: massless 4-particle scattering

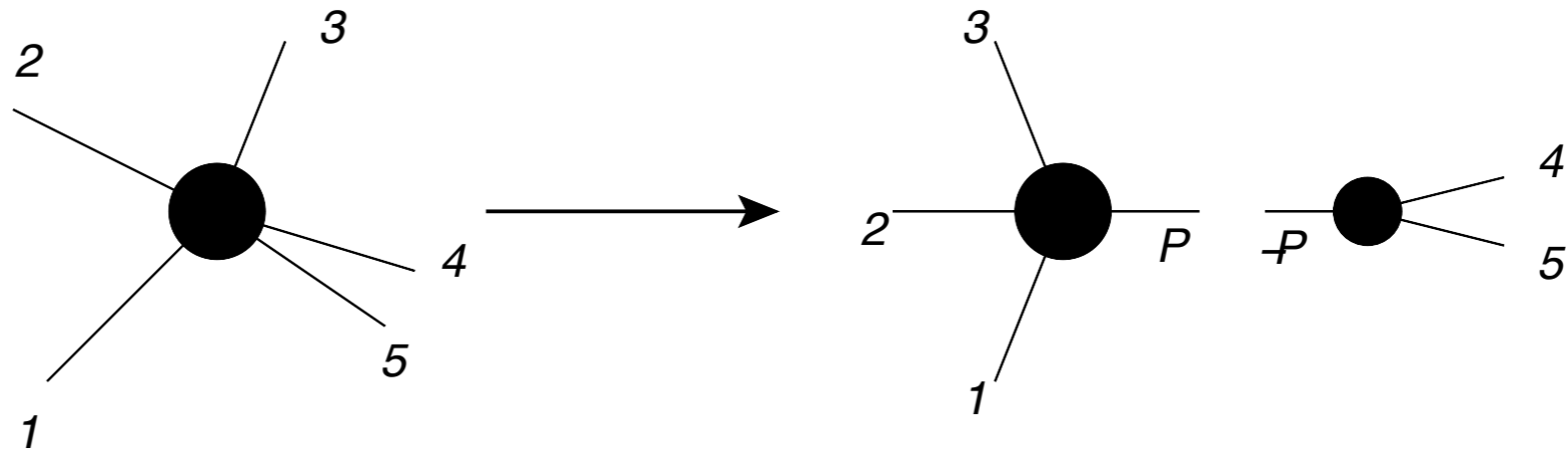


singular points correspond to

$$s = 0, \quad t = 0, \quad u = -s - t = 0$$

# Constraints from symmetries and physical properties

- impose all known symmetries on ansatz
- universal behavior in (singular) limits
- soft, collinear limits



- high-energy, Regge limit
- constraints on discontinuities (e.g. Steinmann relations)

# Sample applications

- Six-particle amplitudes in  $N=4$  sYM known to high loop orders

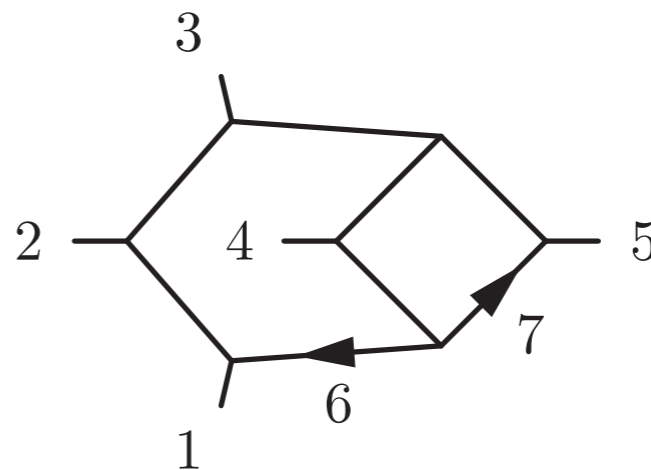
[Caron-Huot, Dixon, McLeod, von Hippel, 2016]

- Applications to quantities in effective field theory

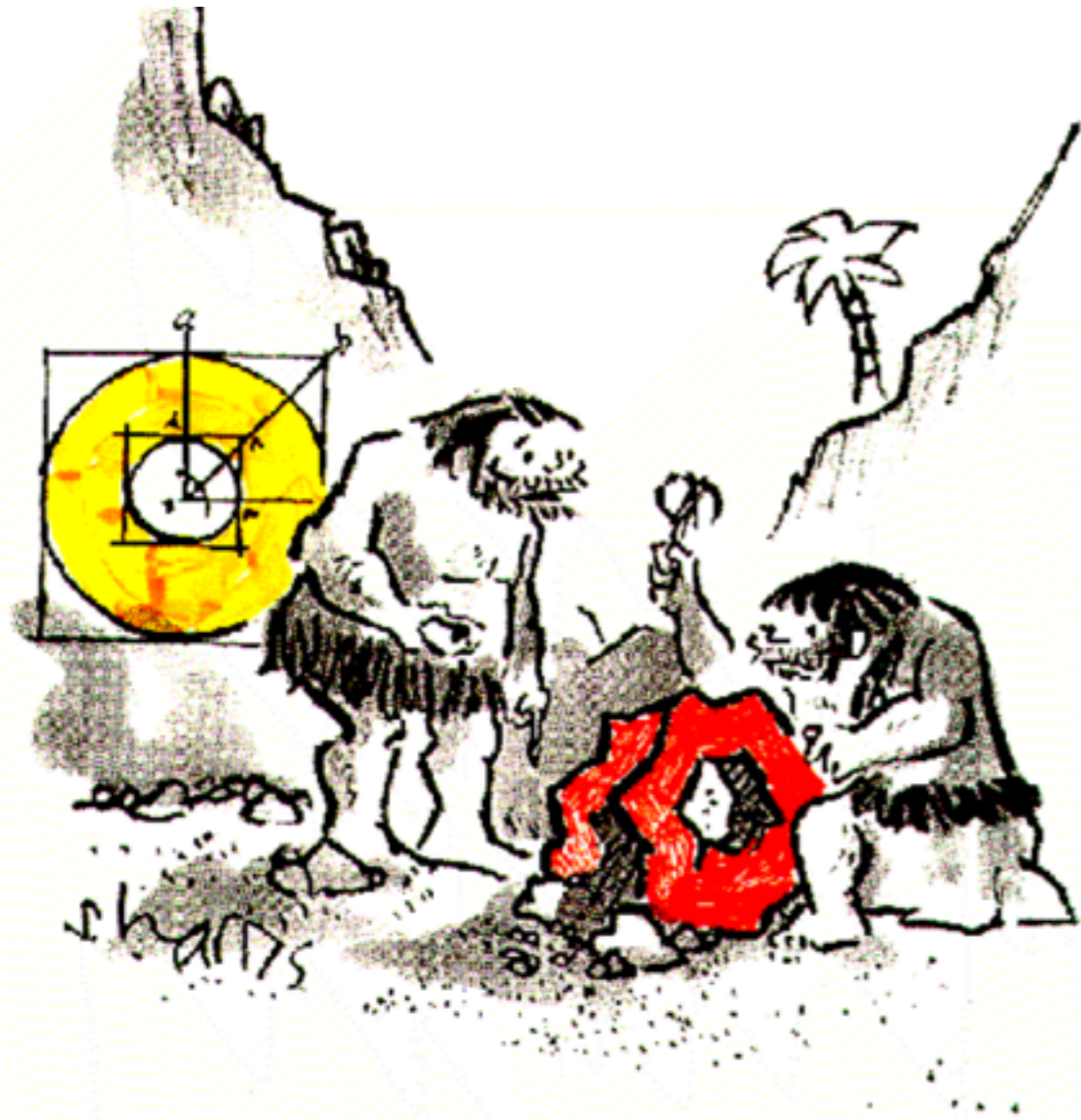
[Li, Zhu, 2017]

- First application to non-planar five-particle integrals in QCD

[Chicherin, JMH, Mitev, 2017]



# Comment on novel methods



*"I guess there'll always be a gap between science and technology."*

- N=4 sYM exciting laboratory for developing ideas
- with refinements, applications to QCD possible
- e.g. open door to 2-loop QCD amplitudes, needed for LHC physics

# Conclusion

- the same hidden symmetry governs several important problems:

- motion of planets
- hydrogen atom
- elementary particle interactions



- amplitude bootstrap



scattering amplitudes determined from symmetries, analytic properties, and physical limits



# Thank you!



(most) illustrations by Joy Katzmarzik,  
[www.leap4joy.de](http://www.leap4joy.de)

# Fruitful interplay of research fields

phenomenology of elementary particles

fundamental aspects of quantum field theory

scattering amplitudes

mathematics  
number theory  
algebraic geometry

string theory, AdS/CFT  
integrable systems



ETH zürich



6-10 july  
**amplitudes 2015**  
INTERNATIONAL CONFERENCE  
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