

The search for neutrinoless double beta decay

Michel Sorel



Colloquium

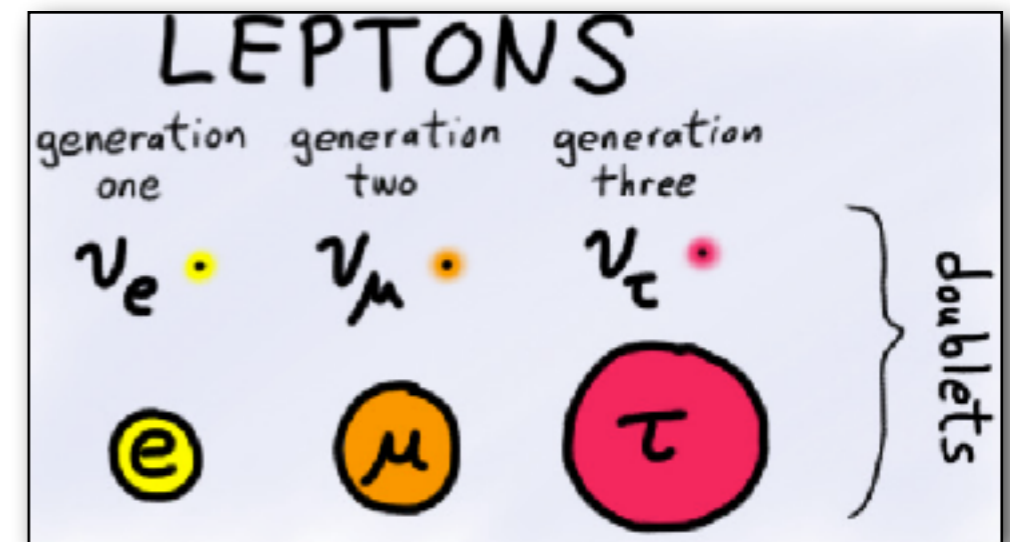
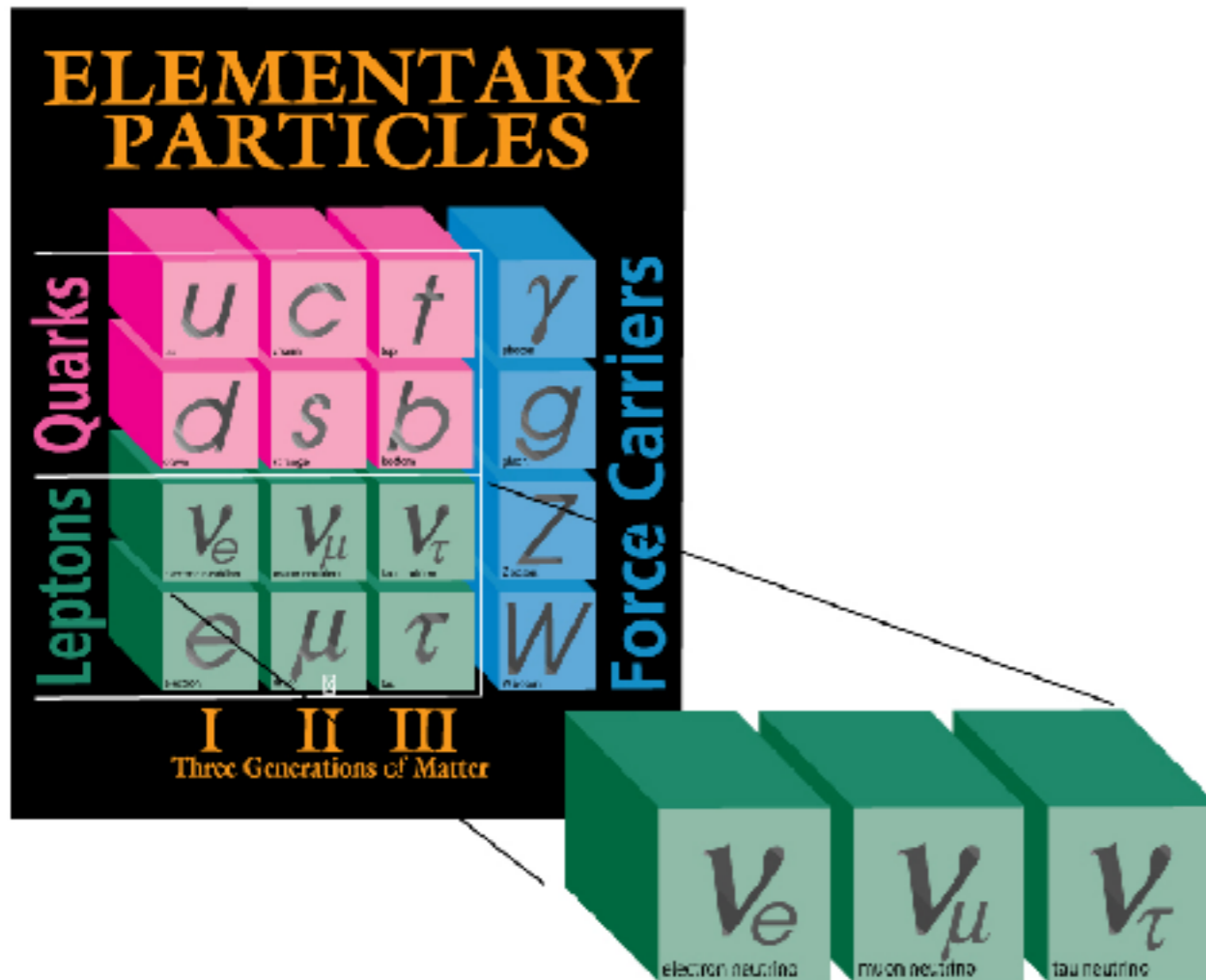
Dip. Fisica Università degli Studi di Torino and INFN Sez. di Torino

May 19th, 2017

Chapter 1:

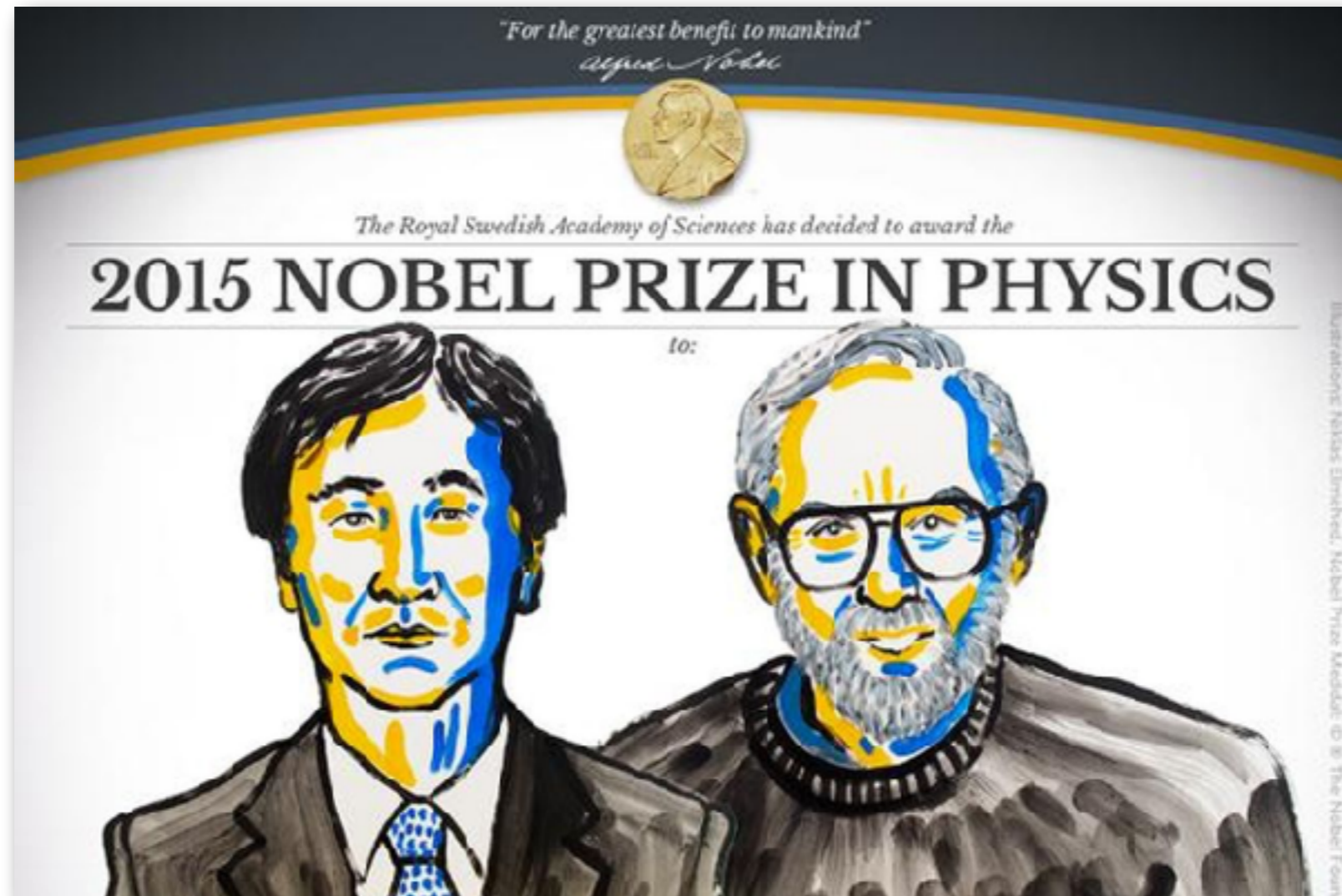
Neutrinos and the flavour puzzle

Neutrinos in the Standard Model



- **Standard Model neutrinos:** 3 types, only weakly interacting, massless

Late '90s: cracks in the Standard Model

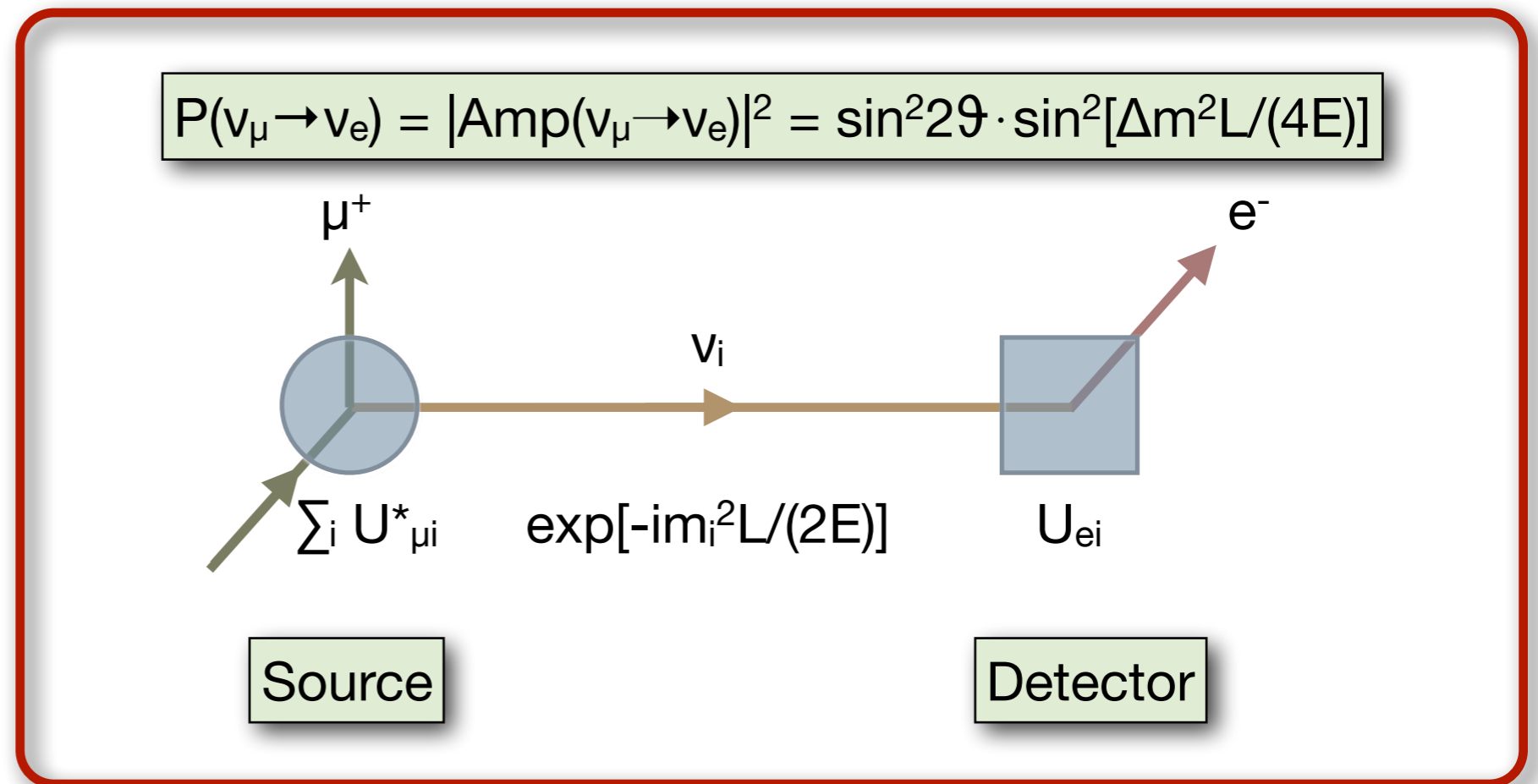
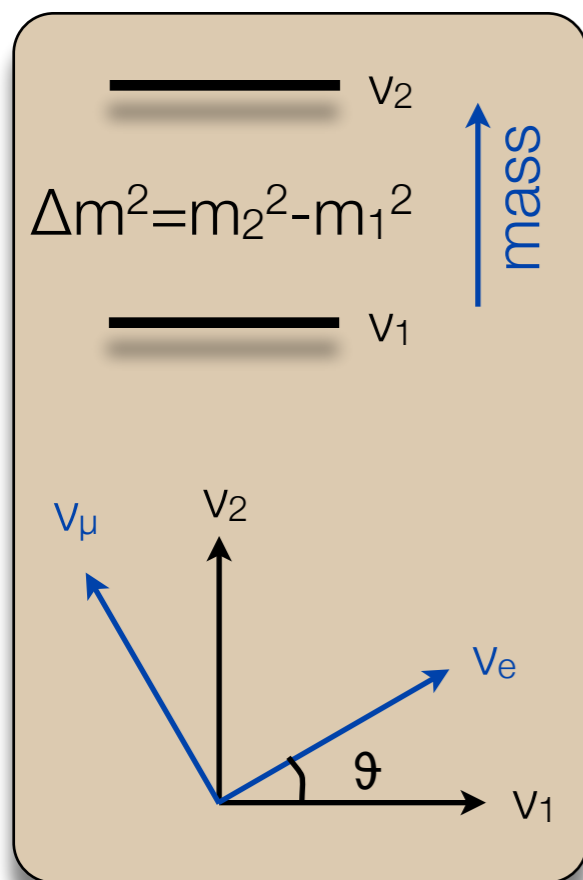


Nobel Prize in Physics 2015:

to Takaaki Kajita (Super-Kamiokande) and Arthur B. McDonald (SNO)
“for the discovery of neutrino oscillations, which shows that neutrinos have mass”

Neutrino oscillations

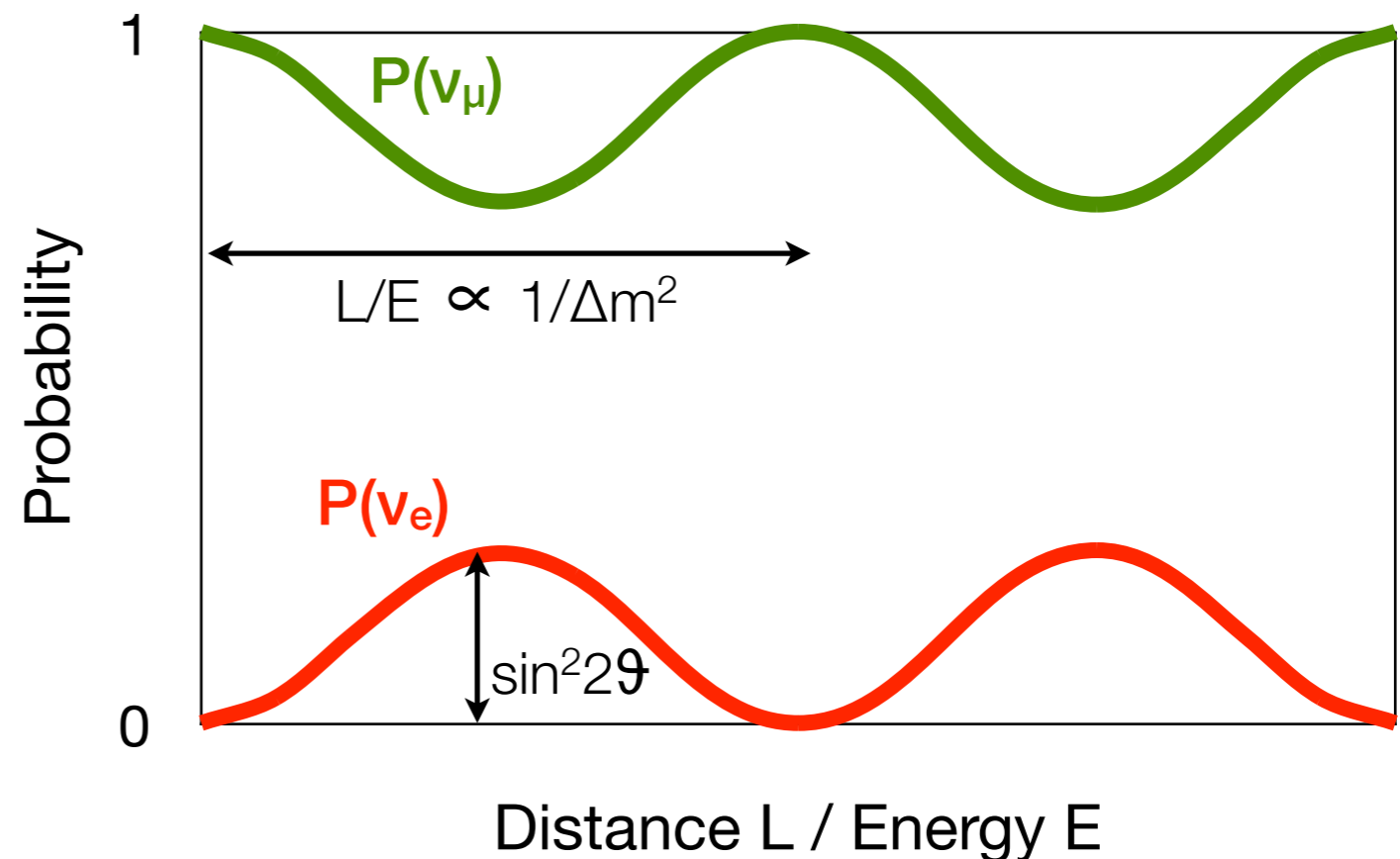
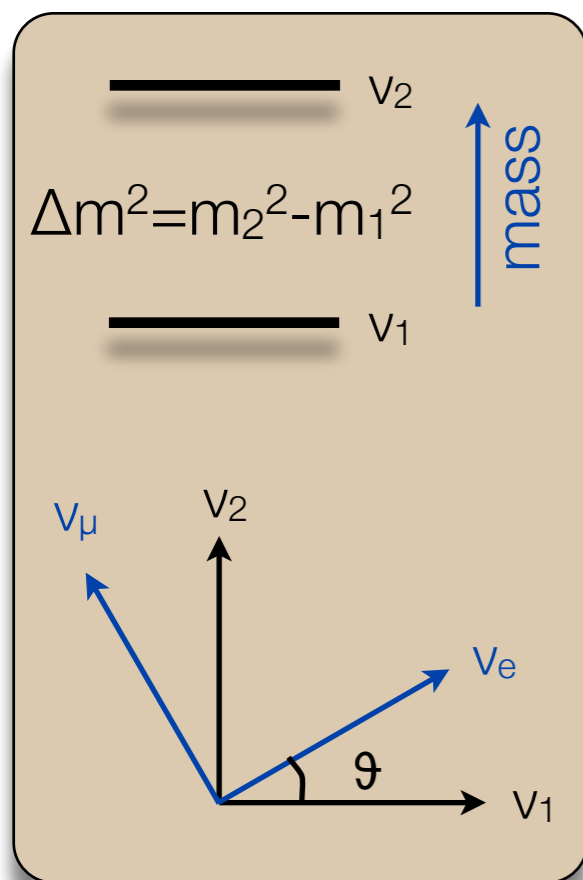
- Neutrinos change flavour as they propagate following oscillatory pattern
- Neutrino oscillation implies massive neutrinos and neutrino mixing



2-neutrino mixing example, for ν_μ beam with energy E

Neutrino oscillations

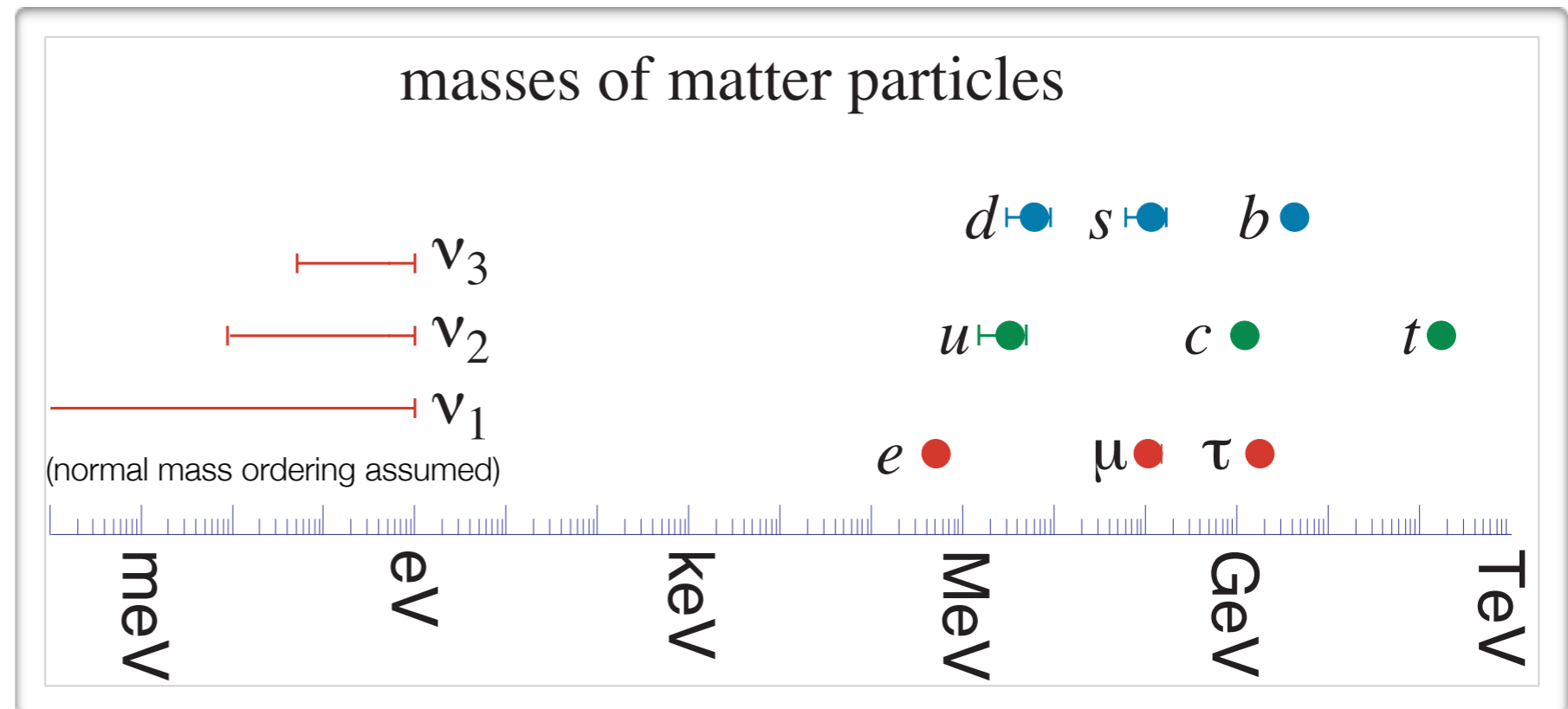
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2-neutrino mixing example, for ν_μ beam with energy E

Neutrinos and the flavour puzzle

- **Neutrino mass:** non-zero and small
- **3 ν flavour mixing:** large



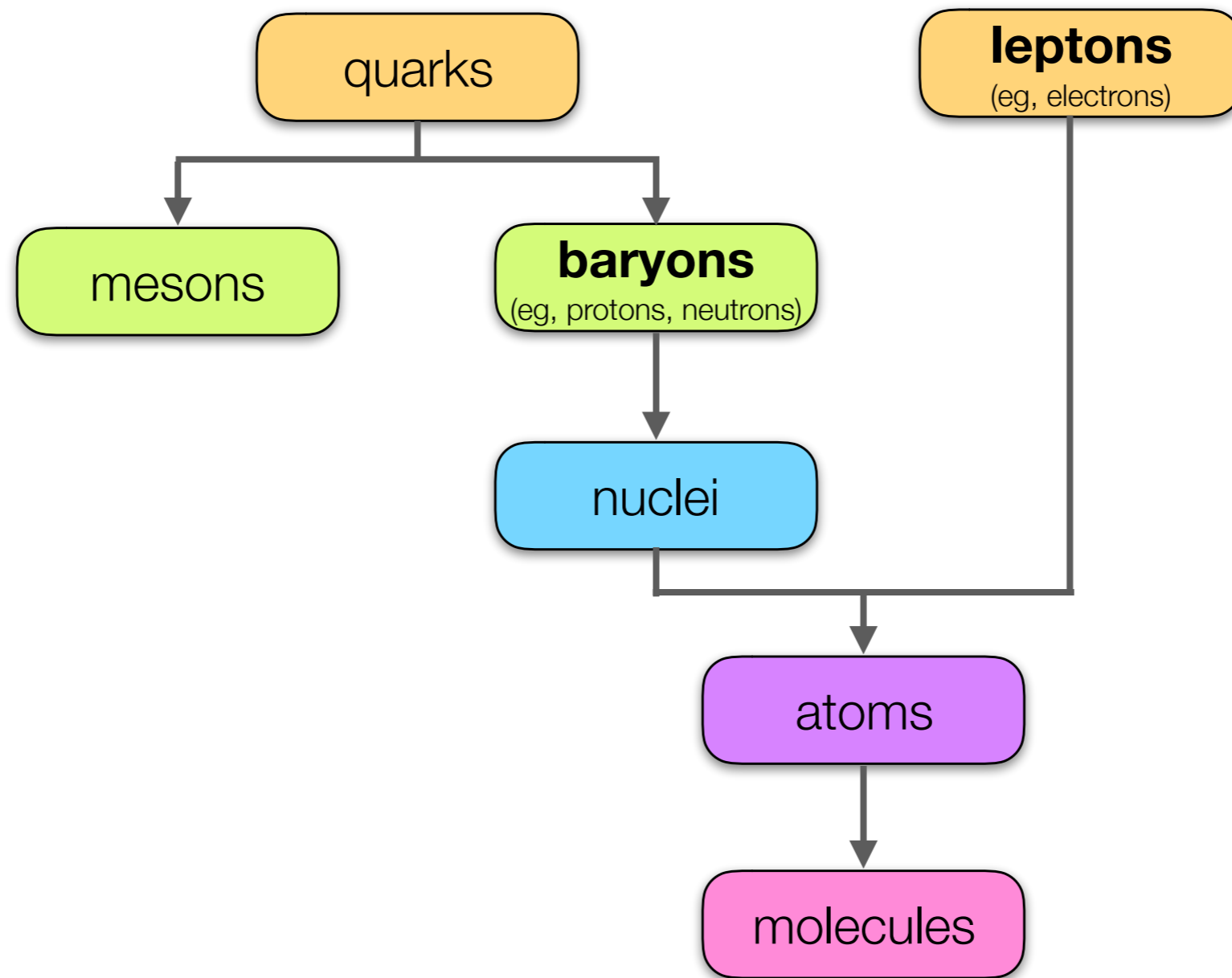
- **Flavour puzzle:**

- Why three generations?
- Why hierarchical masses?
- Origin of mixing pattern?

Chapter 2:

*T*he baryogenesis puzzle

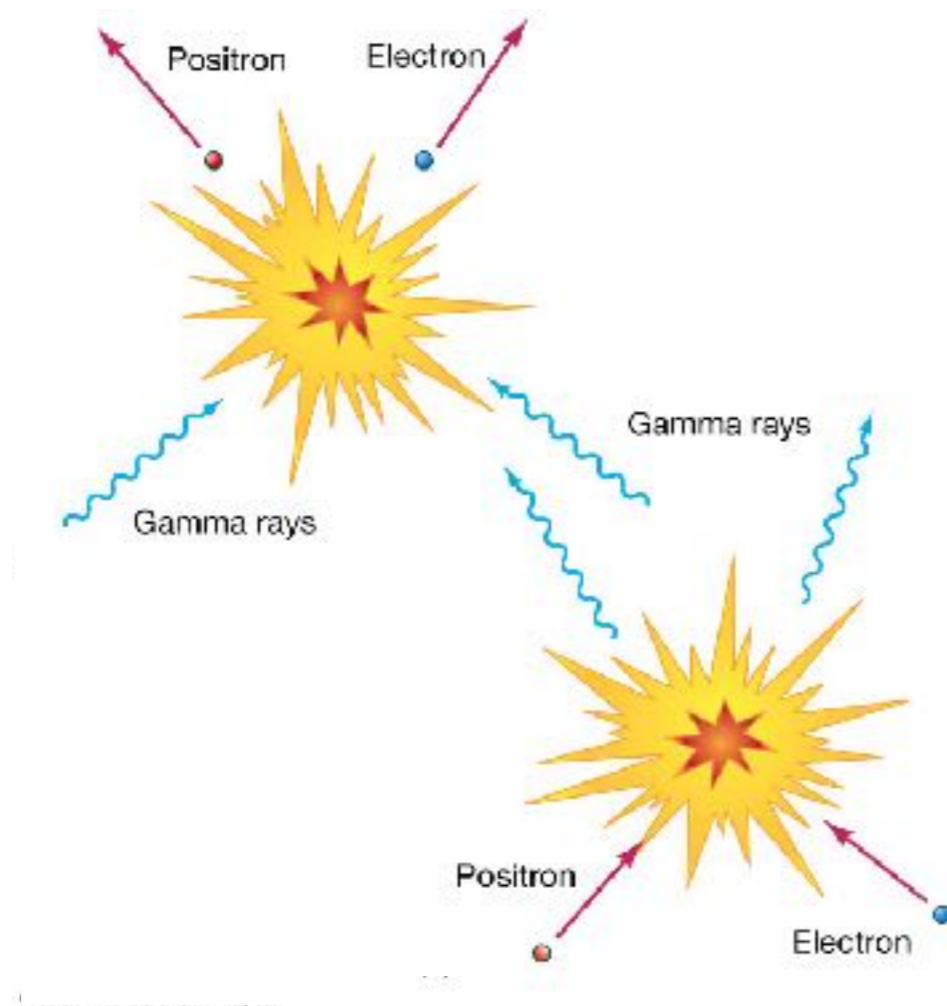
The building blocks of matter



- Ordinary matter requires both baryons and leptons

The first matter in the Universe

- Early Universe dominated by high-energy radiation
- Photons with enough energy to produce particle-antiparticle pairs
- Matter and antimatter thought to be equally abundant at first



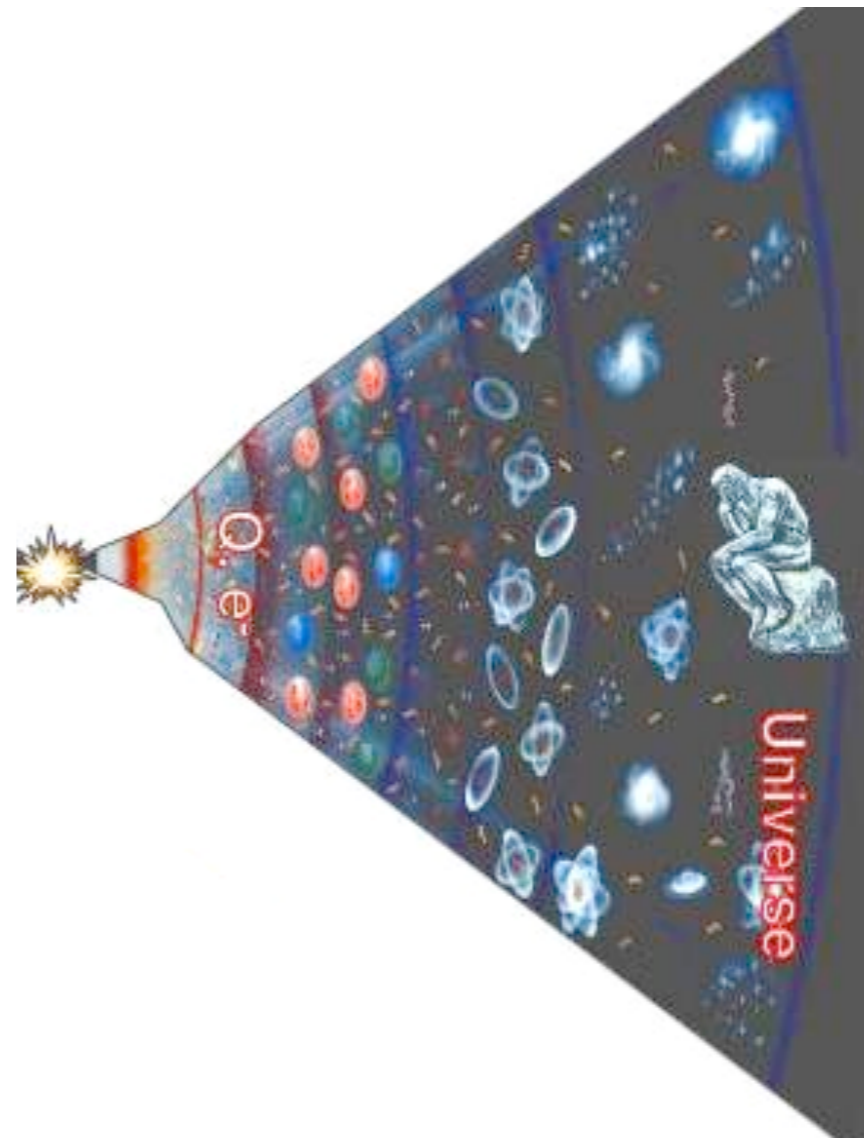
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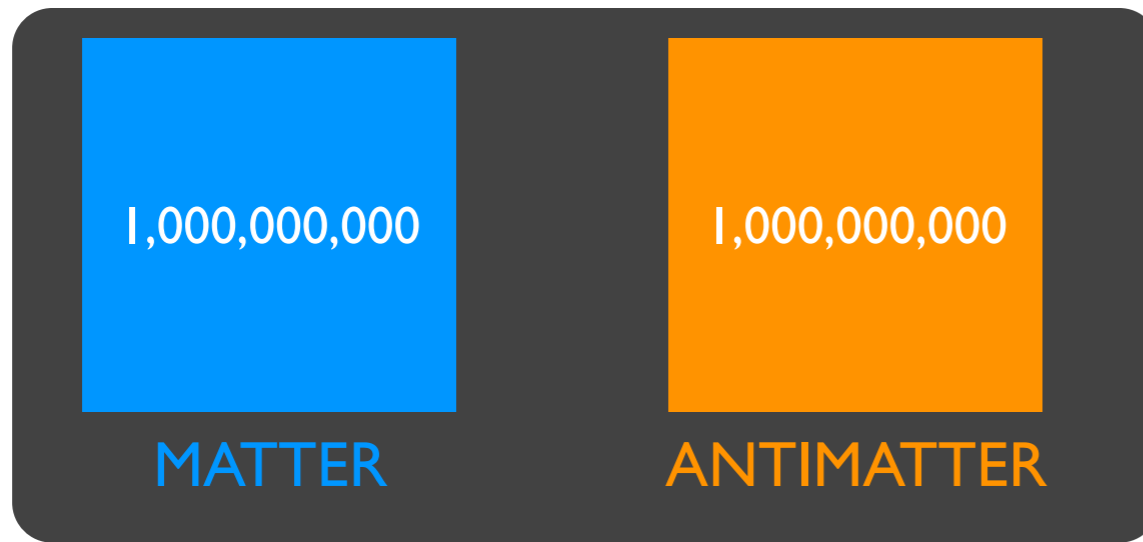
The current matter in the Universe

- No convincing evidence to date for complex antimatter in space
- Search for anti-nuclei with AMS experiment: **Anti-He / He $\approx 10^{-8}$**

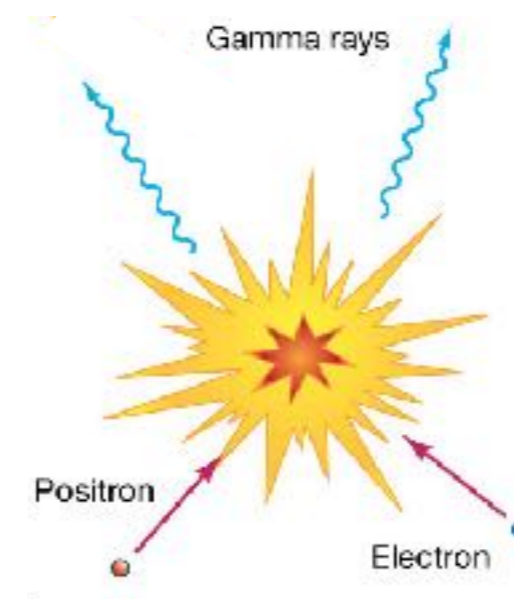
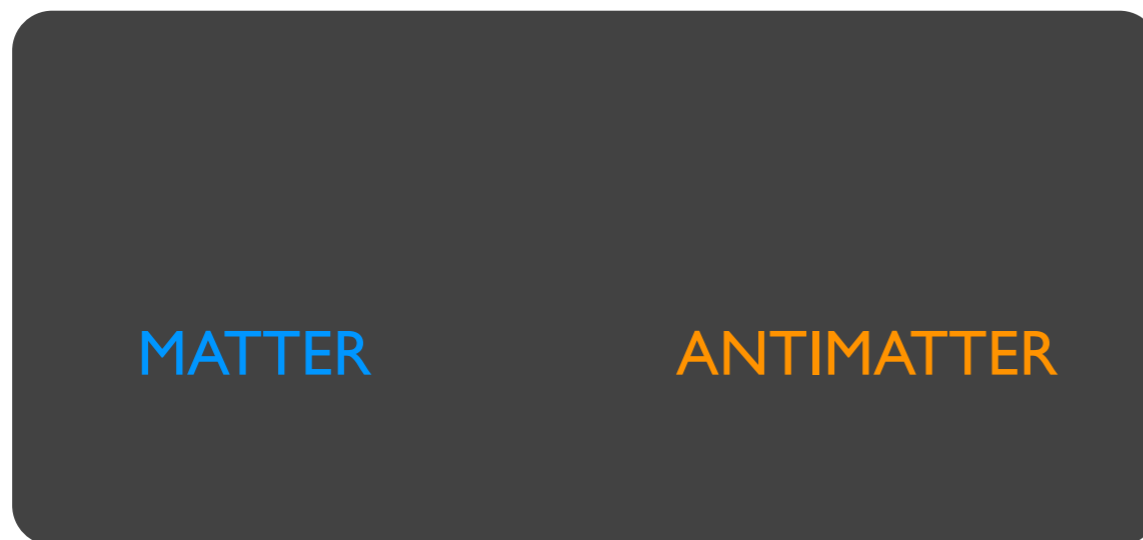


In fact, why do we see matter at all?

Baryon symmetry

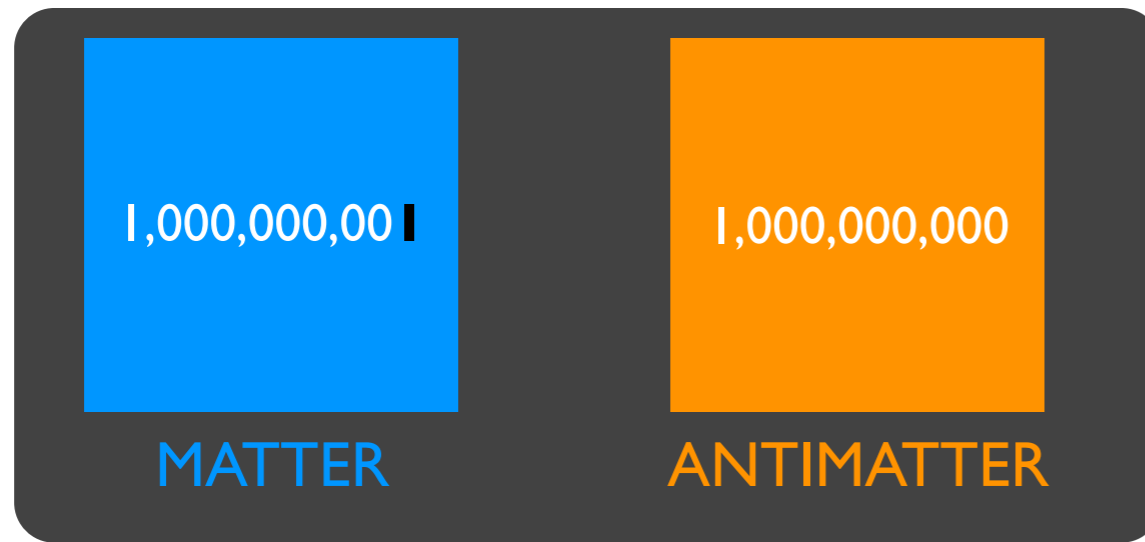


No Baryogenesis

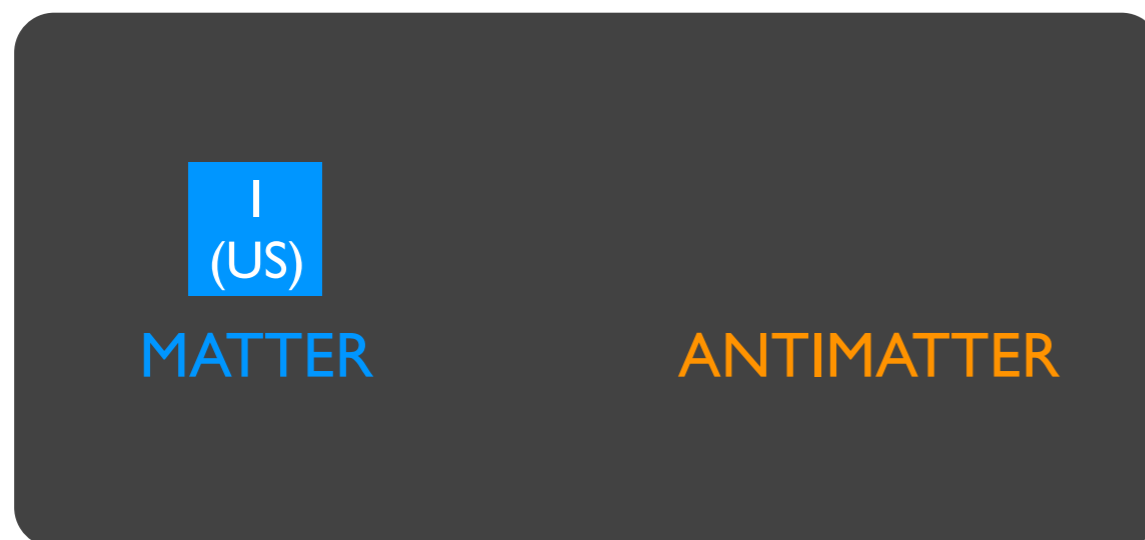


In fact, why do we see matter at all?

Baryon asymmetry



Baryogenesis



- Some physics process slightly changed matter/antimatter equilibrium in favor of matter, shortly after Big Bang

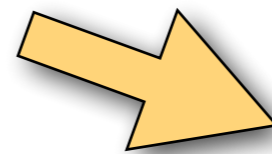
- All antimatter annihilated with matter, leaving only matter: **birth of baryons**

Baryon asymmetry in the early Universe

Experimental evidence

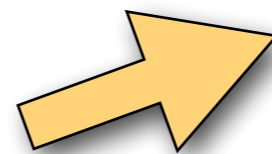
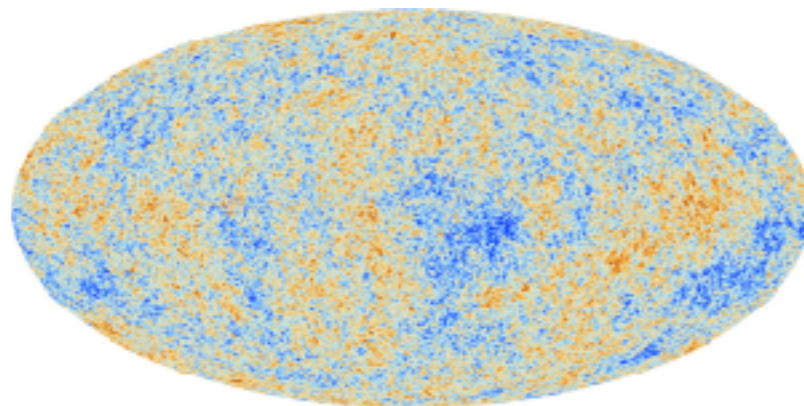
- The baryon asymmetry per unit volume, normalised to the photon number density, has not changed since a few secs after Big Bang
- It has been accurately measured via multiple probes:

BBN ($t \approx \text{min}$):



$$(n_B - n_{\bar{B}}) / n_\gamma = 6 \times 10^{-10}$$

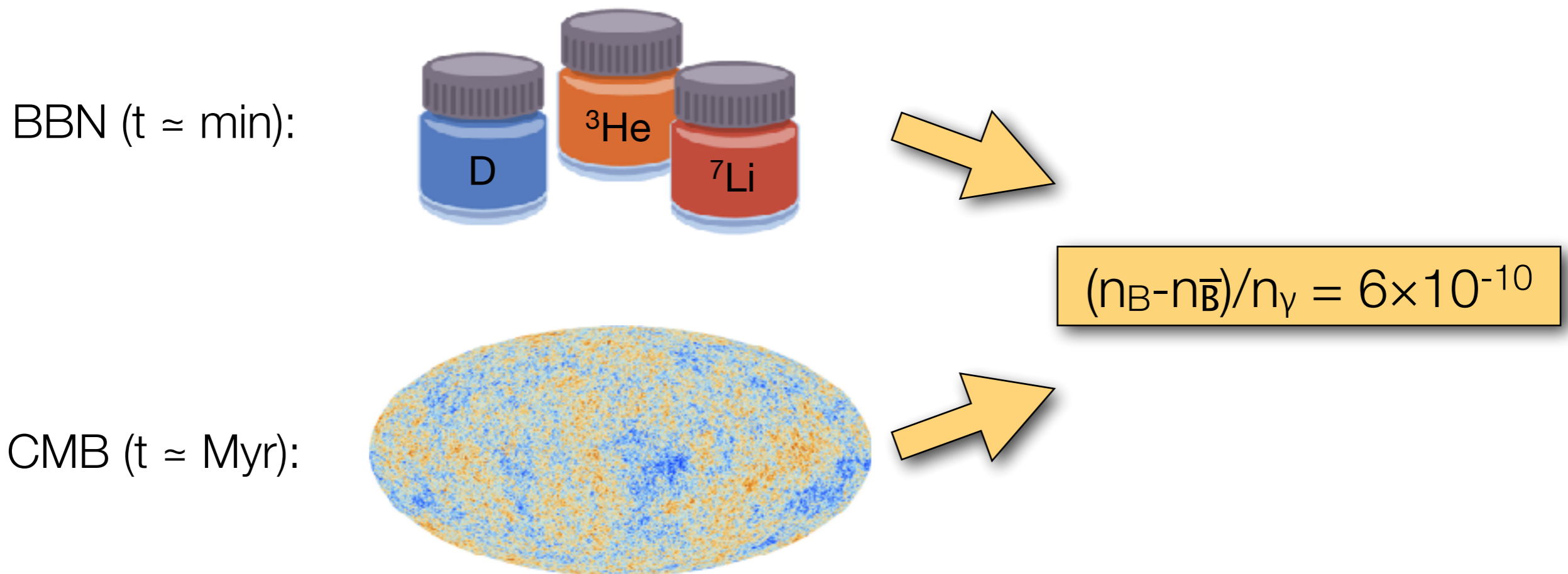
CMB ($t \approx \text{Myr}$):



Baryon asymmetry in the early Universe

Experimental evidence

- The baryon asymmetry per unit volume, normalised to the photon number density, has not changed since a few secs after Big Bang
- It has been accurately measured via multiple probes:



What process caused this baryon/antibaryon asymmetry?

Chapter 3:

*M*ajorana neutrinos to the rescue

Is a neutrino its own antiparticle?

YES



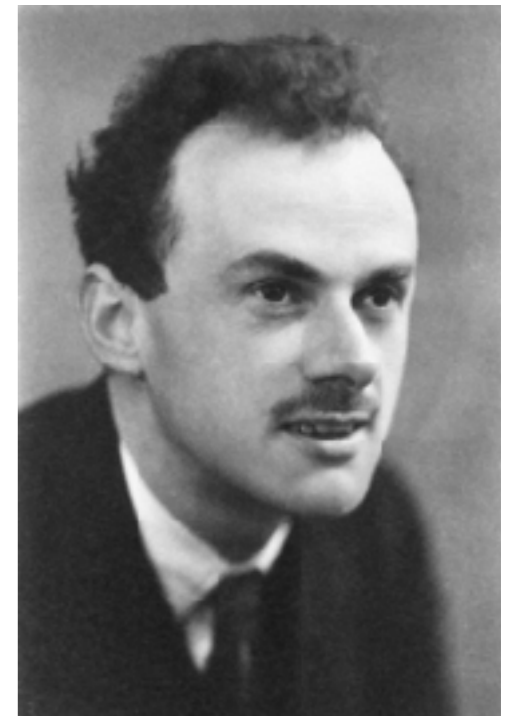
Majorana
Neutrino



NO



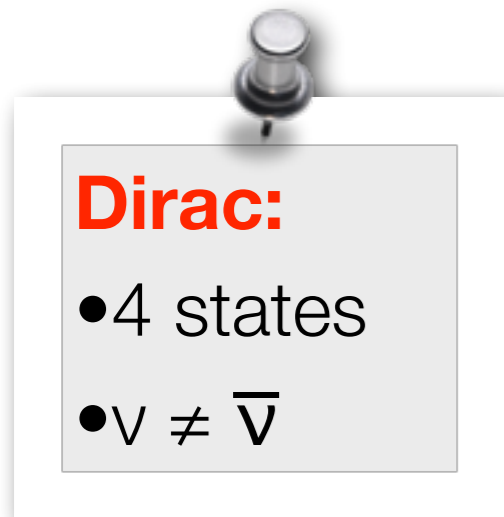
Dirac
Neutrino



- Both possibilities exist for the neutrino
- A Majorana neutrino would be unlike any other fundamental fermion: **a new form of matter**

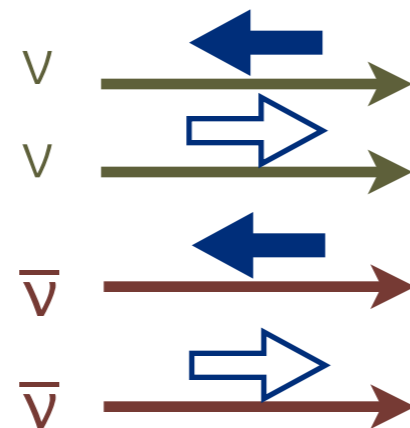
Difference between Dirac and Majorana neutrinos

Idealised neutrino scattering experiment



Dirac:

- 4 states
- $\nu \neq \bar{\nu}$



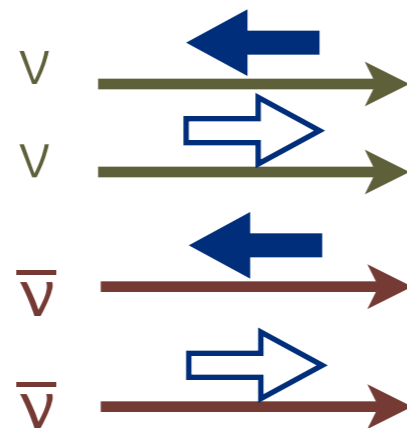
Helicity	Conserved Lepton Number	Lepton production rate	Anti-lepton production rate
-1/2	+1	1	0
+1/2	+1	$(m/E)^2 \ll 1$	0
-1/2	-1	0	$(m/E)^2 \ll 1$
+1/2	-1	0	1

Difference between Dirac and Majorana neutrinos

Idealised neutrino scattering experiment

Dirac:

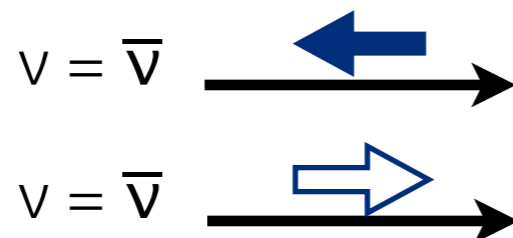
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+1/2	-1	0	1

Majorana:

- 2 states
- $\nu = \bar{\nu}$



Helicity	Conserved Lepton Number	Lepton production rate	Anti-lepton production rate
-1/2	none	1	$(m/E)^2 \ll 1$
+1/2	none	$(m/E)^2 \ll 1$	1

A prime candidate for small neutrino mass

The see-saw mechanism

- Neutrino mass matrix, with both Majorana (M) and Dirac (m_D) terms:

$$\begin{pmatrix} 0 & m_D \\ m_D & M \end{pmatrix}$$

The diagram shows a 2x2 matrix with elements 0, m_D , m_D , and M . Four dashed arrows indicate transitions: a blue arrow from ν_L to $\bar{\nu}_R$ pointing to the top-left element (0); a red arrow from ν_L to ν_R pointing to the top-right element (m_D); a red arrow from $\bar{\nu}_R$ to $\bar{\nu}_L$ pointing to the bottom-left element (m_D); and a blue arrow from ν_R to $\bar{\nu}_L$ pointing to the bottom-right element (M).

- Majorana terms induce $|\Delta L| = 2$ lepton number violating processes and imply $\nu = \bar{\nu}$
- M : **a new physics scale**

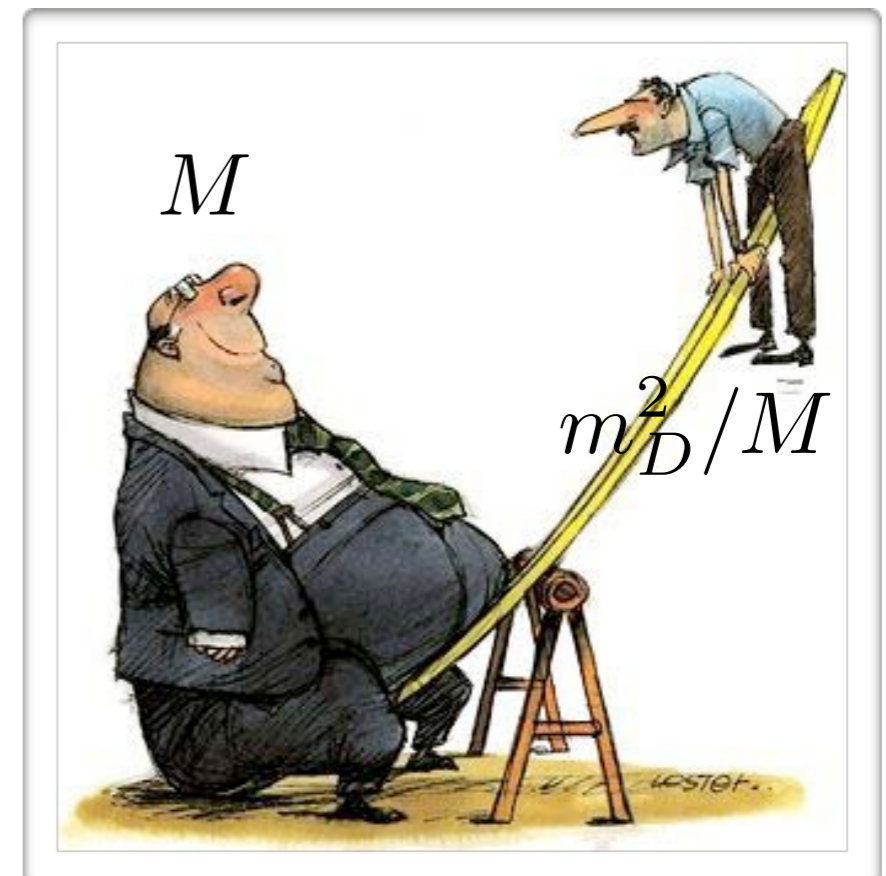
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The see-saw mechanism

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- Majorana terms induce $|\Delta L| = 2$ lepton number violating processes and imply $\nu = \bar{\nu}$
- M : **a new physics scale**
- "See-saw" limit $m_D \ll M$ explains small neutrino masses, which indirectly probe new physics scale

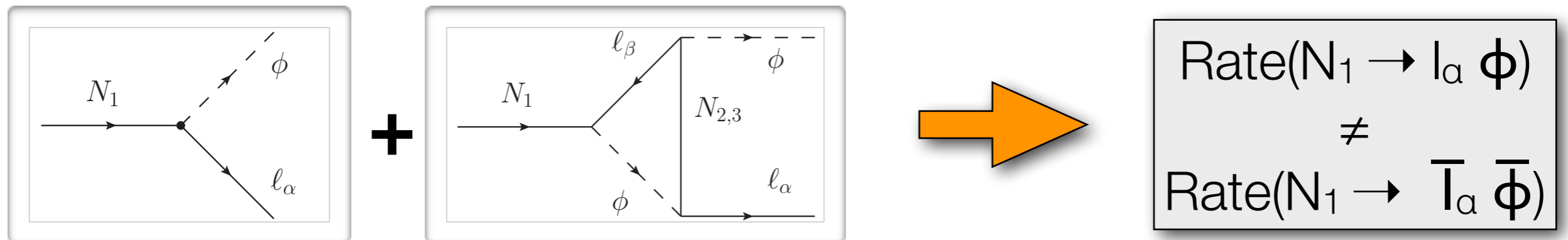
$$\begin{pmatrix} 0 & m_D \\ m_D & M \end{pmatrix}$$

ν_L to $\bar{\nu}_R$ (top-left to top-right)
 $\bar{\nu}_R$ to $\bar{\nu}_L$ (bottom-left to bottom-right)
 ν_L to ν_R (top-left to bottom-right)
 ν_R to $\bar{\nu}_L$ (bottom-right to top-left)



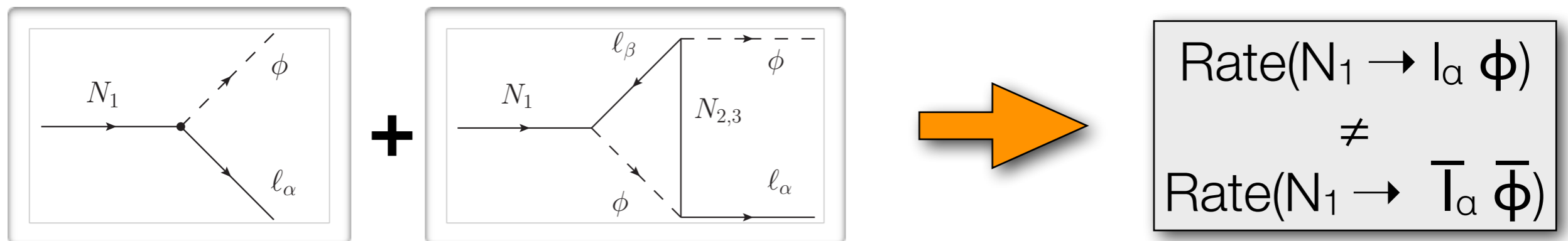
A prime candidate for the baryon asymmetry

- Baryon asymmetry possibly induced by a lepton asymmetry: *leptogenesis*
- *Decay of heavy Majorana neutrinos* ideal for leptogenesis, if CP is violated

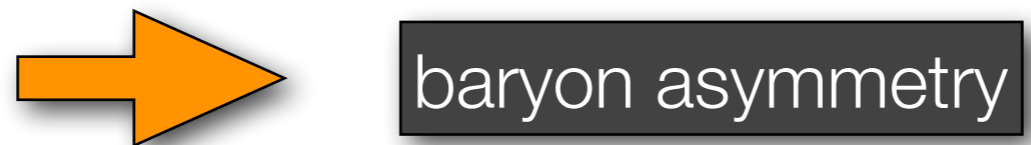


A prime candidate for the baryon asymmetry

- Baryon asymmetry possibly induced by a lepton asymmetry: *leptogenesis*
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- Unequal number of leptons and antileptons is later transferred to baryons:



How to find out if neutrinos are Majorana?

Play...

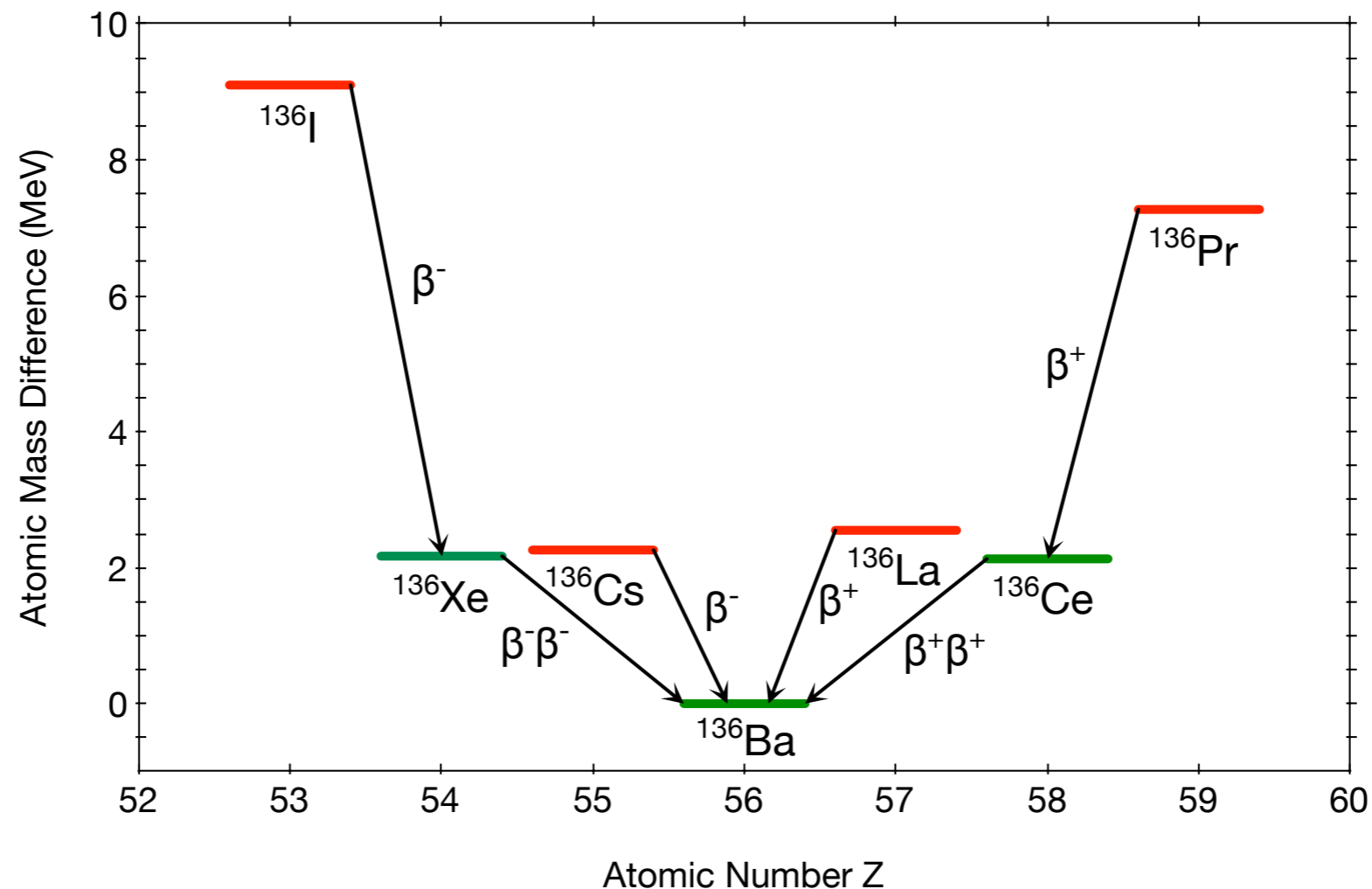


Chapter 4:

Searching for Majorana neutrinos with $\beta\beta 0\nu$

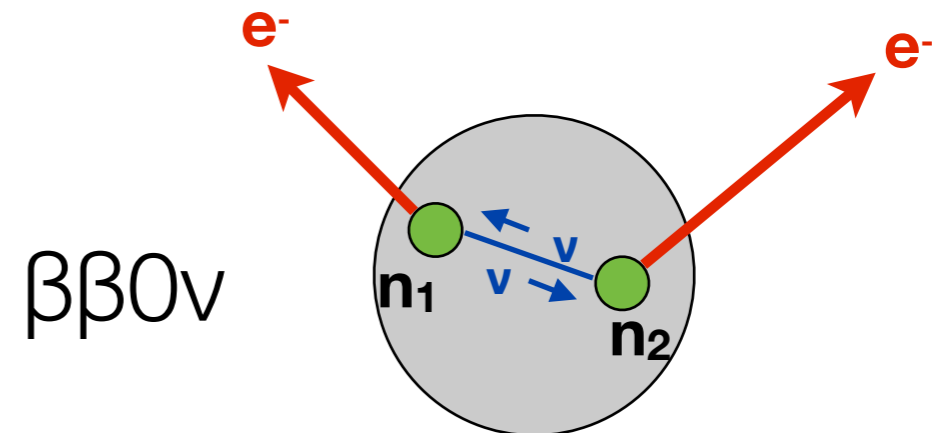
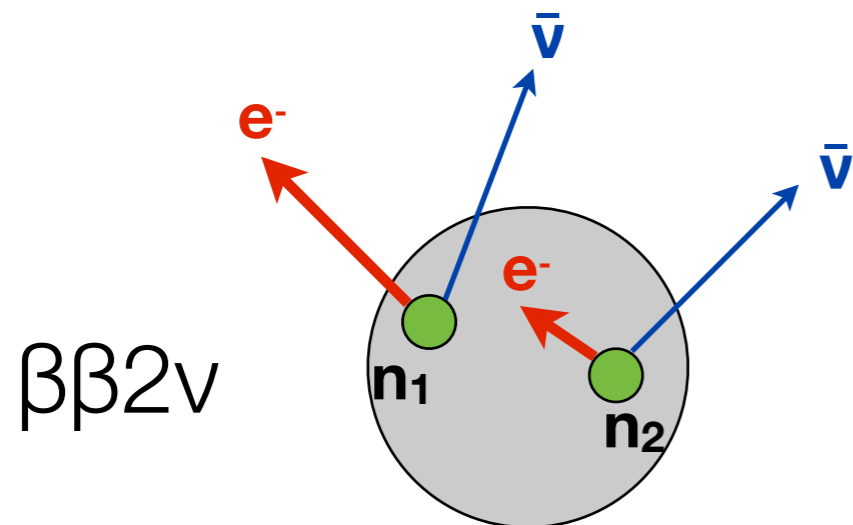
Nuclear double beta decay

- Nuclear $(Z,A) \rightarrow (Z+2,A)$ transition with emission of two electrons. Second order process mediated by the weak interaction
- This process exists in 35 nuclides due to nuclear pairing interaction
→ favours energetically the **even-even isobars** over the **odd-odd ones**.



Double beta decay modes

- Two basic decay modes:



Two neutrino mode

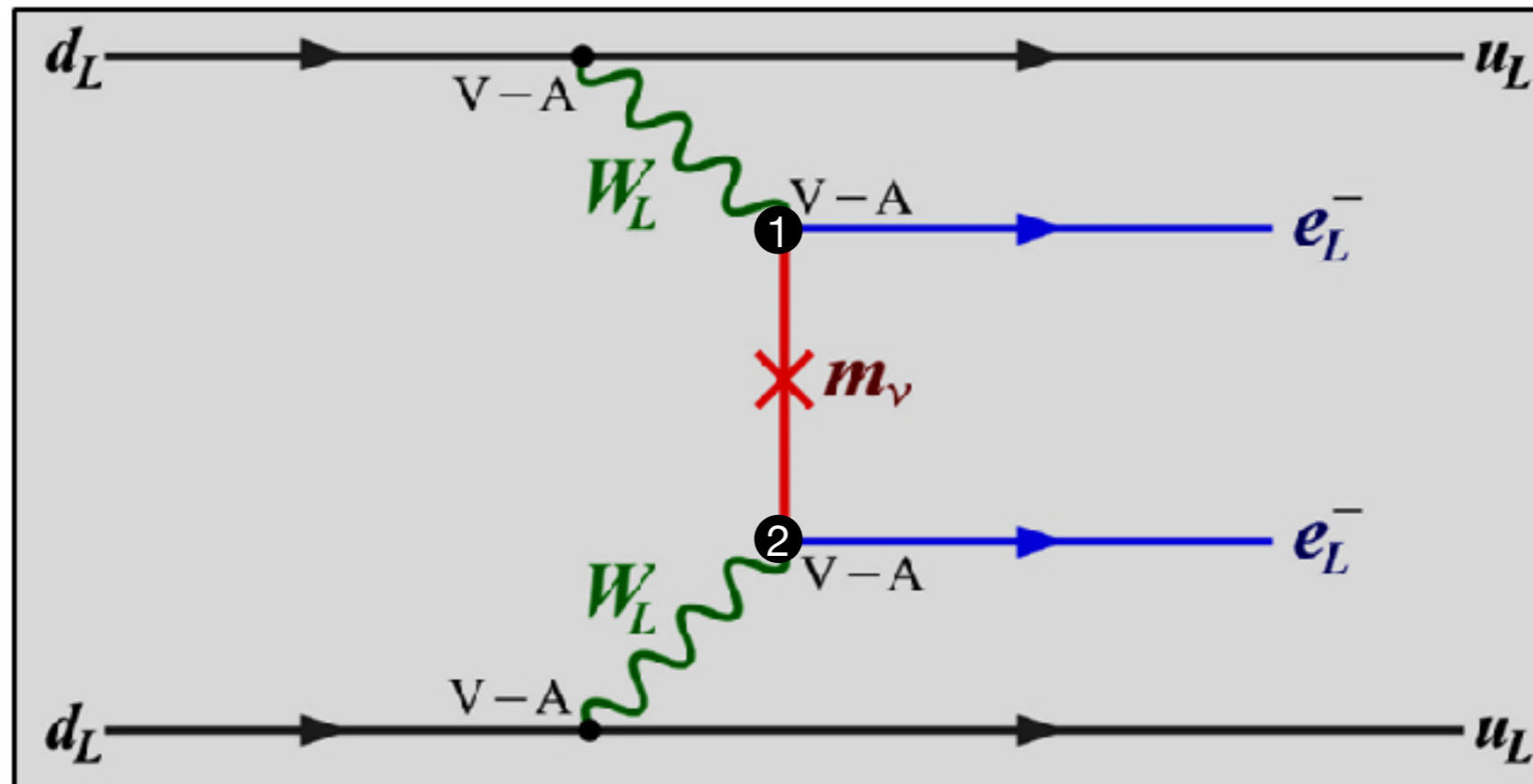
- Observed in several nuclei
- 10^{19} - 10^{21} yr half-lives
- Standard Model allowed
- Conserves lepton number

Neutrinoless mode

- Not observed yet in Nature
- $>10^{26}$ yr half-lives
- Would signal BSM physics
- Violates **L** by two units

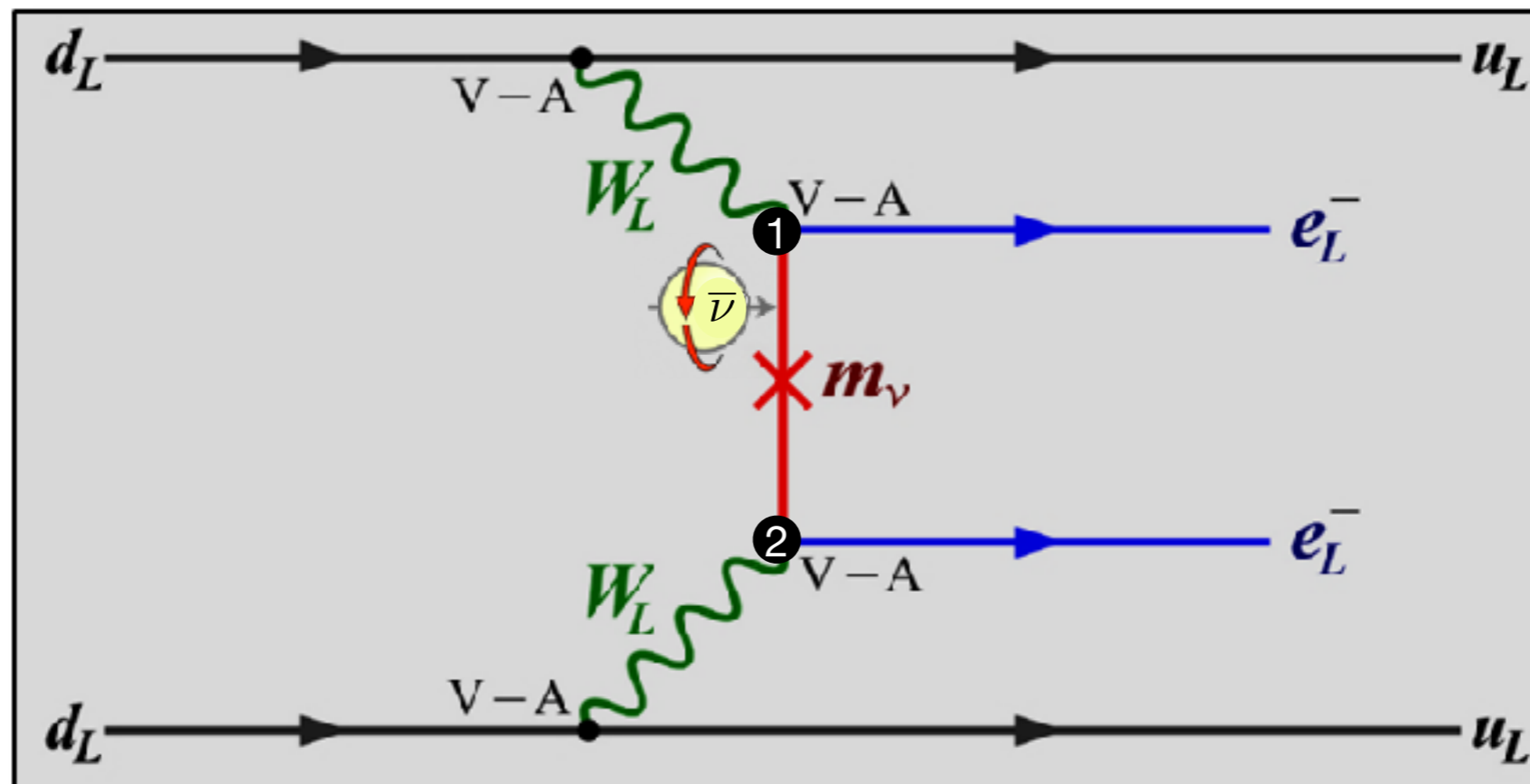
$\beta\beta 0\nu$ and Majorana neutrinos

- $\beta\beta 0\nu$ evidence would imply that neutrinos are massive, Majorana, particles



$\beta\beta 0\nu$ and Majorana neutrinos

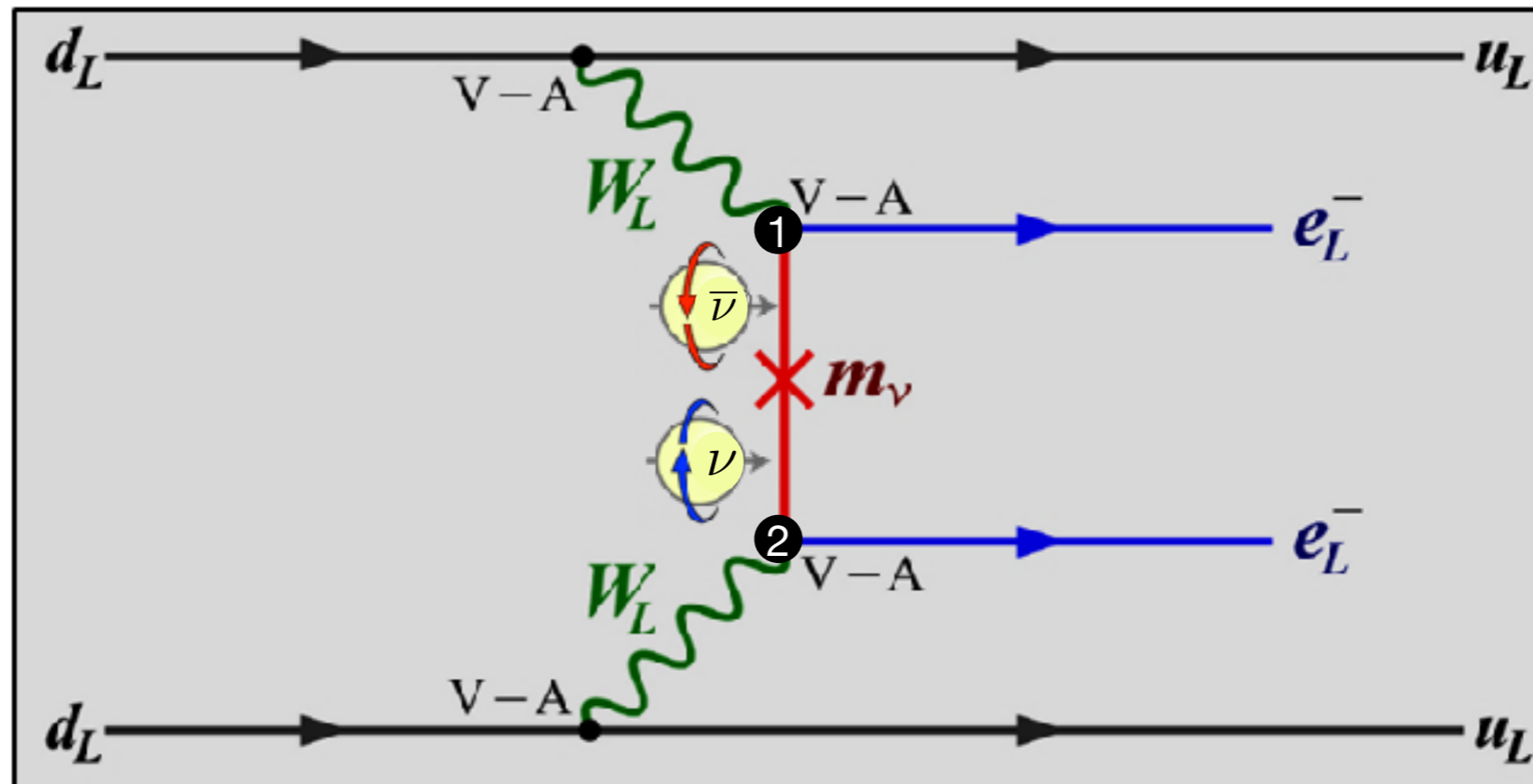
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Emitted in **1**, in association with electron, with almost total positive helicity

$\beta\beta 0\nu$ and Majorana neutrinos

- $\beta\beta 0\nu$ evidence would imply that neutrinos are massive, Majorana, particles



Emitted in **1**, in association with electron, with almost total positive helicity

Only its small, $\mathcal{O}(m/E)$, negative helicity component absorbed in **2**, producing another electron

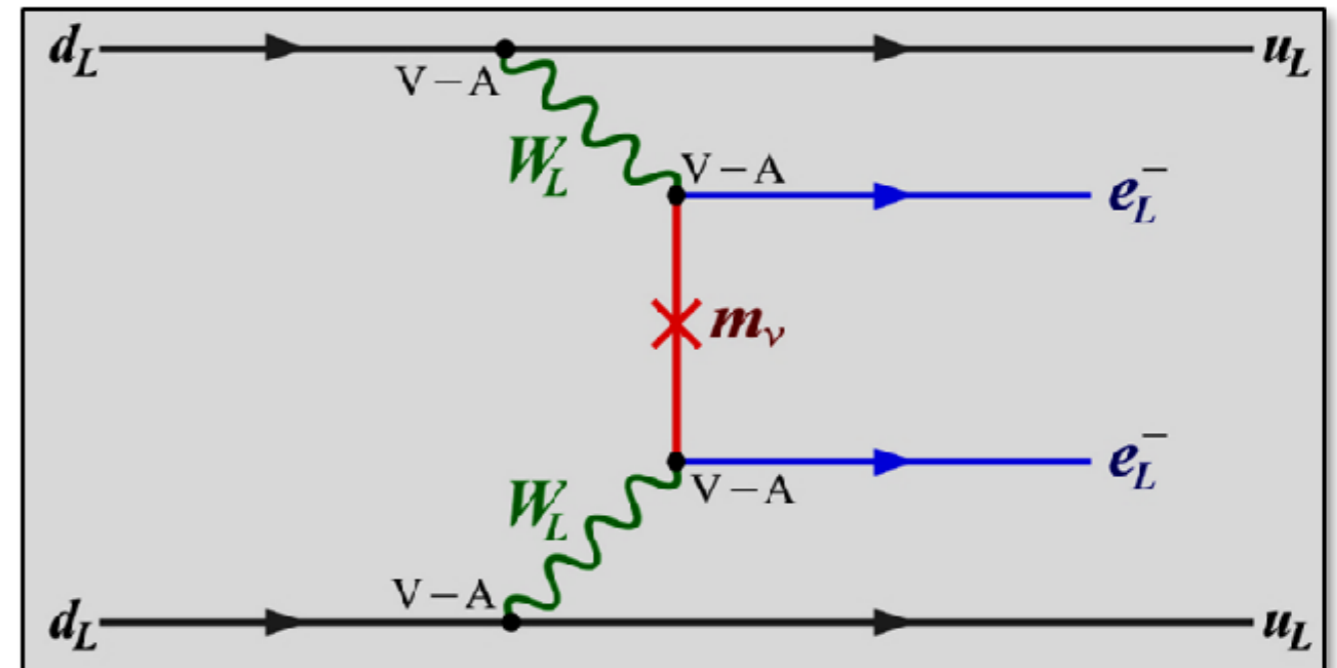
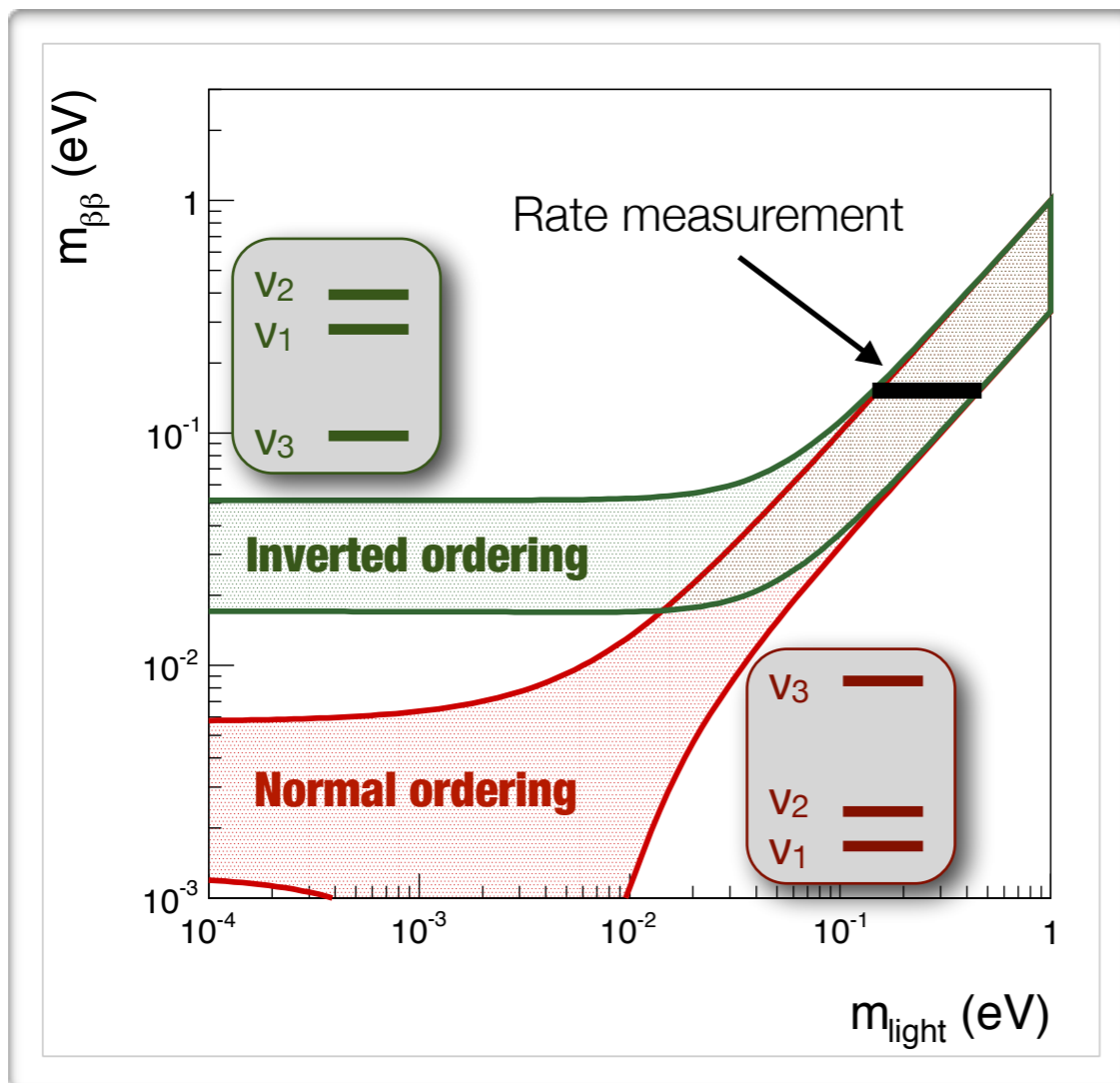
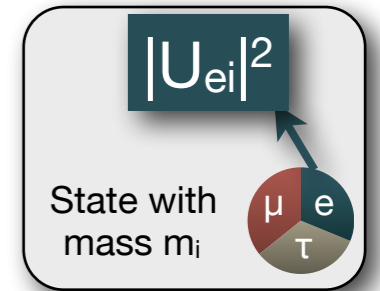
$\beta\beta 0\nu$ and neutrino mass

- $\beta\beta 0\nu$ rate constrains neutrino mass:

$$(\text{Rate})_{\beta\beta 0\nu} \propto m_{\beta\beta}^2$$

Majorana ν mass:

$$m_{\beta\beta} \equiv \left| \sum_i m_i U_{ei}^2 \right|$$



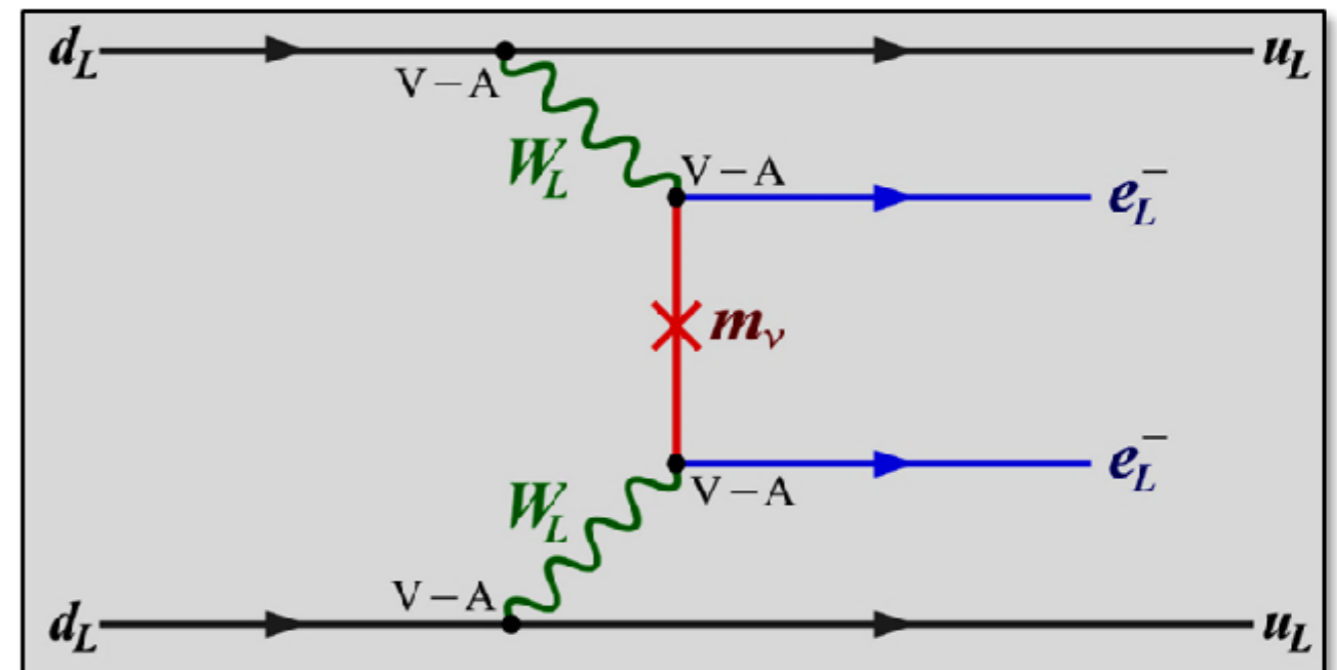
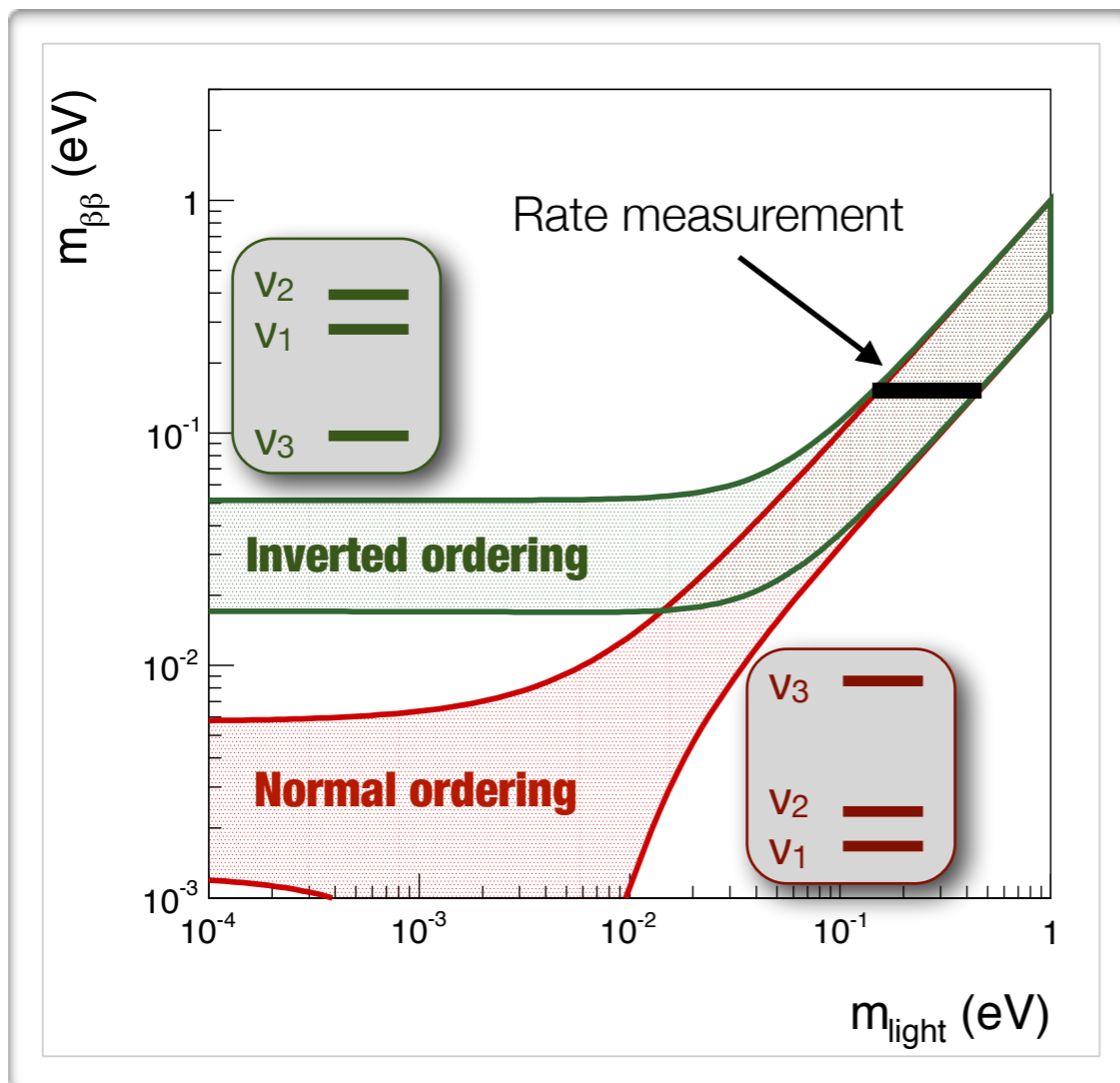
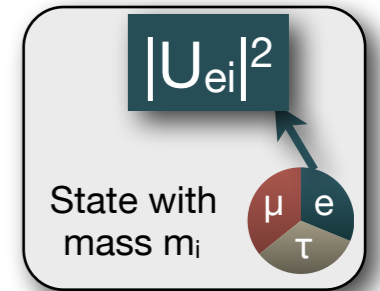
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Majorana ν mass:

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$\beta\beta 0\nu$: most sensitive probe for Majorana neutrinos \rightarrow low $m_{\beta\beta}$ reach!

Chapter 5:

*E*xperimental challenges in $\beta\beta 0\nu$ searches

Facts of life of the double beta decay experimentalist

- Total number of $\beta\beta_{0\nu}$ decays that can be observed in a detector is:

$$N_{\beta\beta_{0\nu}} = \log 2 \cdot \frac{M_{\beta\beta} \cdot N_A}{W_{\beta\beta}} \cdot \varepsilon \cdot \frac{t}{T_{1/2}^{0\nu}}$$

The diagram shows the equation $N_{\beta\beta_{0\nu}} = \log 2 \cdot \frac{M_{\beta\beta} \cdot N_A}{W_{\beta\beta}} \cdot \varepsilon \cdot \frac{t}{T_{1/2}^{0\nu}}$ enclosed in a box. Red arrows point from each variable to its definition: $M_{\beta\beta}$ to 'mass of $\beta\beta$ isotope', N_A to 'Avogadro's constant', $W_{\beta\beta}$ to 'Molar mass of $\beta\beta$ isotope', ε to 'Efficiency', t to 'Exposure time', and $T_{1/2}^{0\nu}$ to ' $\beta\beta_{0\nu}$ half-life'.

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- For ^{136}Xe experiment with 100% efficiency and 1 year exposure time, what is the mass $M_{\beta\beta}$ required to observe only one $\beta\beta_{0\nu}$ decay?

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Diagram illustrating the components of the equation for the total number of $\beta\beta_{0\nu}$ decays observed in a detector:

- $M_{\beta\beta}$: mass of $\beta\beta$ isotope
- N_A : Avogadro's constant
- $W_{\beta\beta}$: Molar mass of $\beta\beta$ isotope
- ε : Efficiency
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- $T_{1/2}^{0\nu}$: $\beta\beta_{0\nu}$ half-life

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- If ^{136}Xe $\beta\beta_{0\nu}$ half-life is $T_{1/2} = 10^{27}$ years, get: **$M_{\beta\beta} = 326 \text{ kg!}$**

Facts of life of the double beta decay experimentalist

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- If ^{136}Xe $\beta\beta_{0\nu}$ half-life is $T_{1/2} = 10^{27}$ years, get: **$M_{\beta\beta} = 326 \text{ kg!}$**
- Life is harder than this: non-perfect efficiencies and backgrounds

Experimental sensitivity to $\beta\beta 0\nu$

- Experiment with no background:

detector efficiency

$$T_{1/2}^{0\nu} \propto \varepsilon \cdot M_{\beta\beta} \cdot t$$

exposure
(mass×time)

Experimental sensitivity to $\beta\beta 0\nu$

- Experiment with no background:

$$T_{1/2}^{0\nu} \propto \epsilon \cdot M_{\beta\beta} \cdot t$$

detector efficiency

exposure (mass×time)

- Experiment with background:

$$T_{1/2}^{0\nu} \propto \epsilon \cdot \sqrt{\frac{M_{\beta\beta} \cdot t}{c \cdot \Delta E}}$$

detector efficiency

exposure (mass×time)

energy resolution

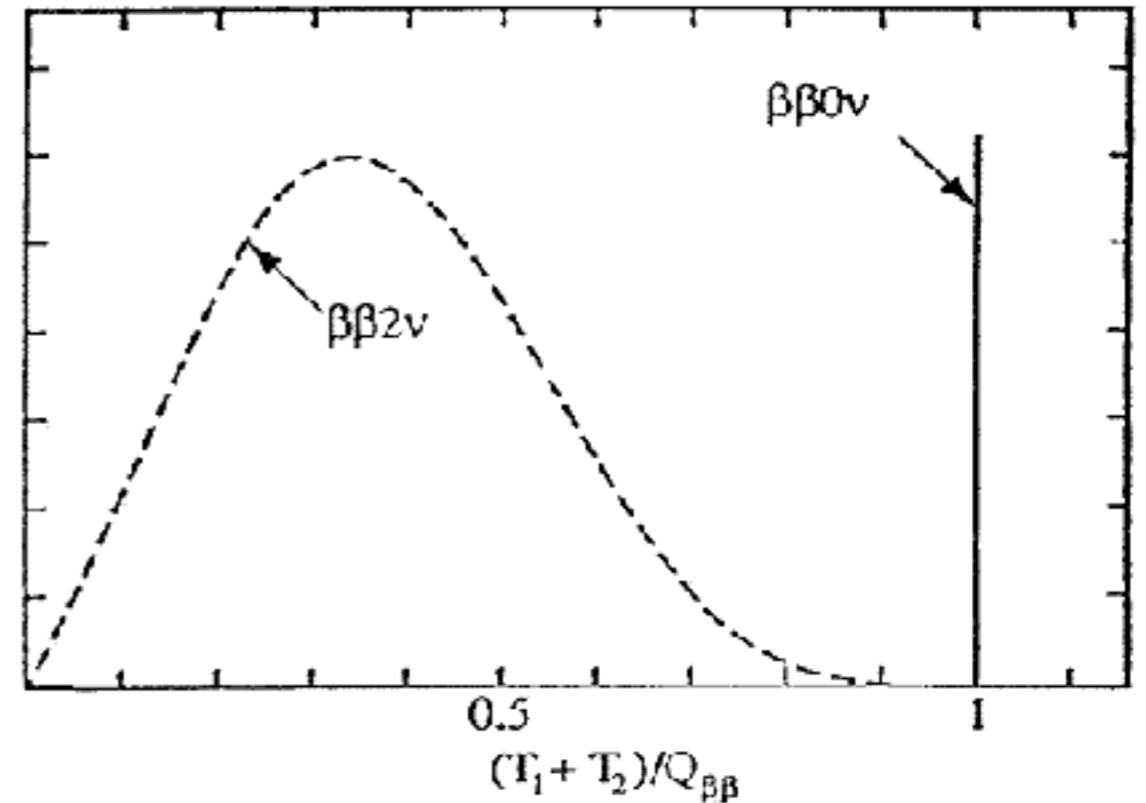
background rate (per unit energy, mass and time)

$\beta\beta 0\nu$ experimental signature

Rare process to be isolated in radio-pure detector underground

1. Calorimetry (A MUST):

- *2 ν mode*: continuous spectrum for sum electron kinetic energy T_1+T_2
- *0 ν mode*: mono-energetic line at $Q_{\beta\beta}$ for T_1+T_2 spectrum

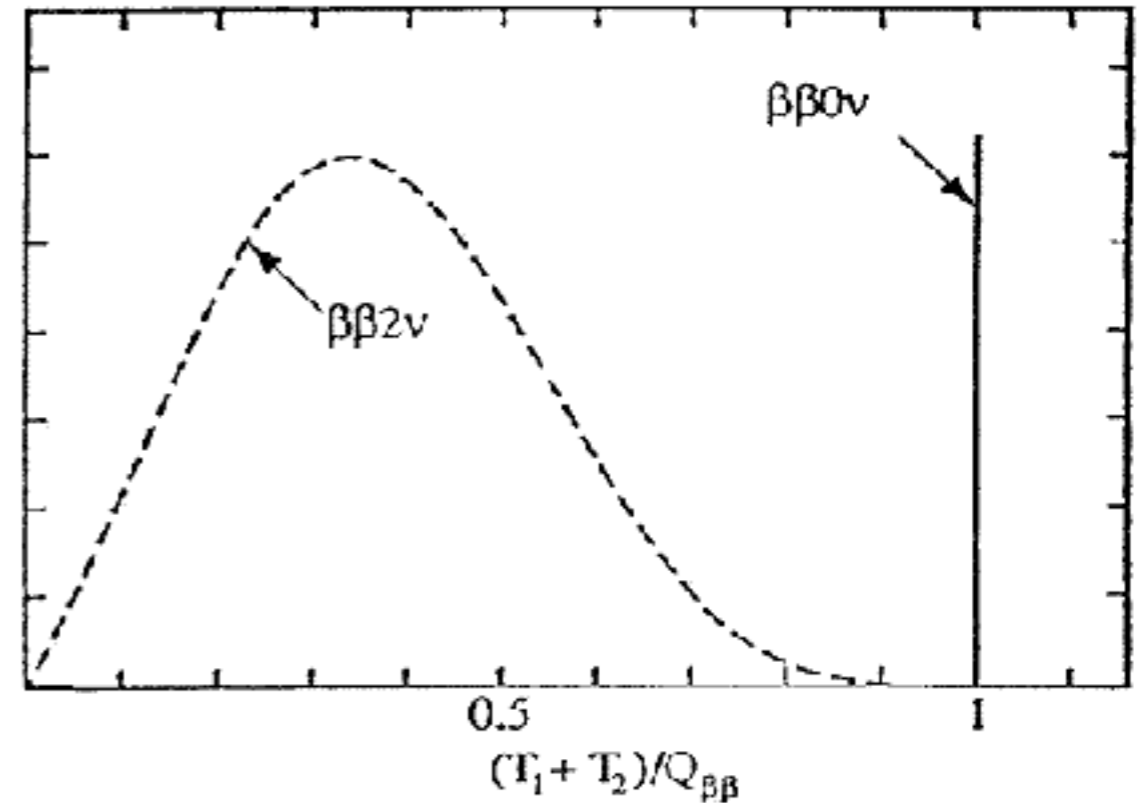


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2. Topology of decay electrons (AN ADDITIONAL HANDLE):

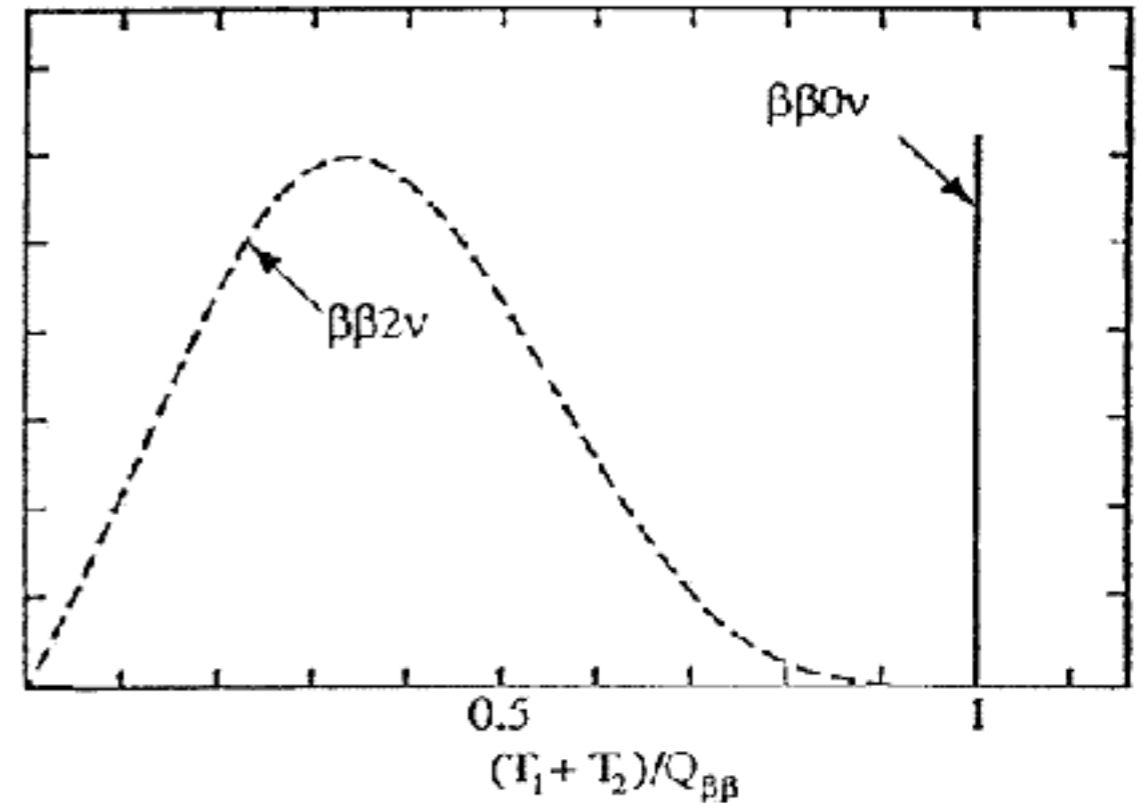
- Two electrons from common vertex

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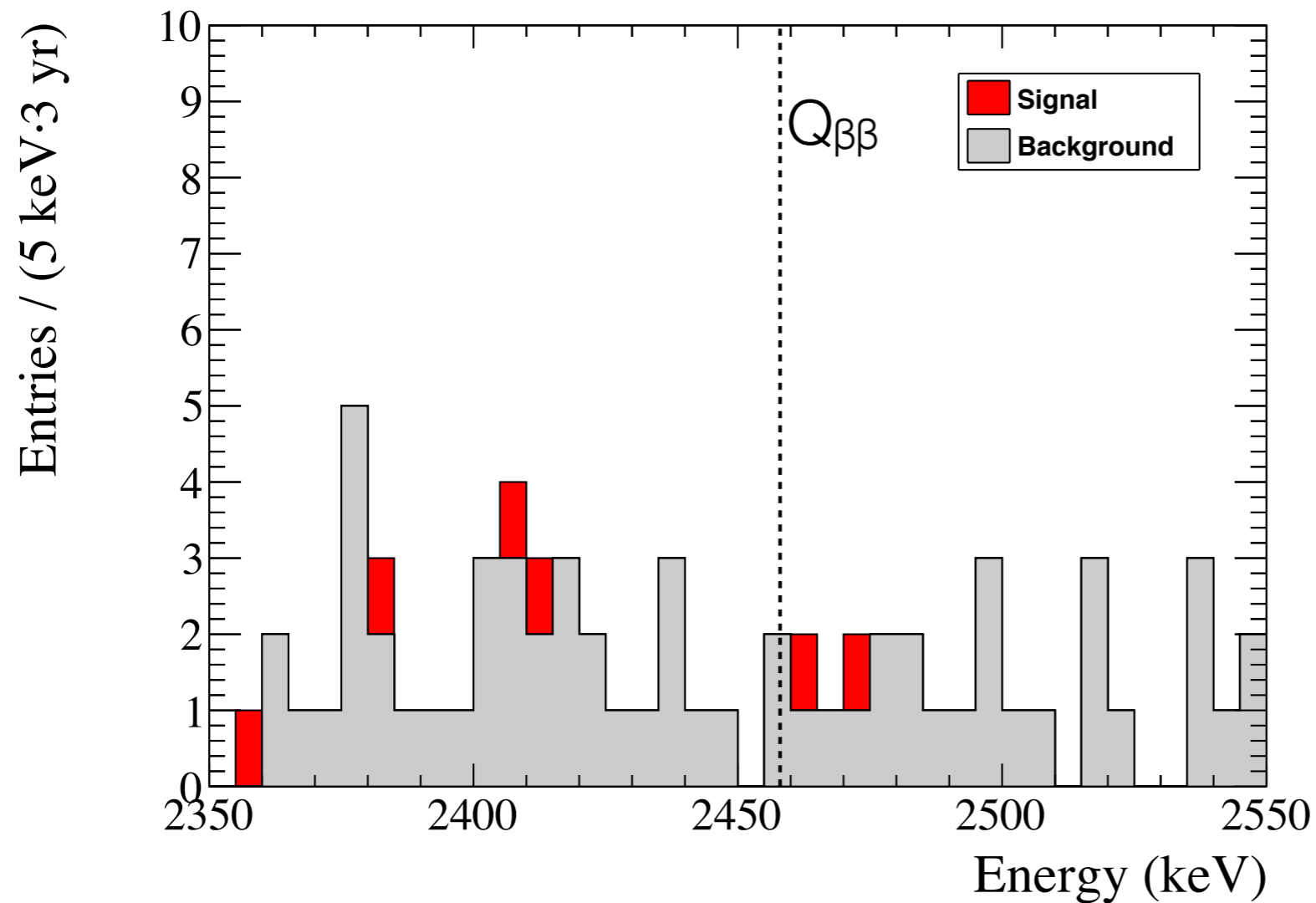
- Two electrons from common vertex

3. Daughter ion tagging (A DREAM):

- Observe ion produced in the decay

How can we uncover the $Q_{\beta\beta}$ peak in the energy spectrum?

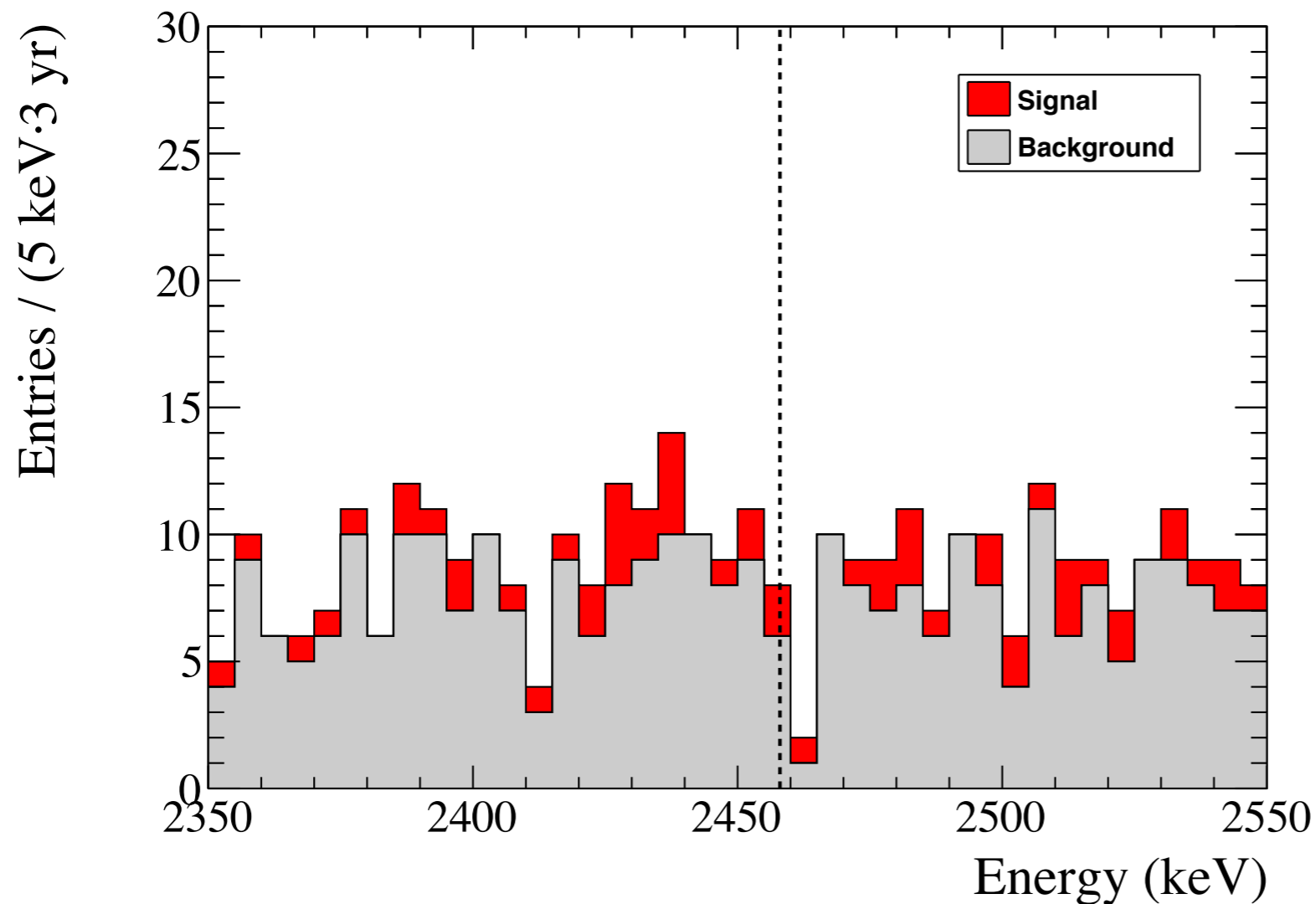
- Typical situation for current-generation detector performance, assuming $T_{1/2}^{0\nu} = 10^{26}$ yr:



- $M = 100$ kg
- $t = 3$ yr
- $T_{1/2} = 10^{26}$ yr
- $\epsilon = 100\%$
- $\Delta E = 10\%$ FWHM
- $c = 10^{-3}$ / (keV·kg·yr)

Improvement no.1: larger detector

$$T_{1/2}^{0\nu} \propto \varepsilon \sqrt{\frac{Mt}{c \Delta E}}$$



- **M = 1,000 kg**
- t = 3 yr
- $T_{1/2} = 10^{26}$ yr
- $\varepsilon = 100\%$
- $\Delta E = 10\%$ FWHM
- **c = $5 \cdot 10^{-4}$ / (keV · kg · yr)**

- More events! Also: signal \propto volume, background \propto surface \rightarrow S/B \nearrow
- Mass scalability depends on chosen $\beta\beta$ isotope

Comparison of $\beta\beta$ isotopes

- $\beta\beta$ isotope choice also affects relationship ($\beta\beta 0\nu$ rate \leftrightarrow Majorana mass):

atomic, nuclear, particle physics

$$1/T_{1/2}^{0\nu} = G^{0\nu} \cdot |M^{0\nu}|^2 \cdot m_{\beta\beta}^2$$

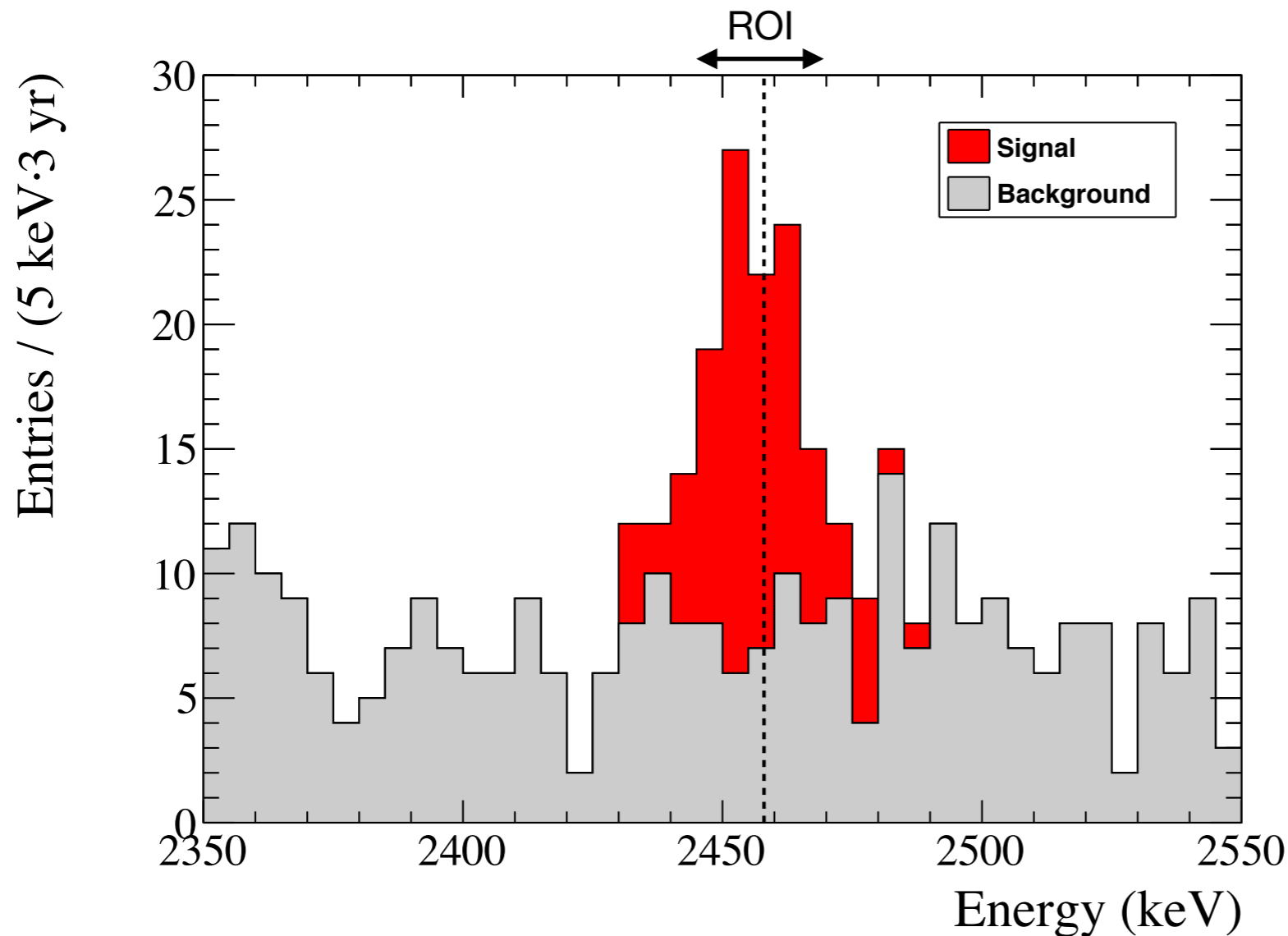
Isotope	Q-value (MeV)	Phase space $G^{0\nu}$ ($\text{yr}^{-1} \text{eV}^{-2}$)	Matrix element $ M^{0\nu} $	Isotopic abundance (%)	Cost (normalized to ^{76}Ge)	Current experiments
^{76}Ge	2.04	3.0×10^{-26}	≈ 4.1	7.8	1	GERDA, Majorana
^{130}Te	2.53	2.1×10^{-25}	≈ 3.6	33.8	0.2	CUORE, SNO+
^{136}Xe	2.46	2.3×10^{-25}	≈ 2.8	8.9	0.1	EXO, KamLAND-Zen, NEXT

The higher, the better

The lower, the better

Improvement no.2: better energy resolution

$$T_{1/2}^{0\nu} \propto \varepsilon \sqrt{\frac{Mt}{c \Delta E}}$$



- $M = 1,000 \text{ kg}$
- $t = 3 \text{ yr}$
- $T_{1/2} = 10^{26} \text{ yr}$
- $\varepsilon = 100\%$
- **$\Delta E = 1\% \text{ FWHM}$**
- $c = 5 \cdot 10^{-4} / (\text{keV} \cdot \text{kg} \cdot \text{yr})$

- Experiments define energy **R**egion **O**f **I**nterest near $Q_{\beta\beta}$. ROI width depends on energy resolution (1 FWHM typical)
- The better the resolution, the lower the background within the ROI!

Energy resolution versus background type

$\beta\beta 0\nu$ backgrounds unrelated to $\beta\beta$ source: contamination of detector components, cosmogenics, etc.

- Can be eliminated in experiment with average energy resolution, provided perfect shielding ($c \sim 0$) available

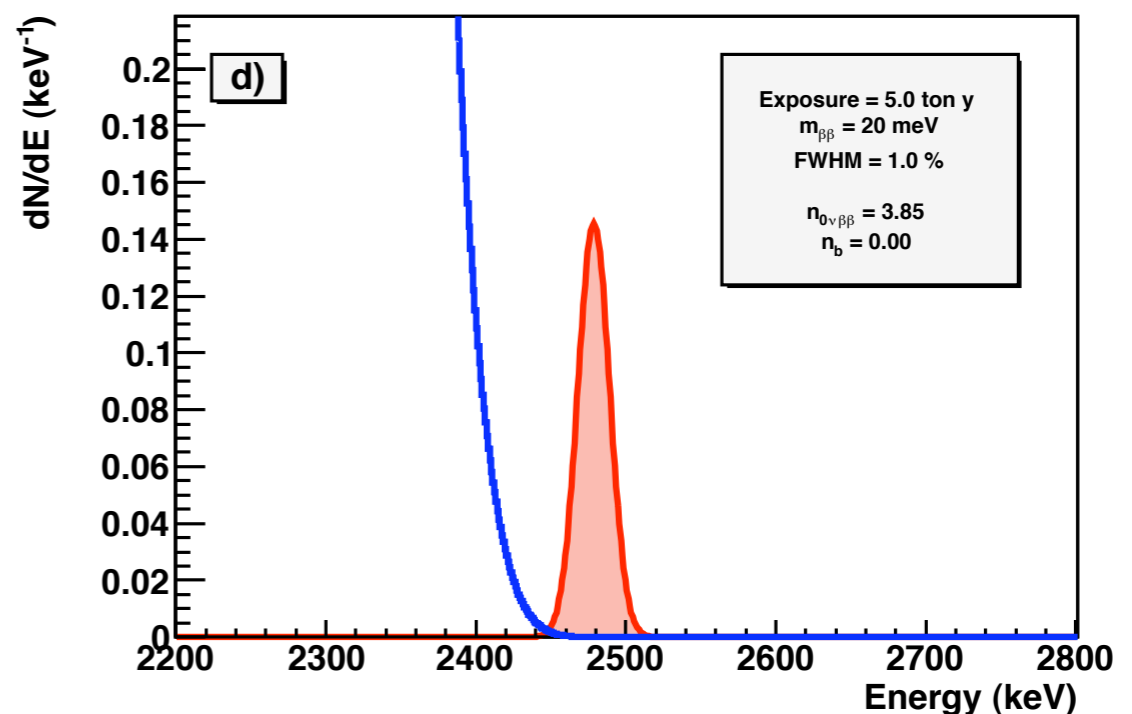
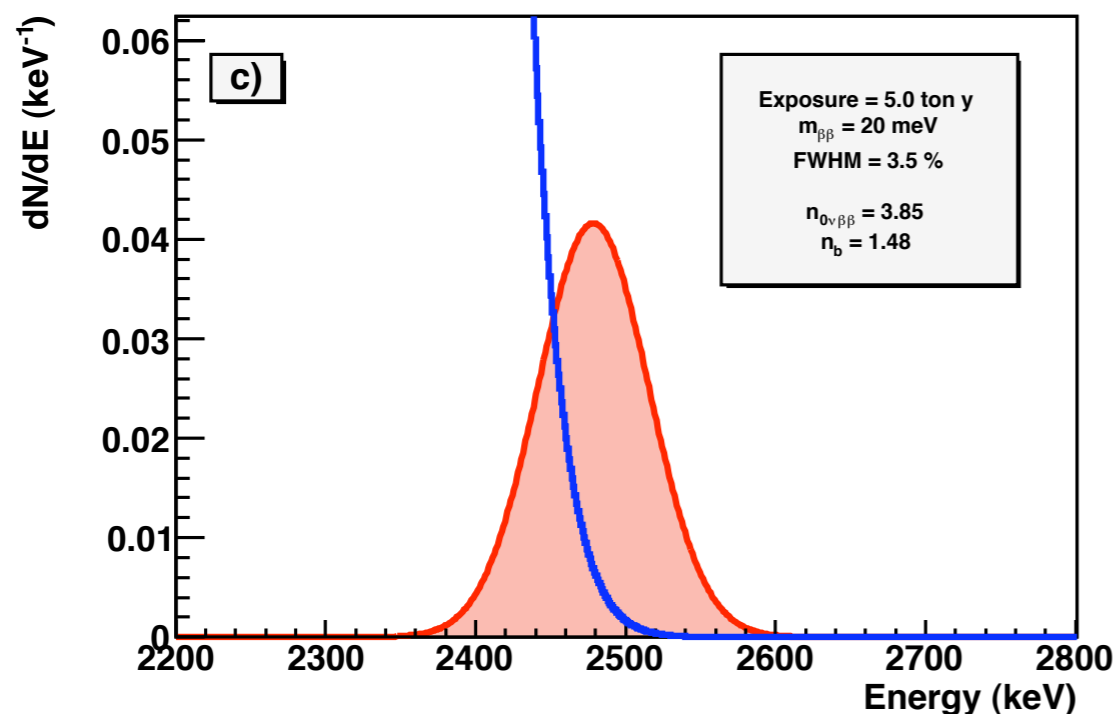
Energy resolution versus background type

$\beta\beta 0\nu$ backgrounds unrelated to $\beta\beta$ source: contamination of detector components, cosmogenics, etc.

- Can be eliminated in experiment with average energy resolution, provided perfect shielding ($c \sim 0$) available

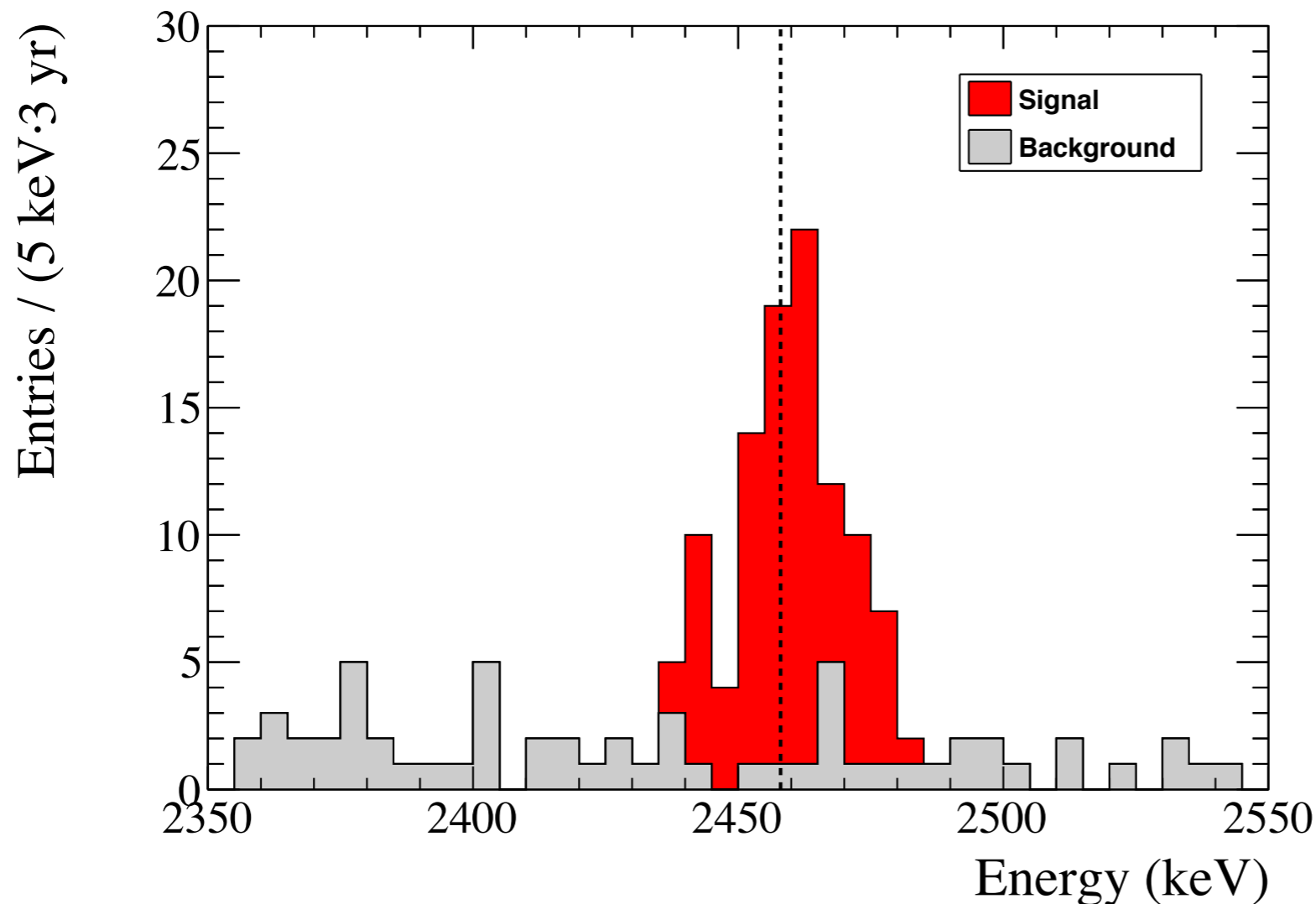
$\beta\beta 0\nu$ backgrounds related to $\beta\beta$ source: $\beta\beta 2\nu$!

- Irreducible background unless resolution is excellent



Improvement no.3: lower background rate

$$T_{1/2}^{0\nu} \propto \varepsilon \sqrt{\frac{Mt}{c \Delta E}}$$

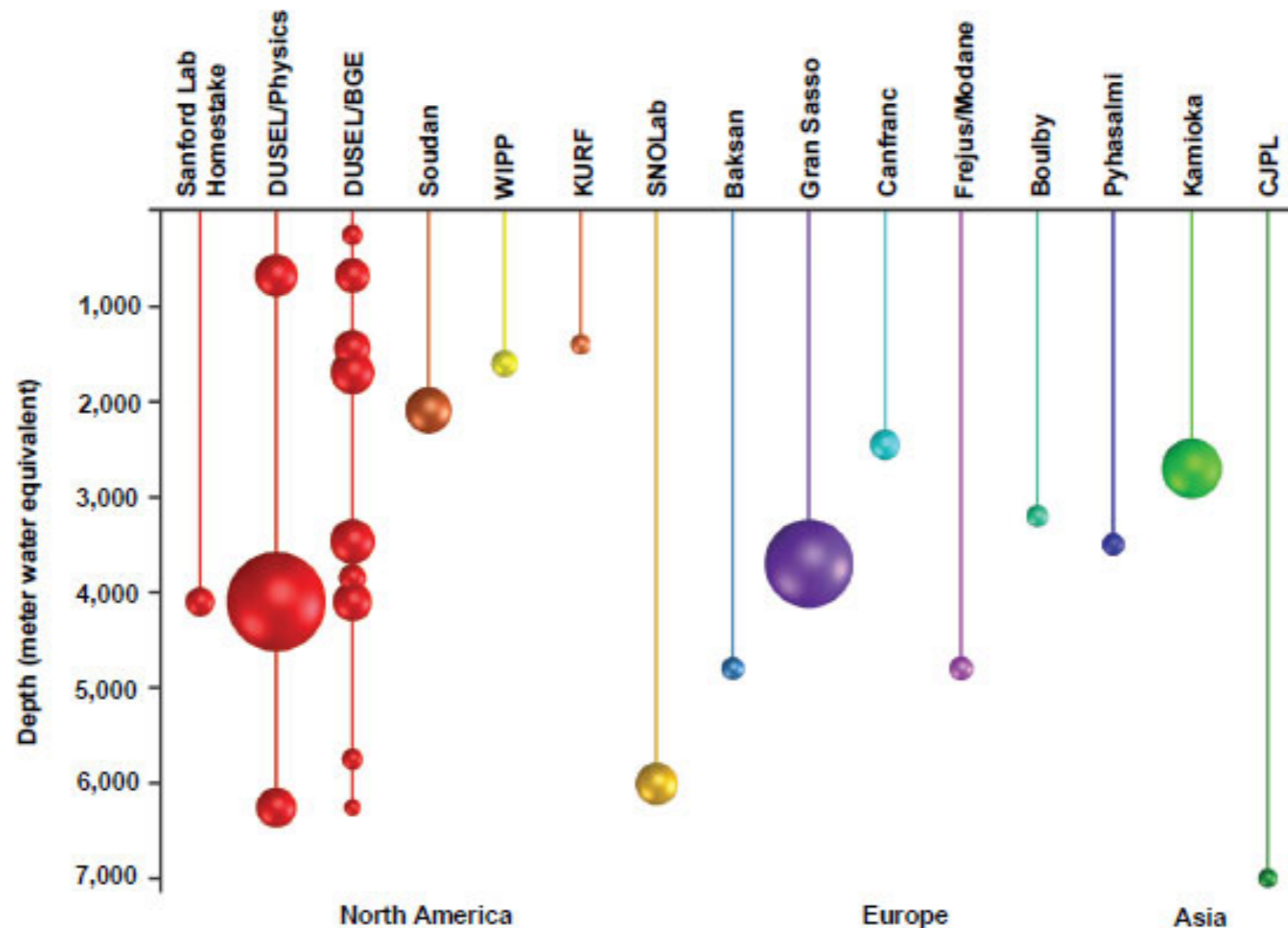


- $M = 1,000 \text{ kg}$
- $t = 3 \text{ yr}$
- $T_{1/2} = 10^{26} \text{ yr}$
- $\varepsilon = 100\%$
- $\Delta E = 1\% \text{ FWHM}$
- $c = 10^{-3} /(\text{keV} \cdot \text{kg} \cdot \text{yr})$

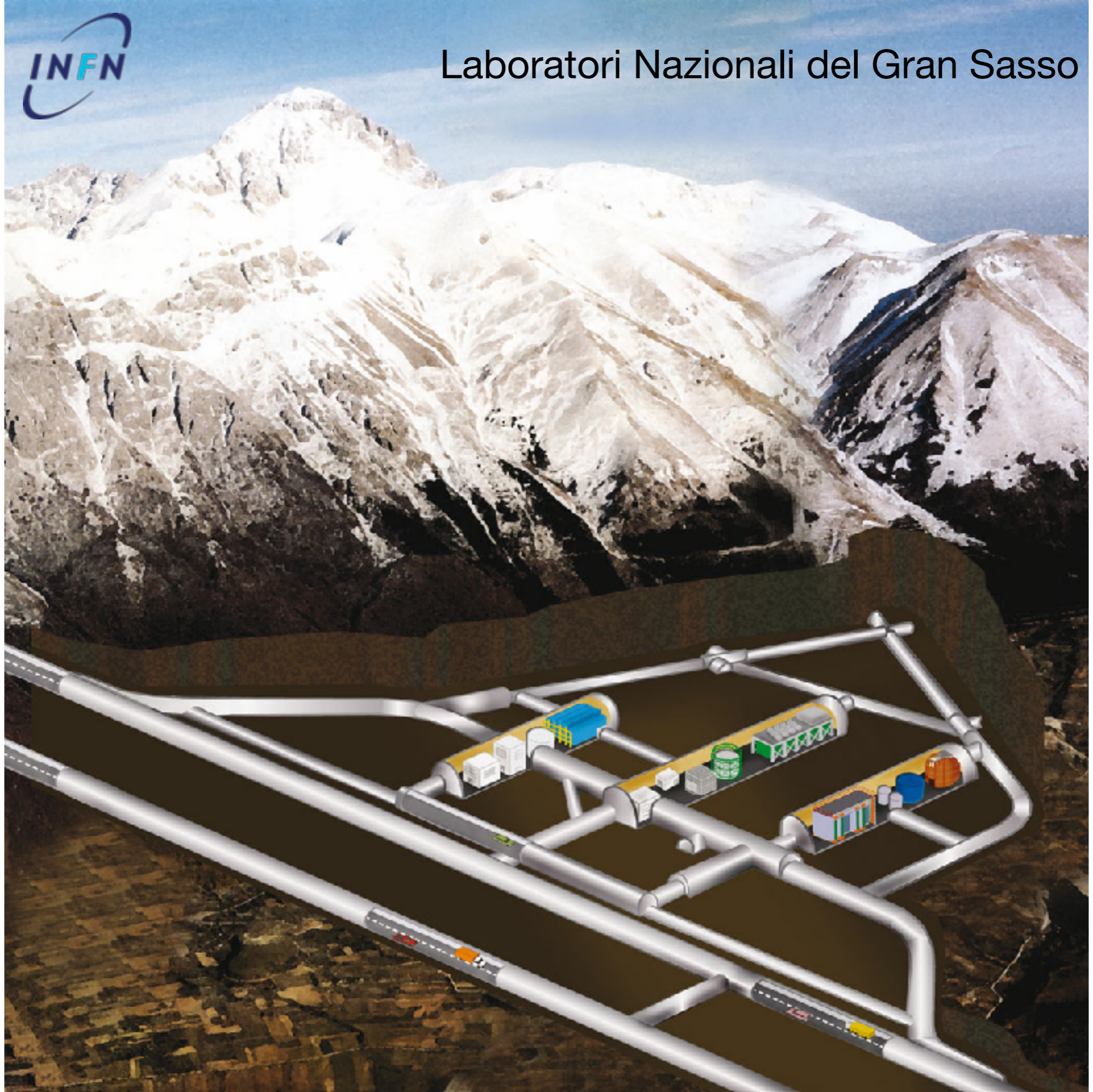
How to lower background rate?

- Go underground and radiopure
- Electron tracking
- Daughter ion tagging?

Underground detectors

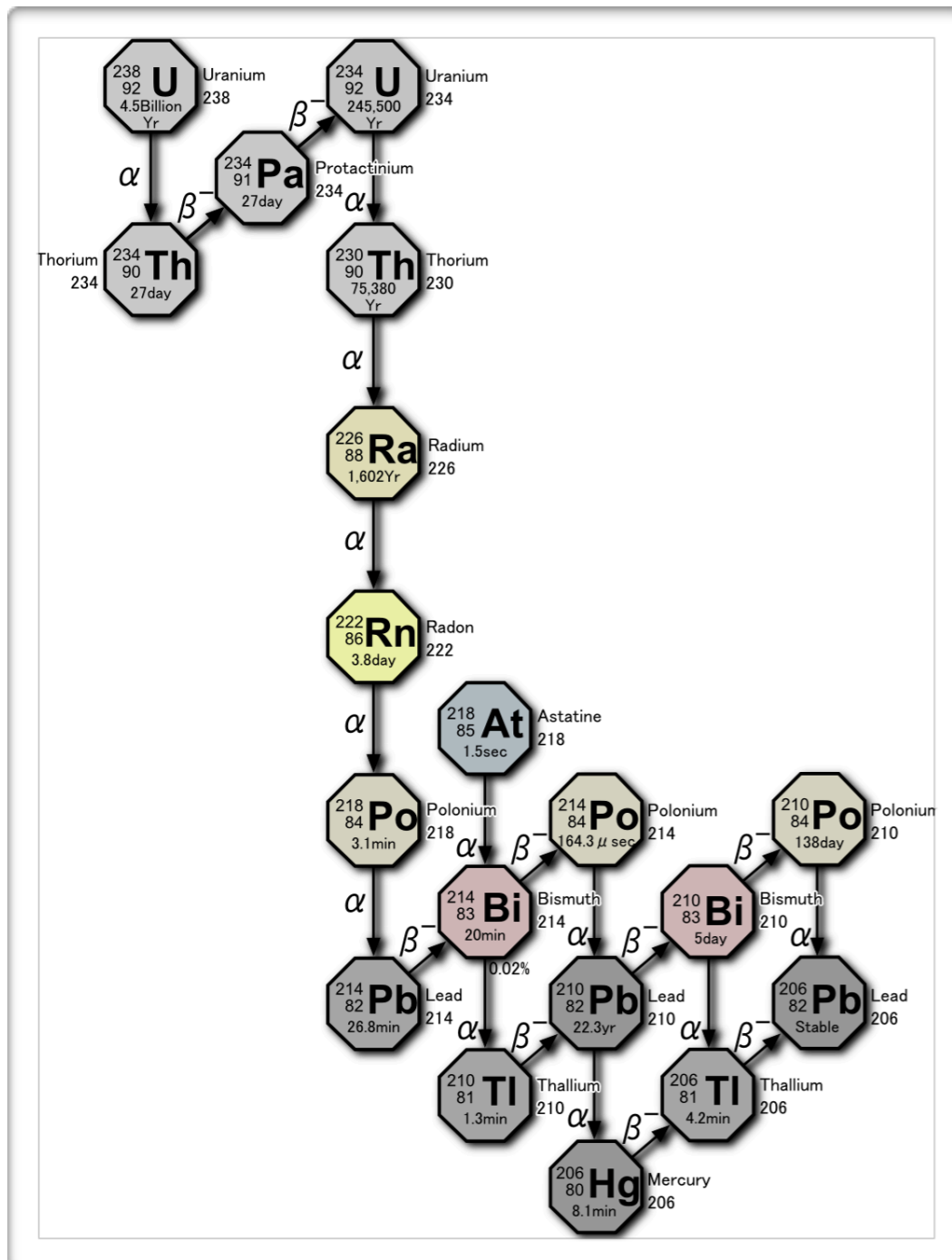


- Some backgrounds originated outside detector by cosmic-ray interactions
- All $\beta\beta 0\nu$ experiments located deep underground, using rock as shield



Radiopure detectors

- Minimise contamination from natural radioactivity in all detector components

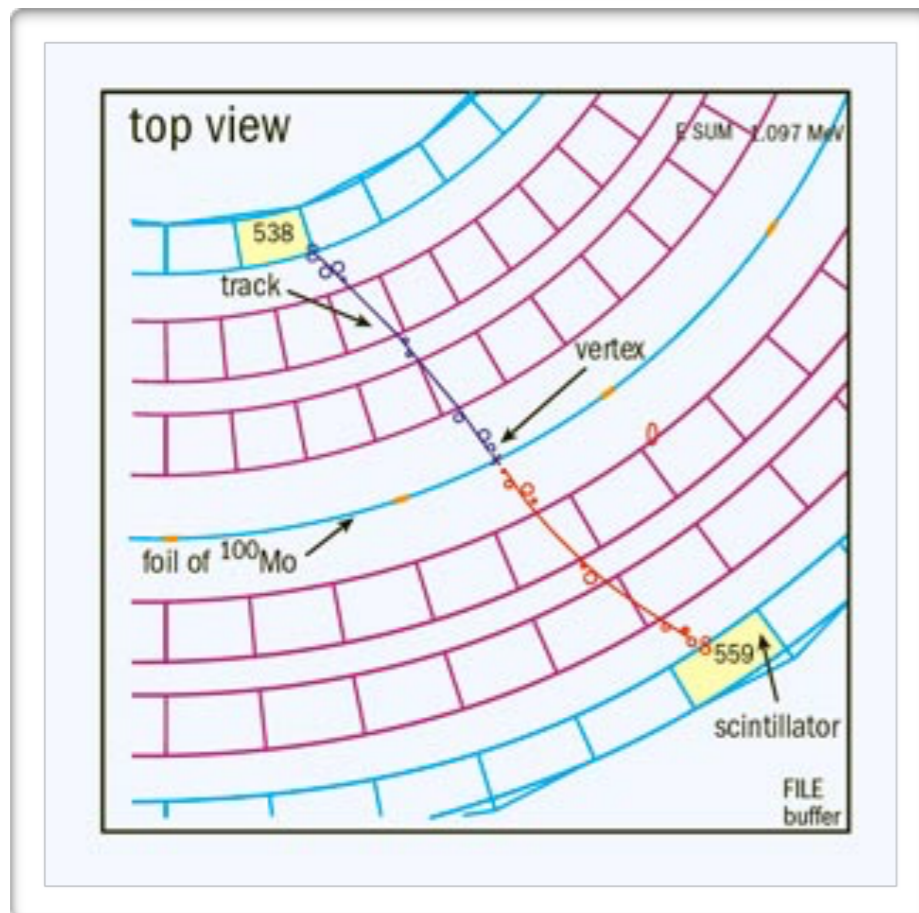


Tracking detectors

- Observe the two stopping electron tracks emitted from common vertex!

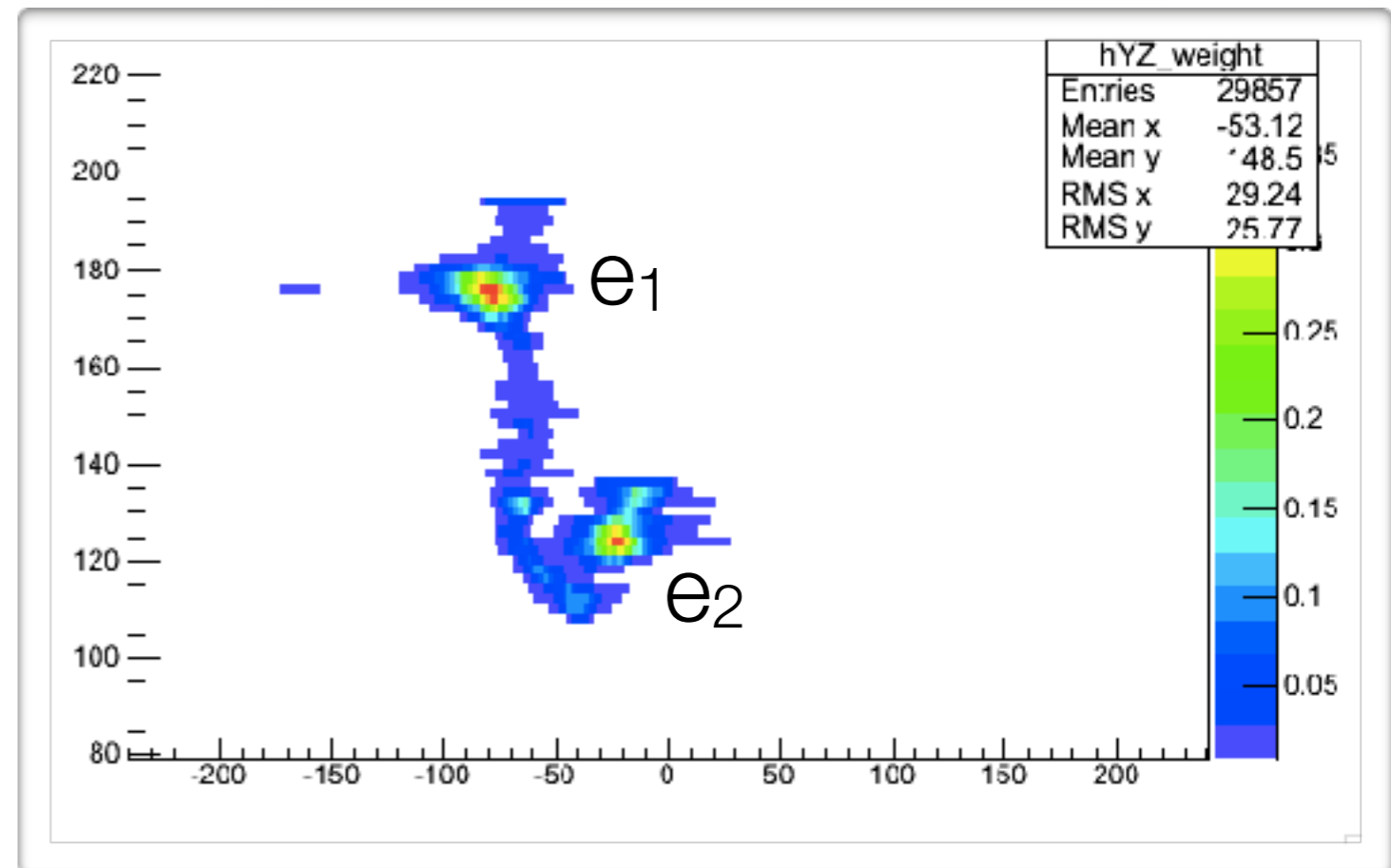
NEMO3 experiment

(drift Geiger cells, He gas at 1 bar)



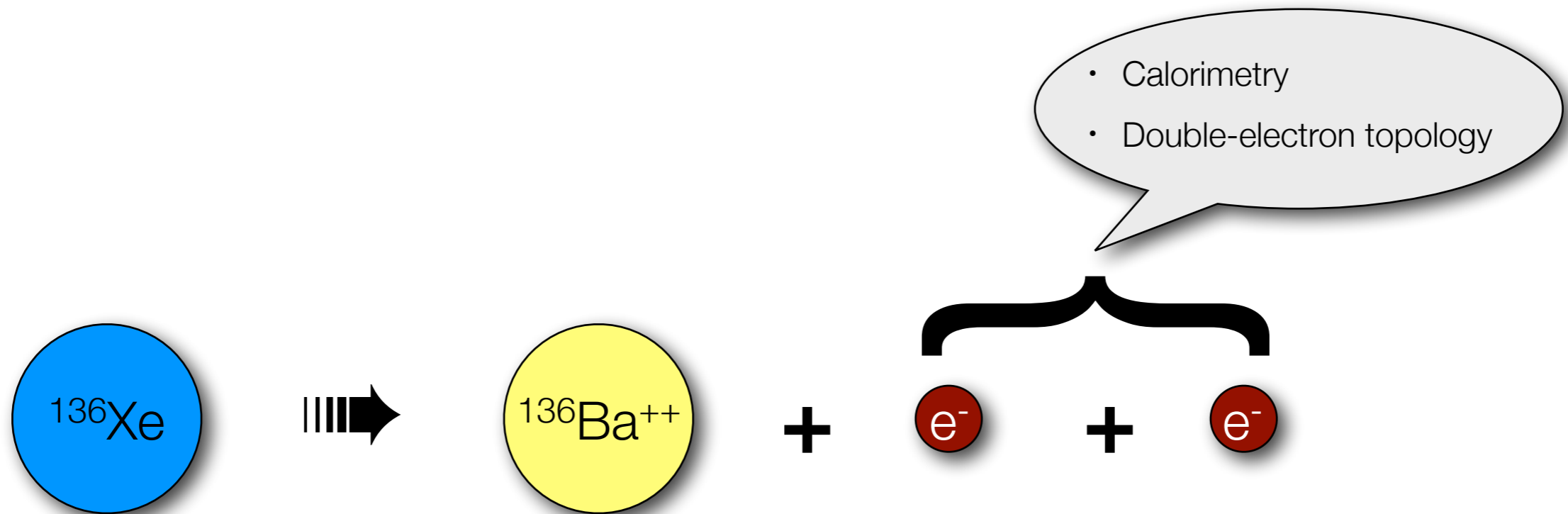
NEXT experiment

(time projection chamber, Xe gas at 7-15 bar)



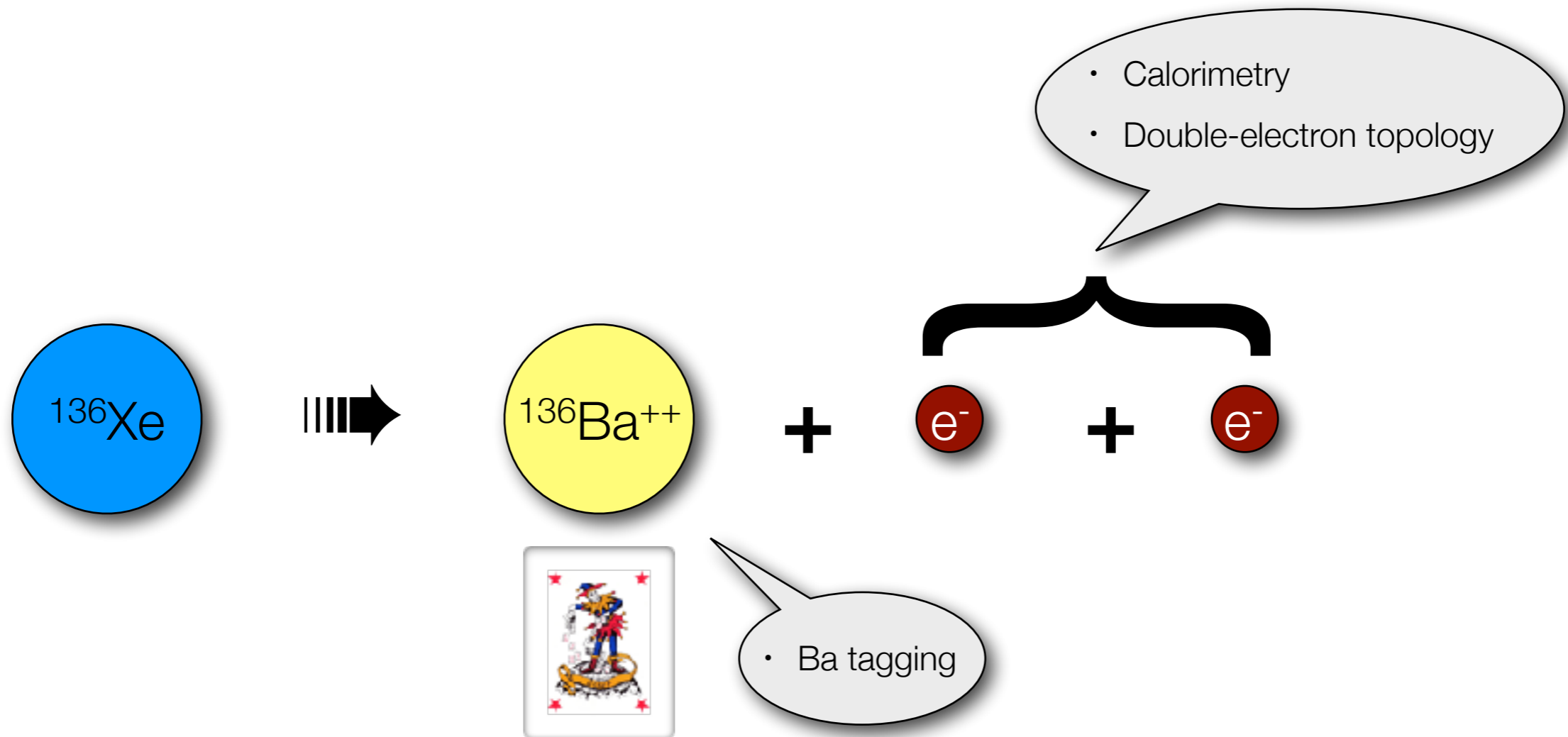
Is daughter ion tagging possible?

- Active R&D in ^{136}Xe experiments (liquid and gas) to detect $^{136}\text{Ba}^{++}$ ion:



Is daughter ion tagging possible?

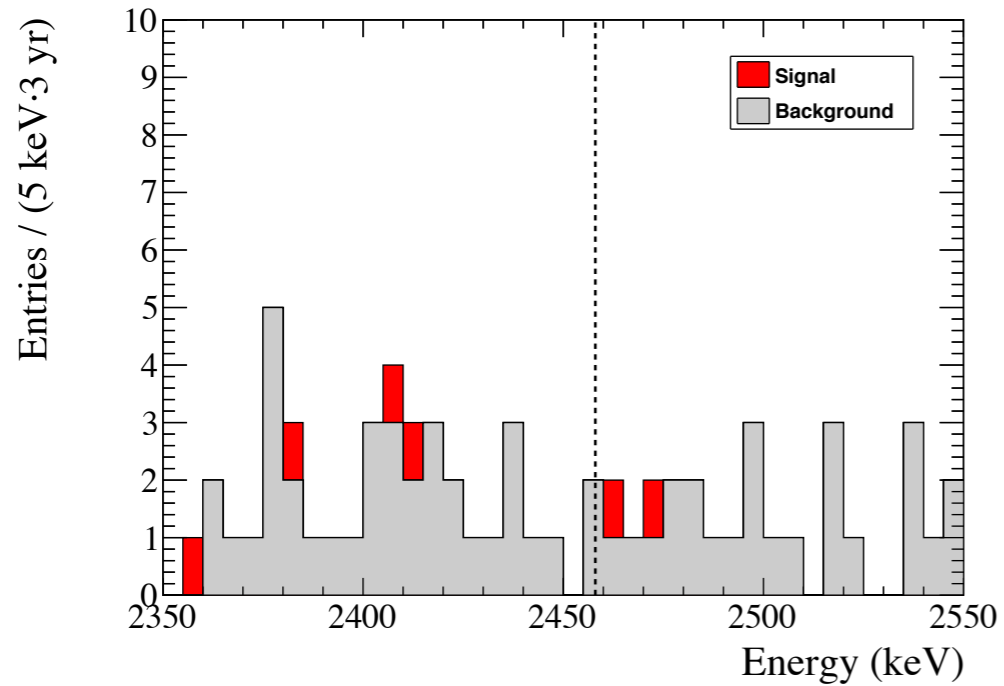
- Active R&D in ^{136}Xe experiments (liquid and gas) to detect $^{136}\text{Ba}^{++}$ ion:



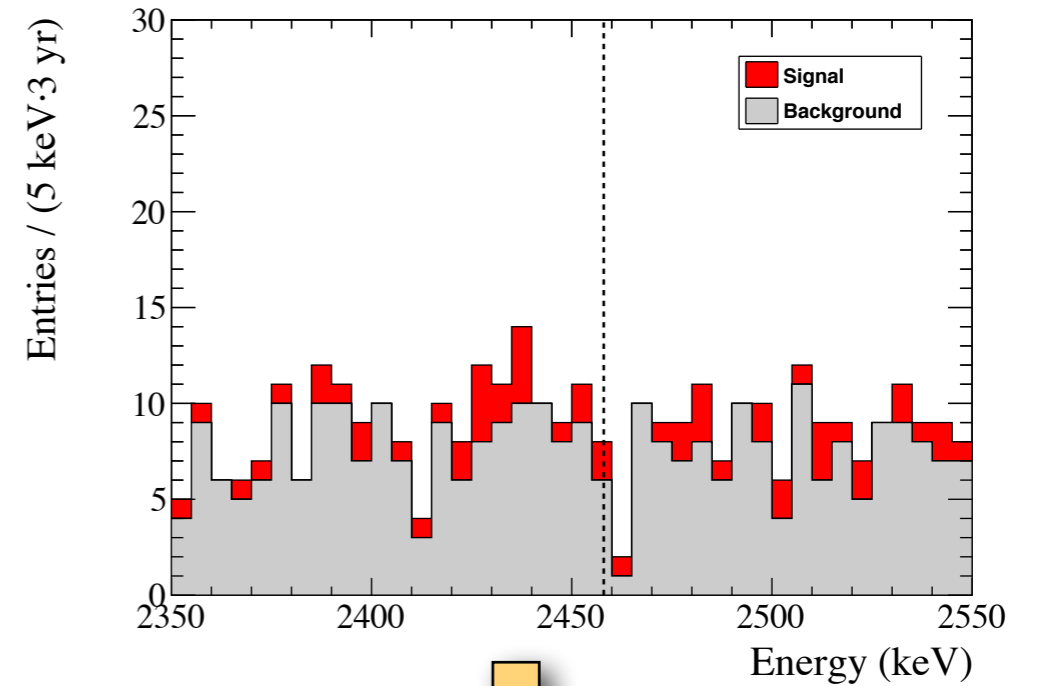
- If successful, one would be left with $\beta\beta 2\nu$ background only!

Putting it all together: Q $_{\beta\beta}$ peak uncovered!

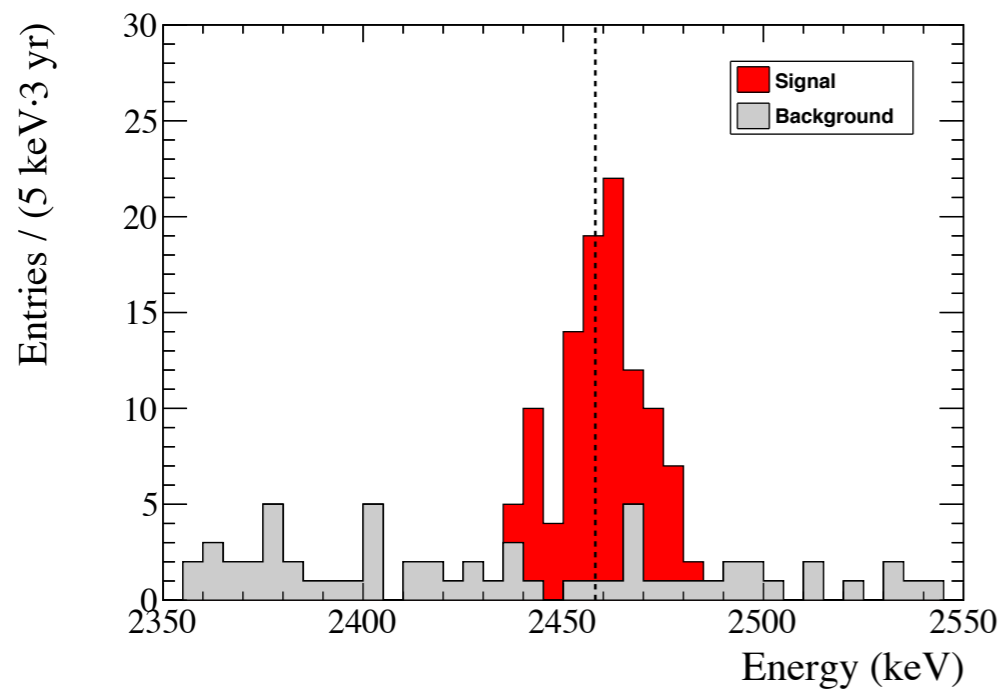
$$T_{1/2}^{0\nu} \propto \epsilon \sqrt{\frac{Mt}{c \Delta E}}$$



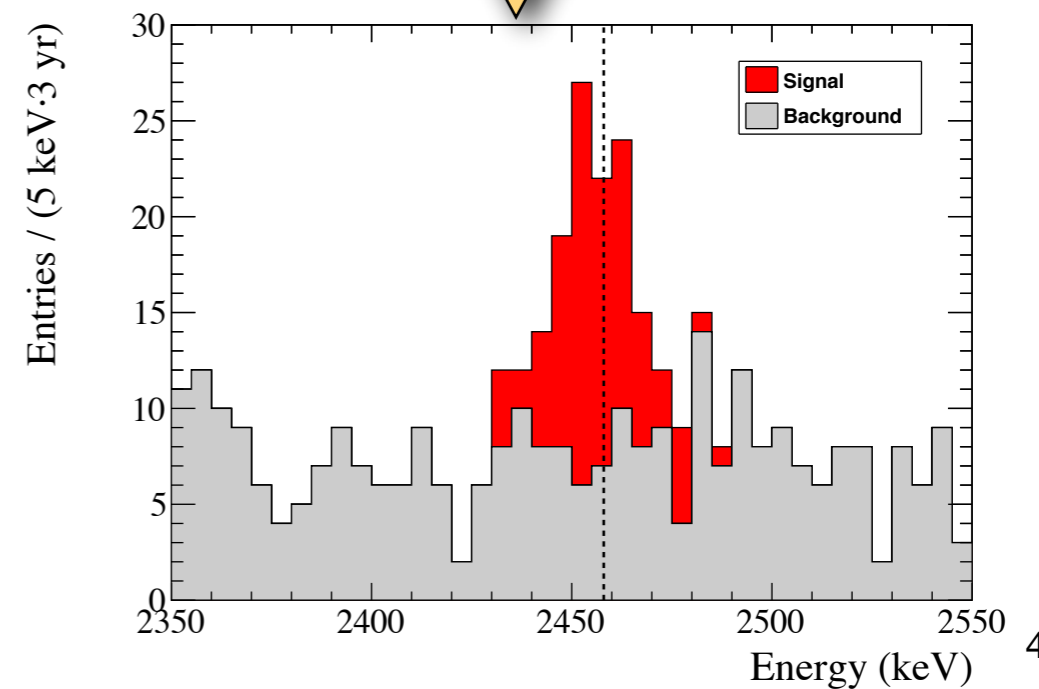
Mt ↗



ΔE ↘



C ↙



Chapter 6:

*T*he present of $\beta\beta 0\nu$ searches

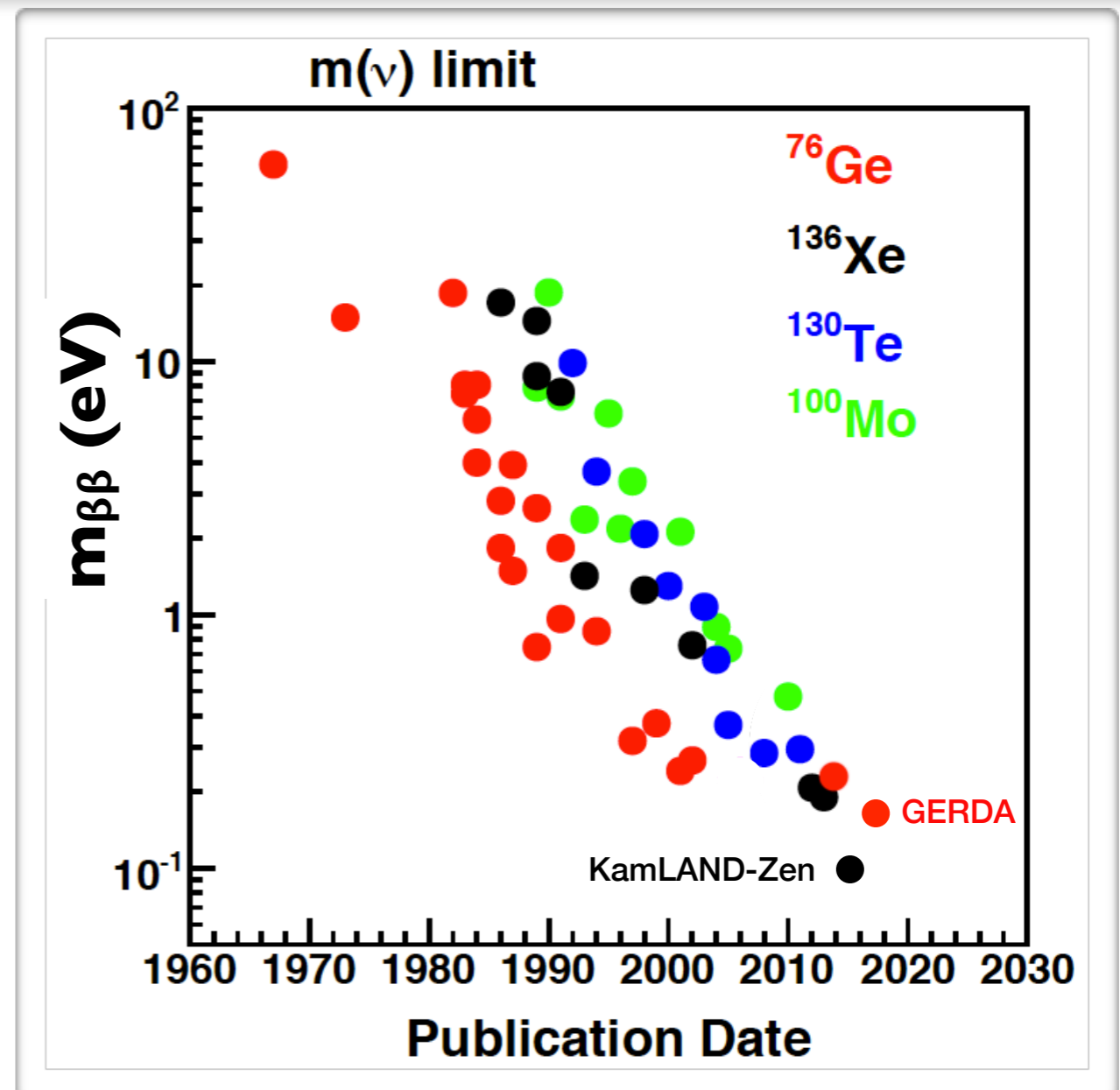
$\beta\beta 0\nu$ experimental status

Main experiments, current generation:



- No convincing evidence for $\beta\beta 0\nu$
- Best limits:

Experiment	$T_{1/2}^{0\nu}$ limit (yr)	$m_{\beta\beta}$ limit (meV)
KamLAND-Zen	$> 1.07 \times 10^{26}$	$< 61-165$
GERDA	$> 5.3 \times 10^{25}$	$< 150-330$





GERDA experiment

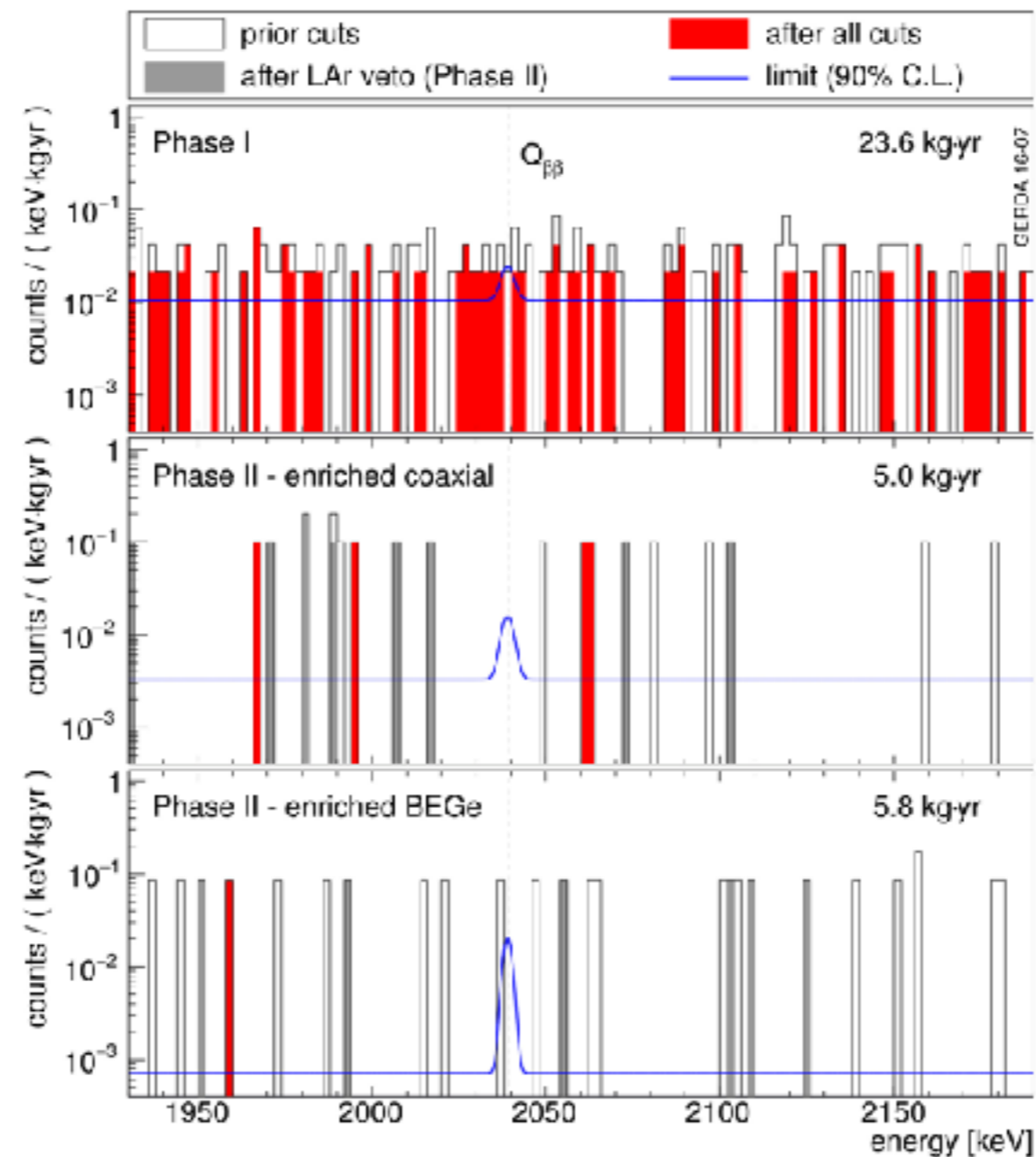
- High-purity germanium diodes enriched in ^{76}Ge immersed in LAr
- **Advantages:** energy resolution, radiopurity → background-free!



GERDA experiment



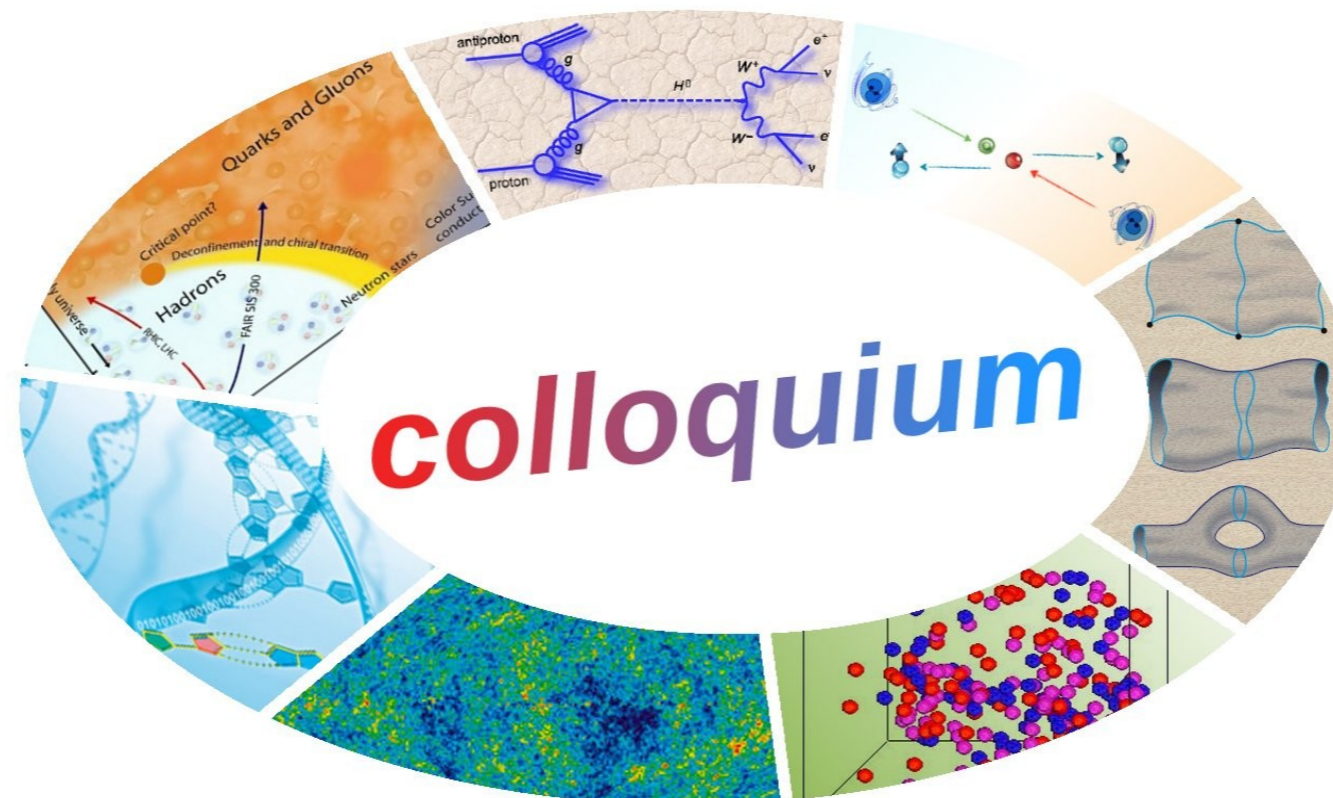
- High-purity germanium diodes enriched in ^{76}Ge immersed in LAr
- **Advantages:** energy resolution, radiopurity \rightarrow background-free!



GERDA experiment



Dipartimento di Fisica e Sezione INFN di Torino

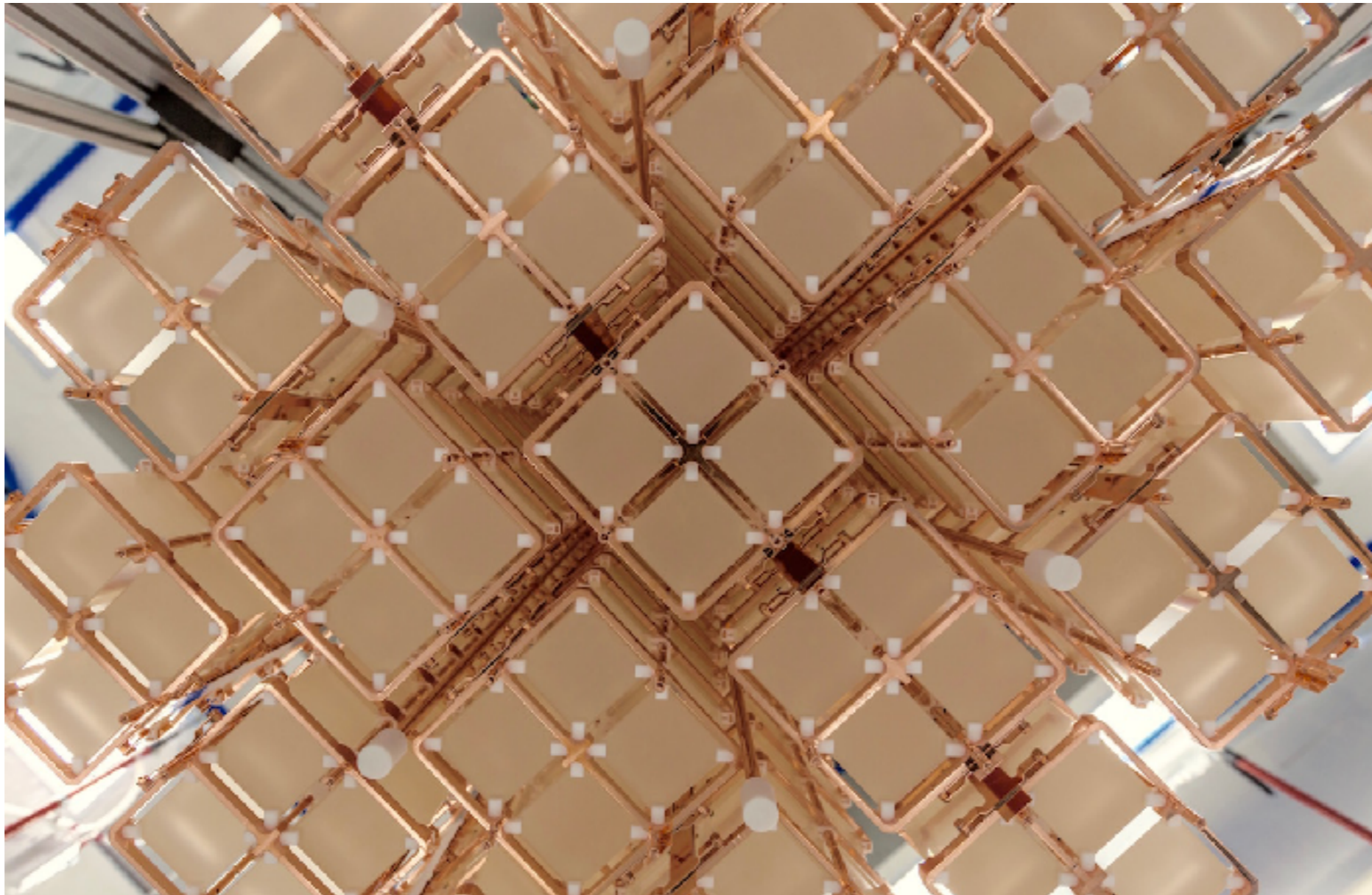


Giovedì 1 Giugno 2017, ore 14:30, **Sala Wataghin**
Riccardo Brugnera
(Università di Padova)

Neutrinoless double-beta decay searches with ^{76}Ge

CUORE experiment

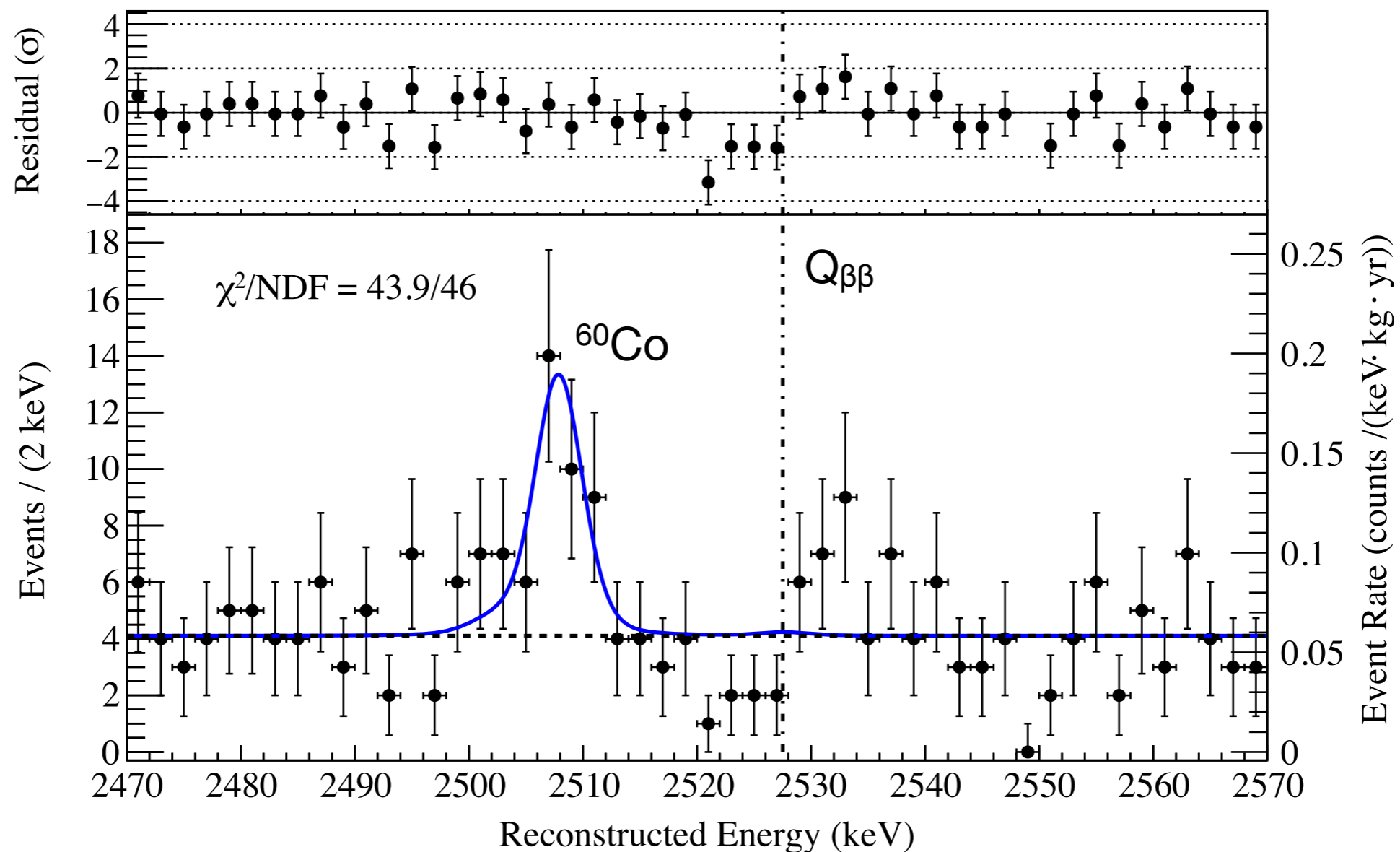
- Towers of TeO_2 crystals. $\beta\beta$ energy measured as temperature increase
- **Advantages:** energy resolution, mass scalability



CUORE experiment

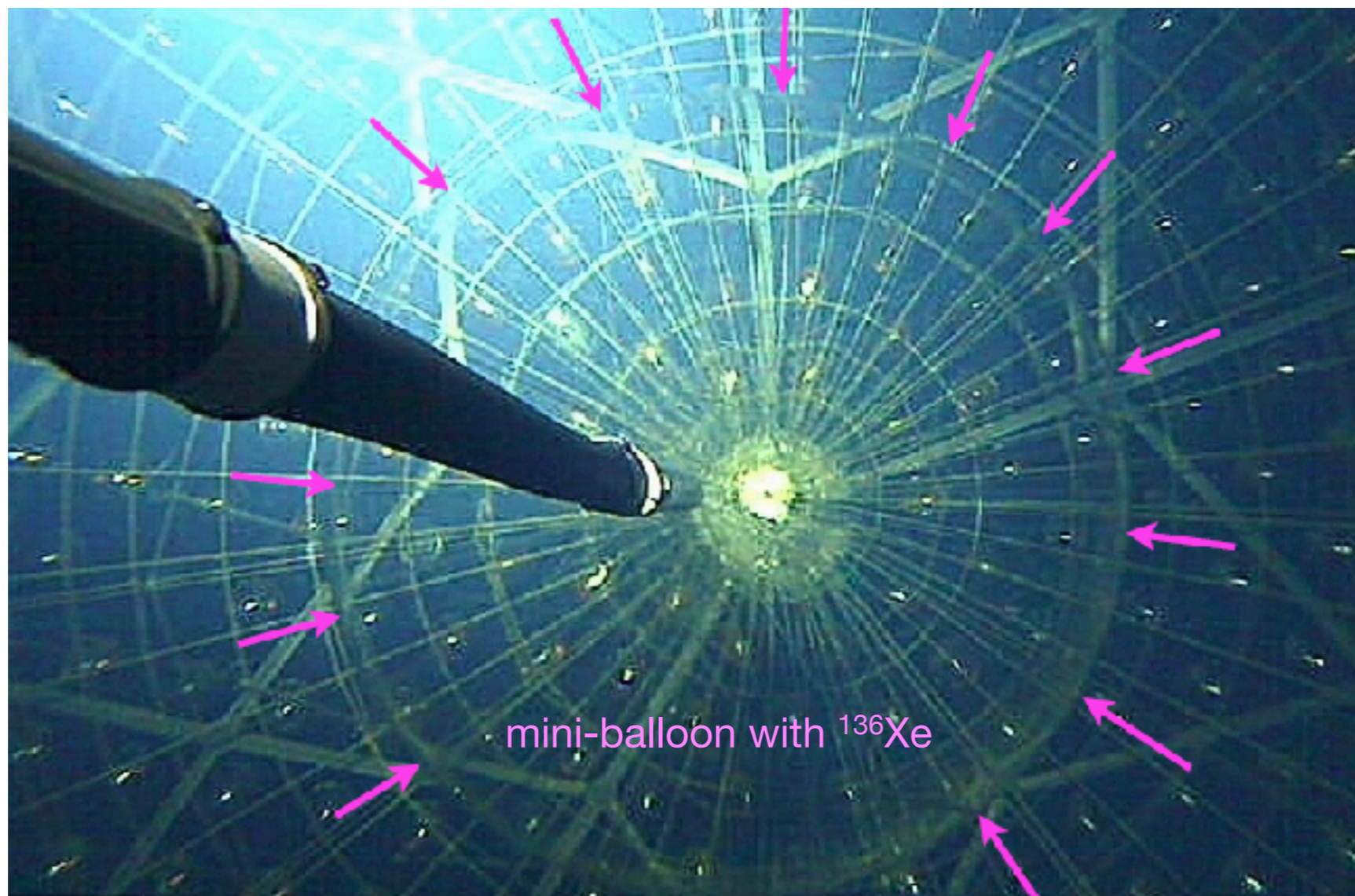


- Towers of TeO_2 crystals. $\beta\beta$ energy measured as temperature increase
- **Advantages:** energy resolution, mass scalability



KamLAND-Zen experiment

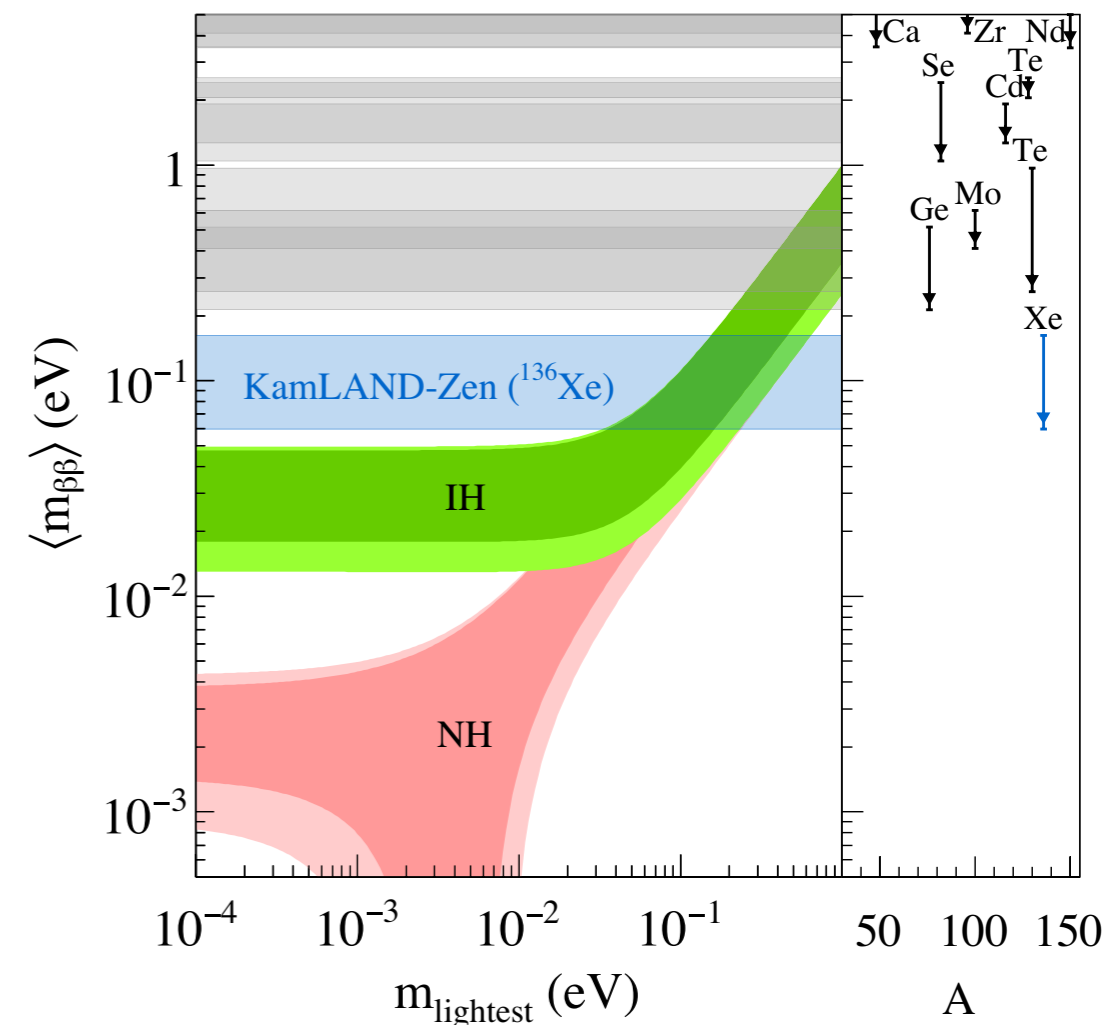
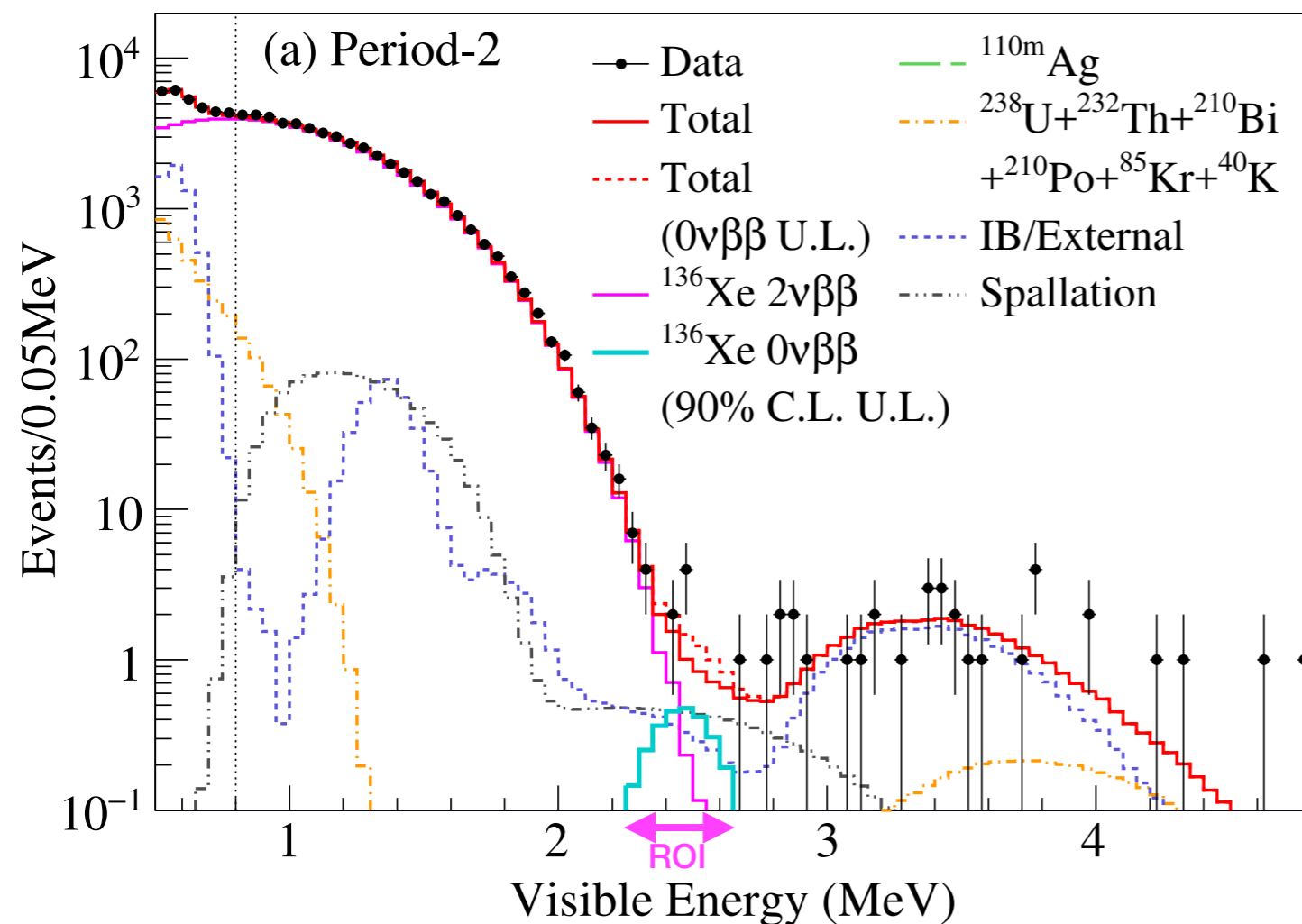
- Liquid scintillator with 300-750 kg of ^{136}Xe gas dissolved in it
- **Advantages:** mass scalability, radiopure, veto region → leading the field





KamLAND-Zen experiment

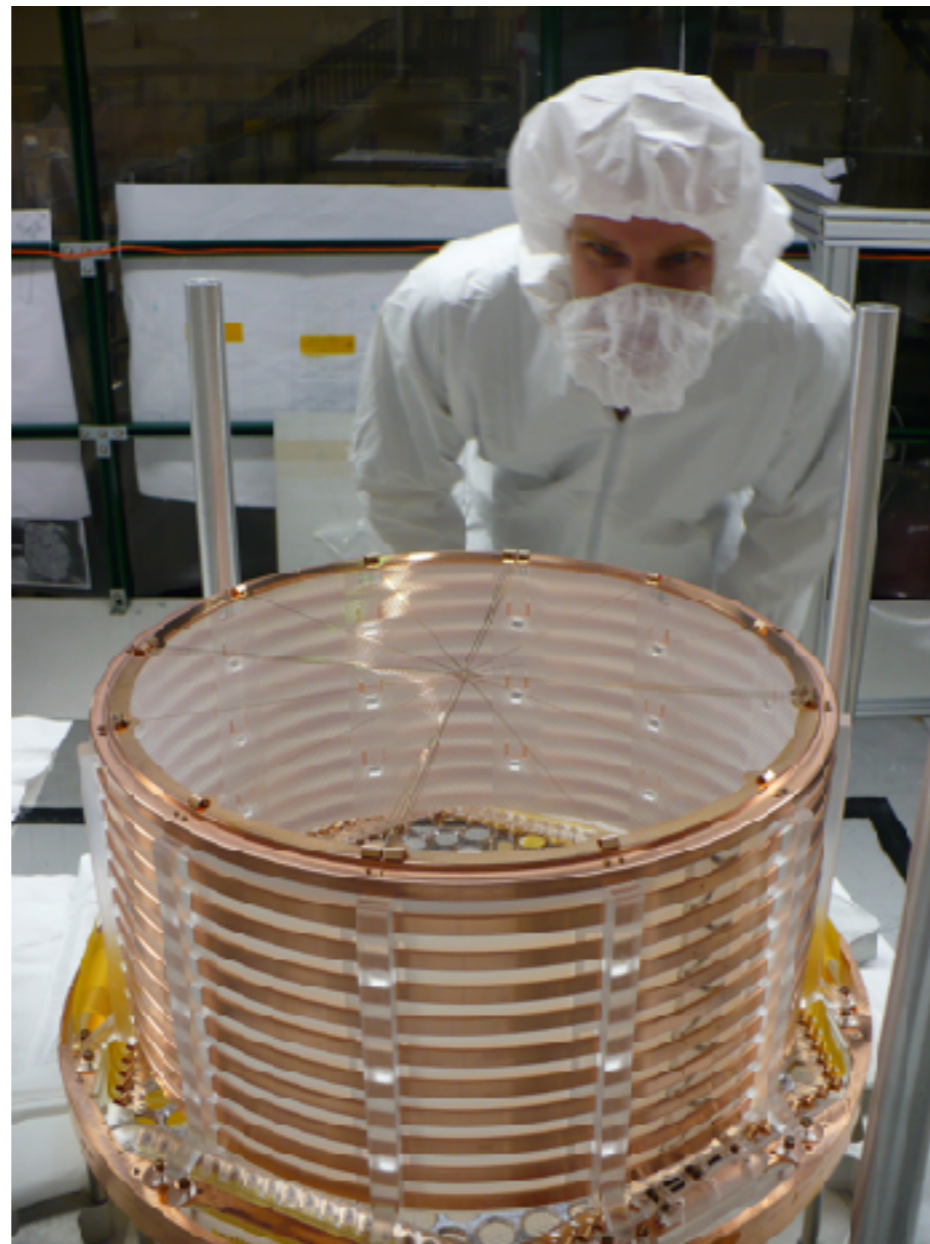
- Liquid scintillator with 300-750 kg of ^{136}Xe gas dissolved in it
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EXO experiment



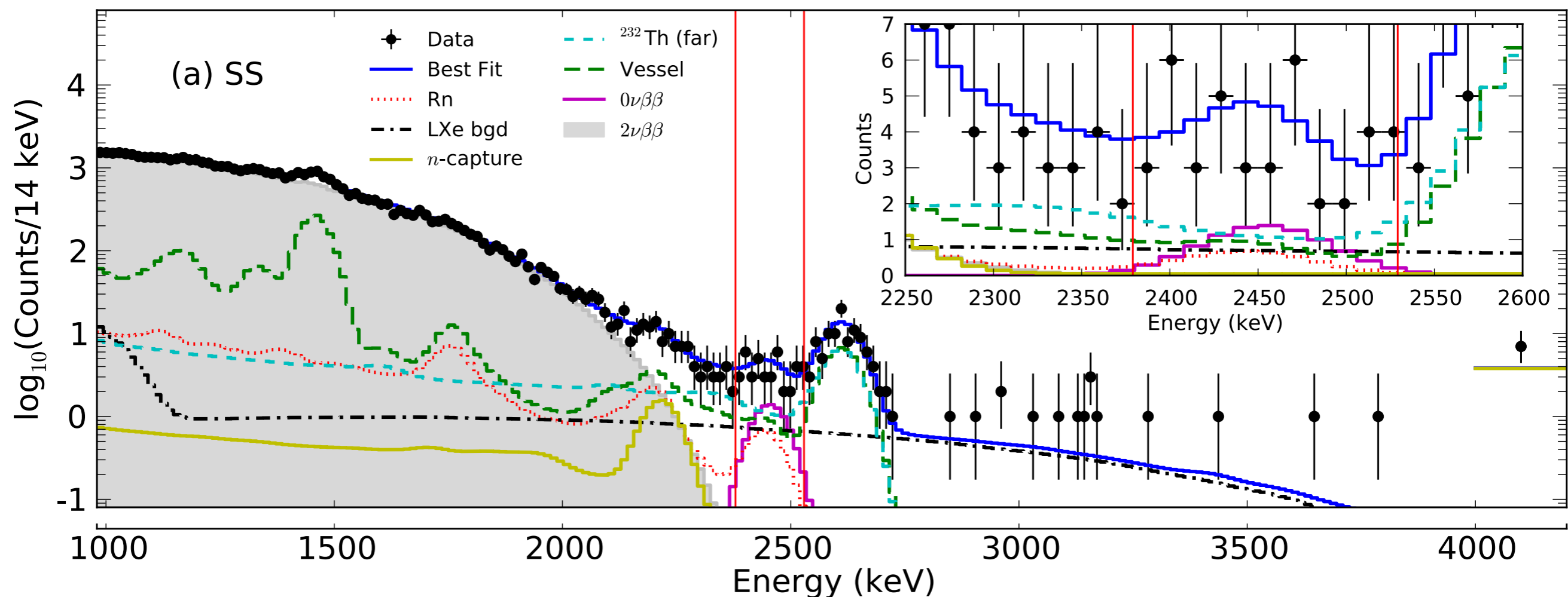
- Cryogenic time projection chamber filled with 80 kg (fiducial) liquid xenon
- **Advantages:** mass scalability, some electron topology



EXO experiment



- Cryogenic time projection chamber filled with 80 kg (fiducial) liquid xenon
- **Advantages:** mass scalability, some electron topology



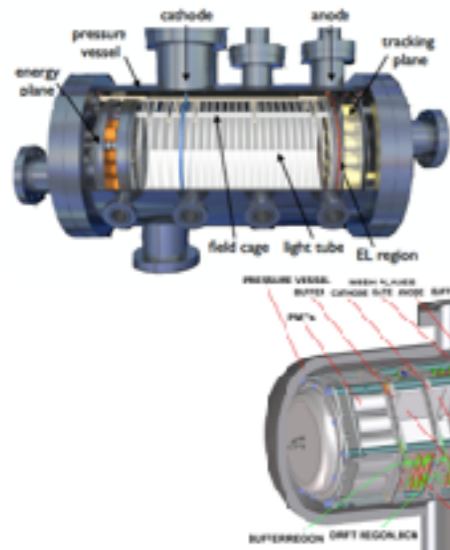
NEXT experiment



- Time projection chamber filled with high-pressure (10-15 bar) ^{136}Xe gas
- **Advantages:** energy resolution, image electron tracks

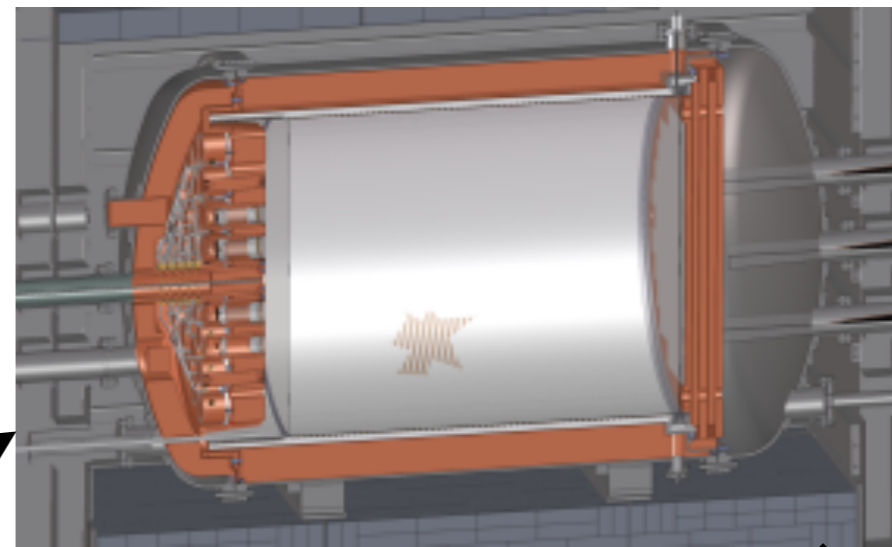


NEXT phases



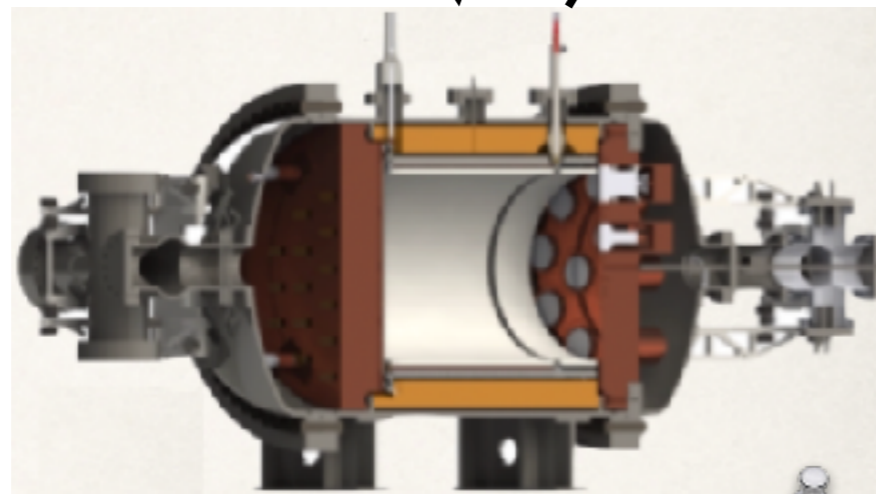
Prototypes (~1 kg)
[2009 - 2014]

Demonstration of detector concept



NEXT-100 (~100 kg)
[2018 - 2020's]

Neutrinoless double beta decay search

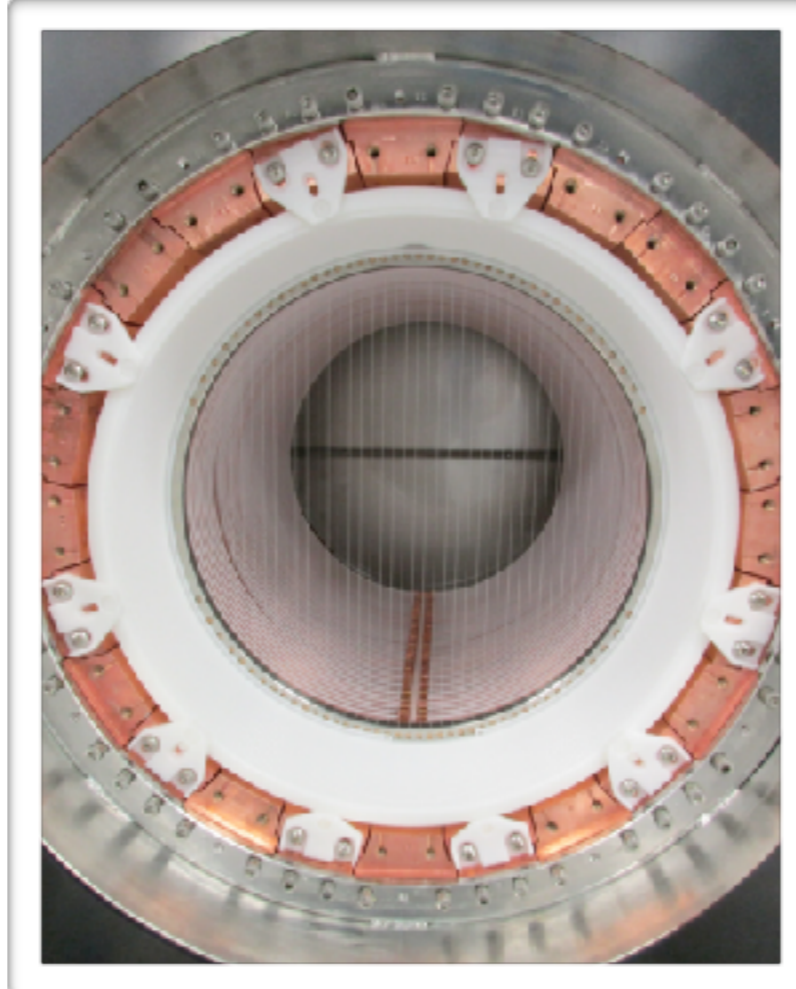
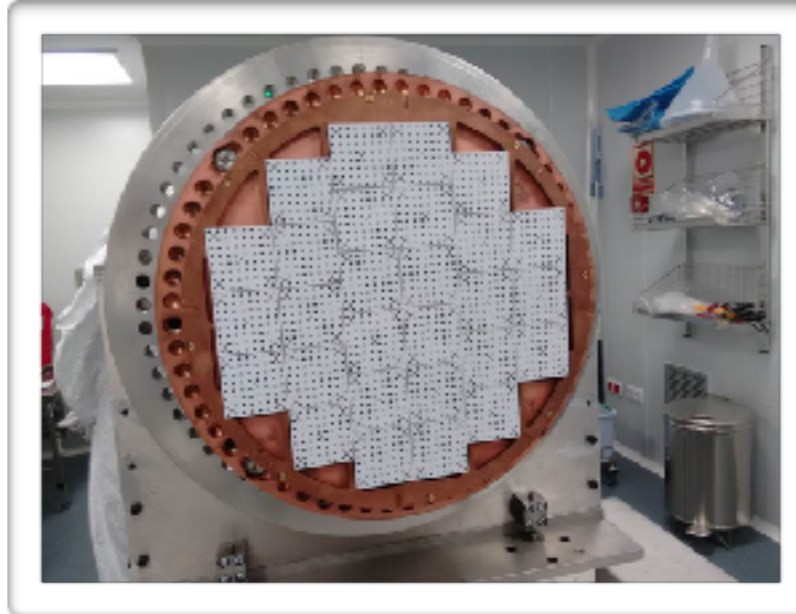


NEXT-NEW (~5 kg)
[2015 - 2018]

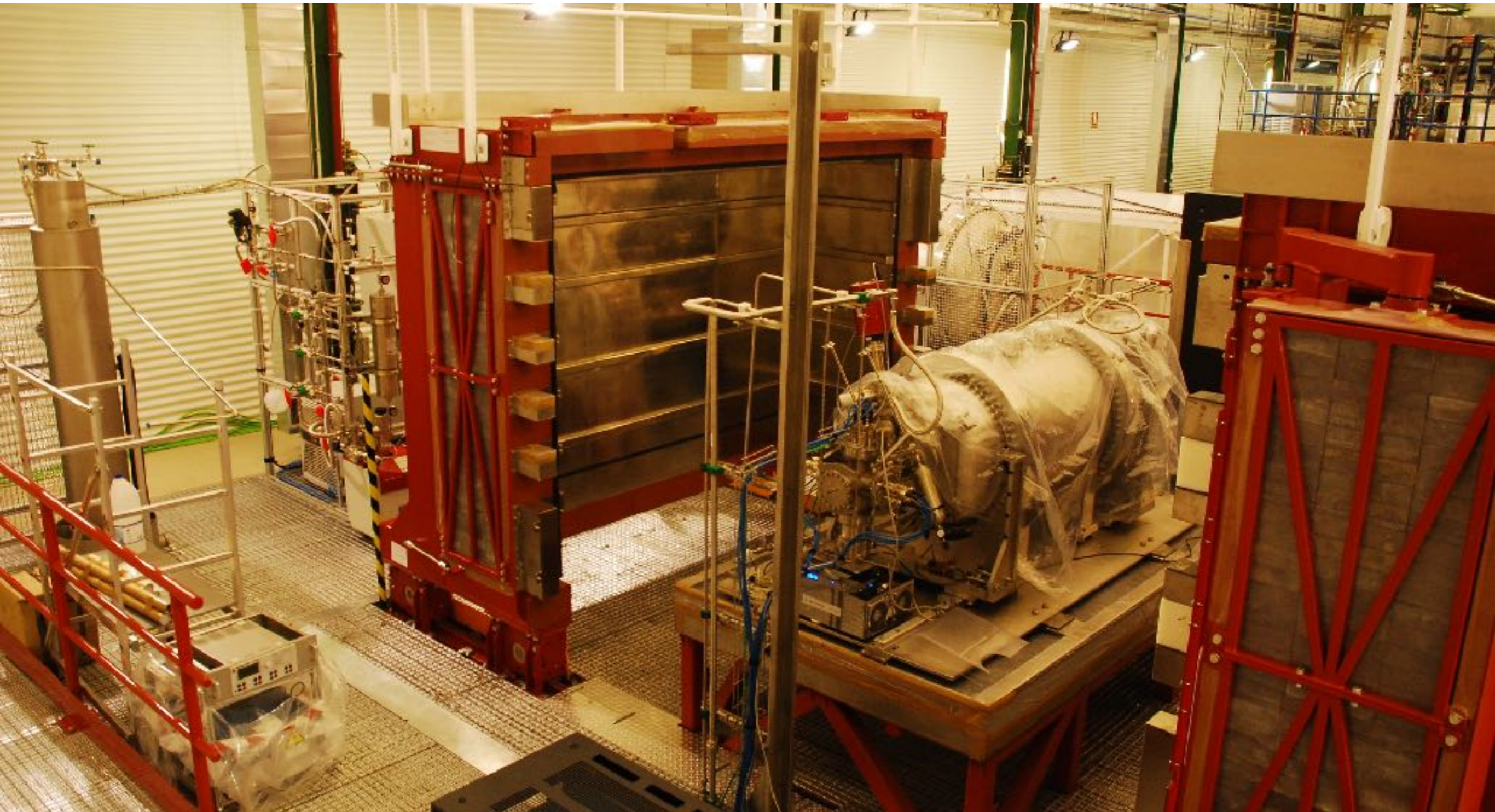
Underground and radio-pure operations, background, $\beta\beta_{2\nu}$

NEXT-tonne (~1000 kg)
[future generation]

NEXT-NEW construction

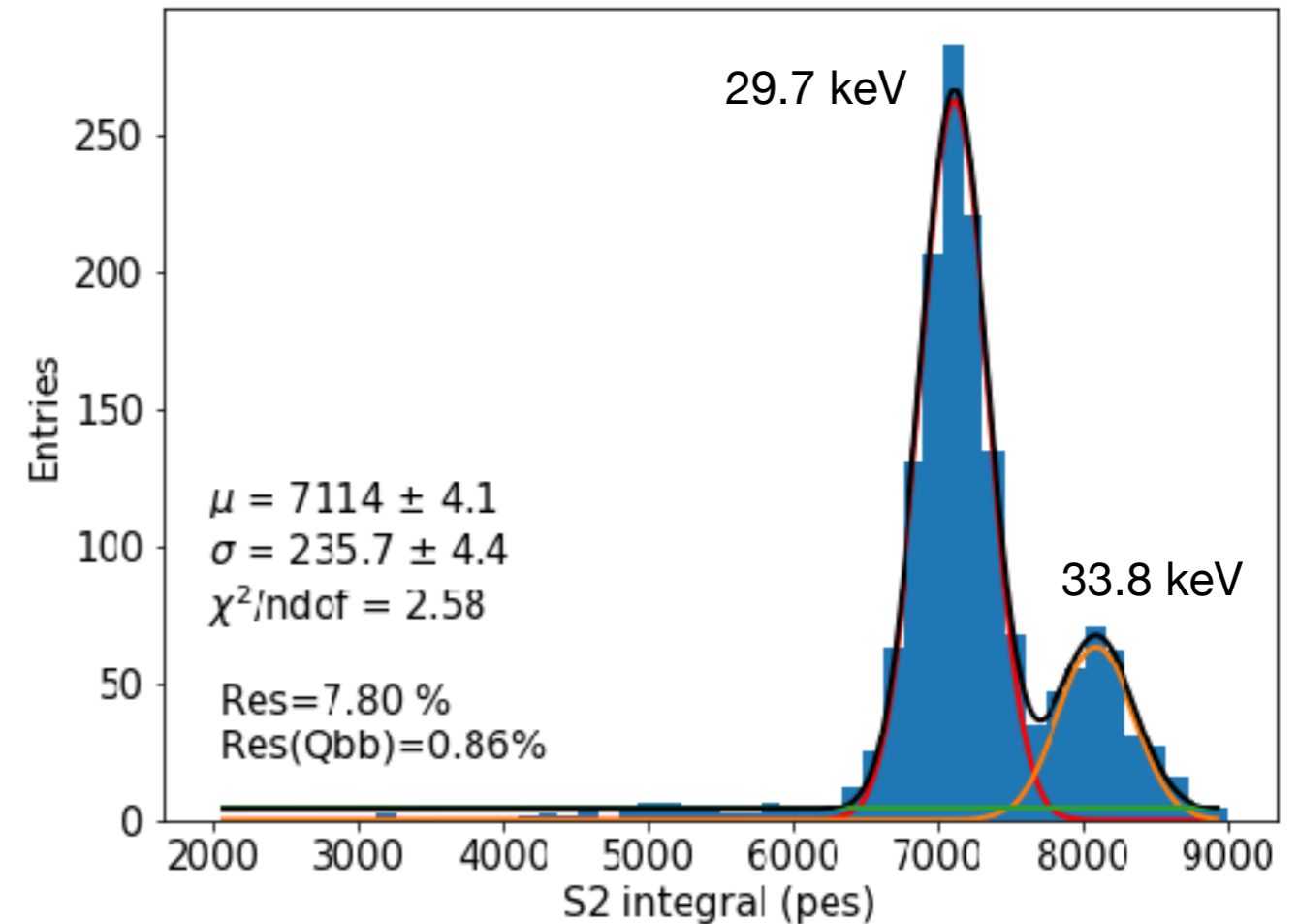
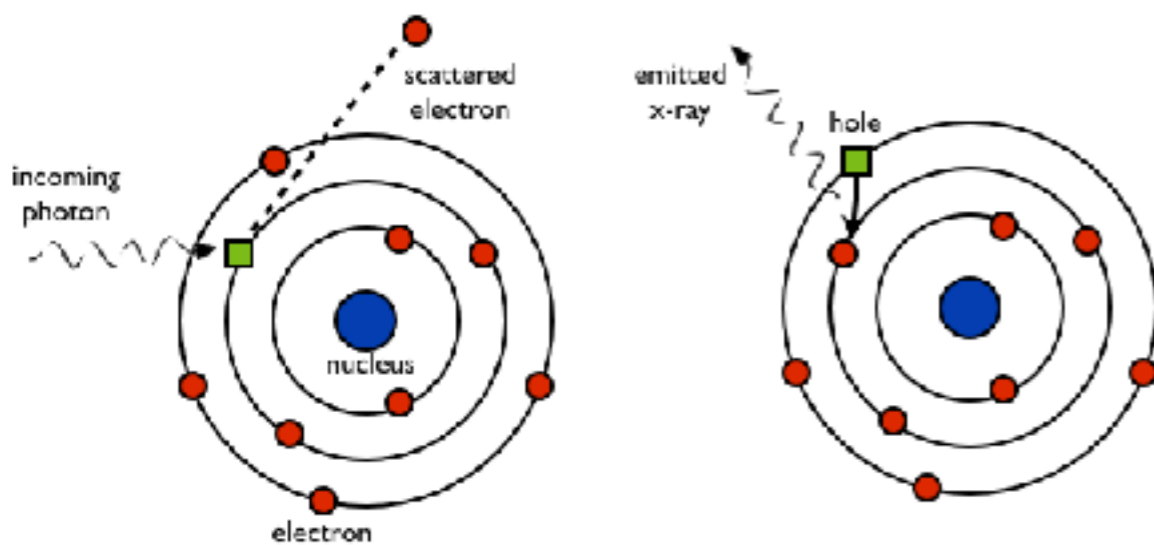


NEXT-NEW installation at the LSC



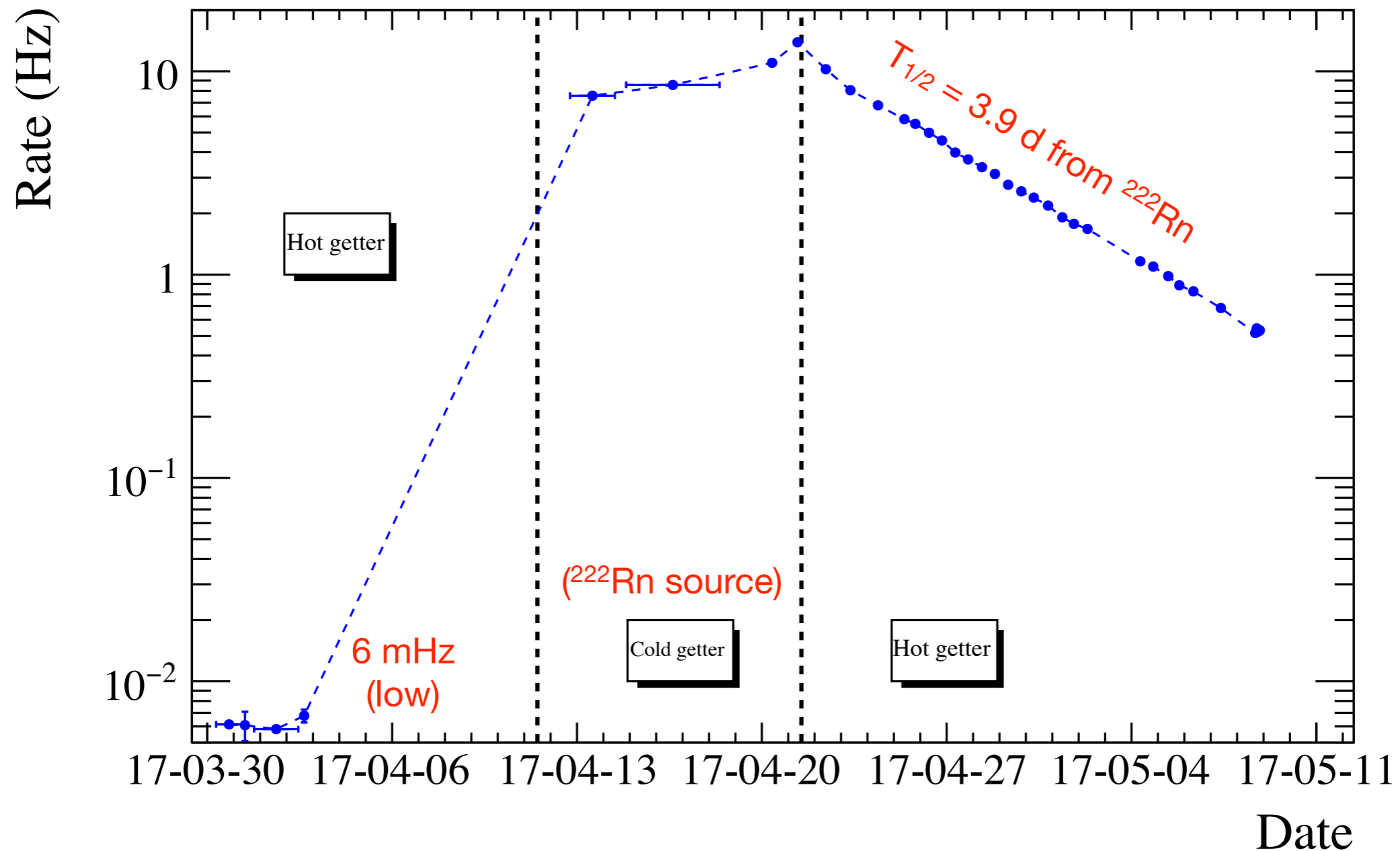
NEXT-NEW first results

- Energy resolution from low-energy xenon X-rays:



NEXT-NEW first results

- Alpha production rate from radon (\rightarrow $\beta\beta 0\nu$ background!):



Chapter 7:

*T*he future of $\beta\beta 0\nu$ searches

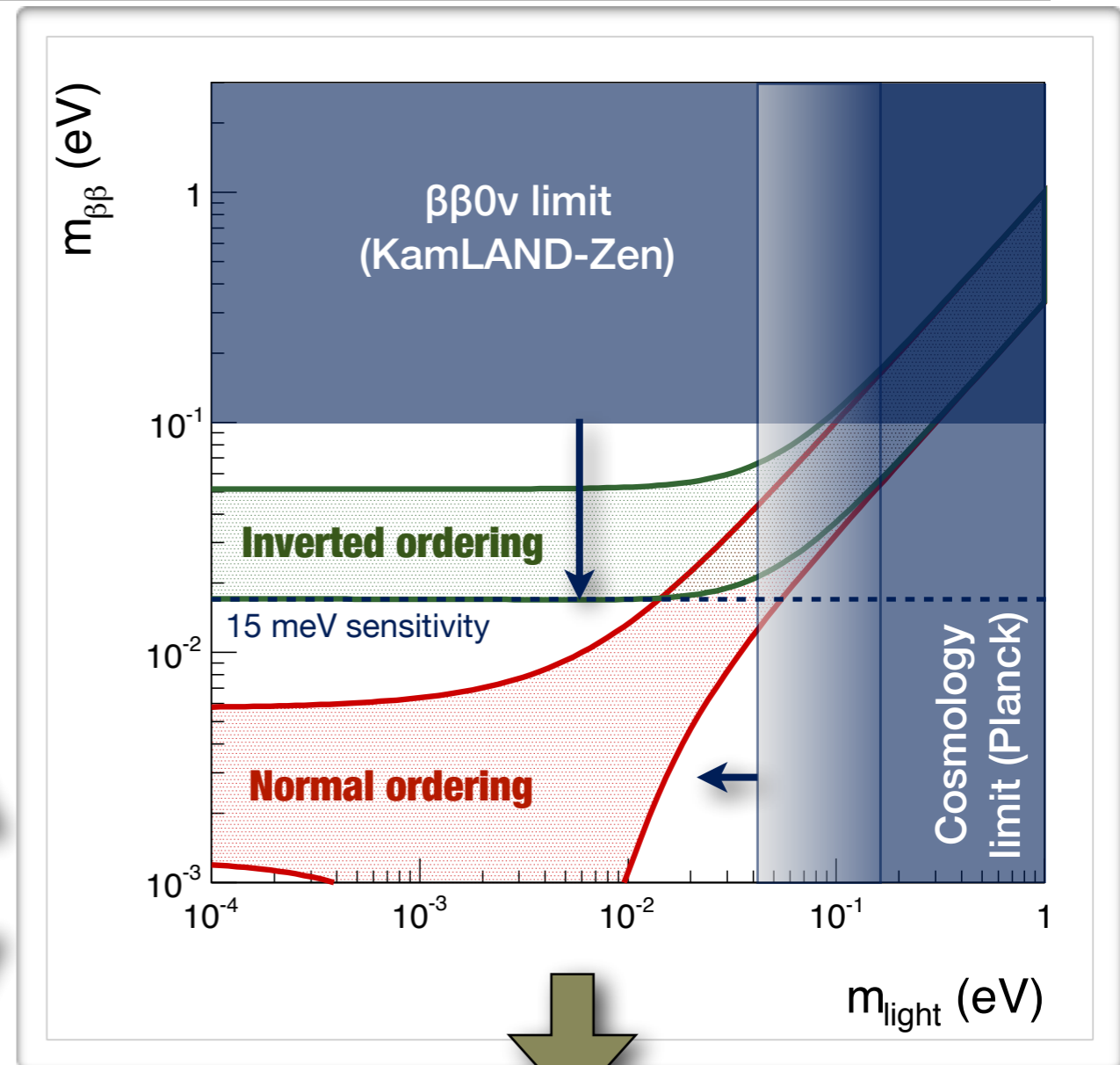
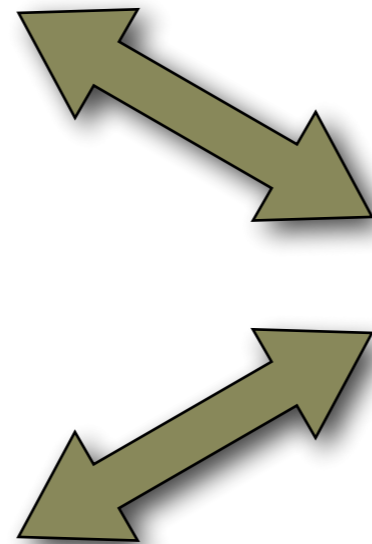
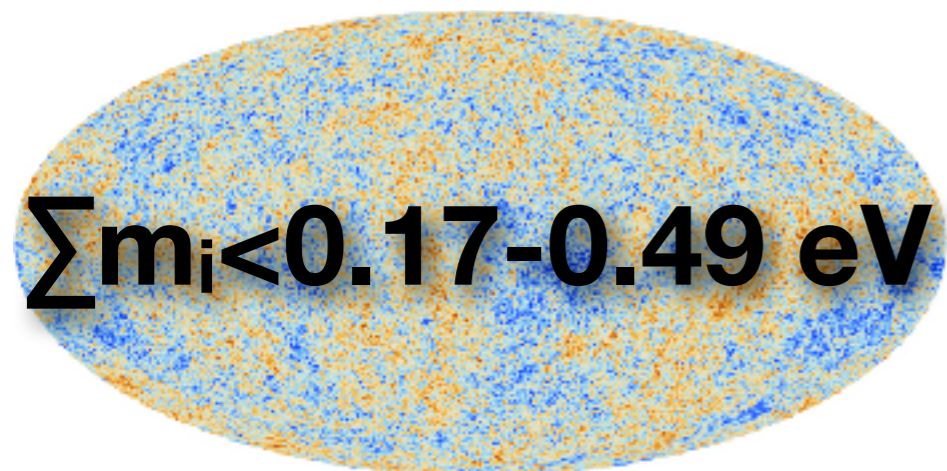
How to move forward?



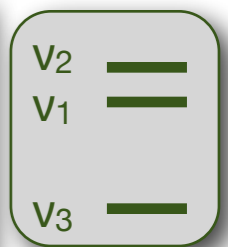
- Strong support to build 2-3 next-generation experiments. Which ones?

Goal for next-generation experiments

15 meV Majorana neutrino mass sensitivity



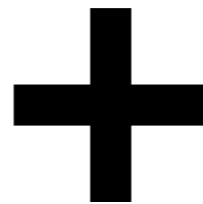
Guaranteed $\beta\beta 0\nu$ discovery if neutrinos are Majorana and have “inverted” mass ordering



Recipe for next-generation experiments



Favour $\beta\beta$ isotopes that can be extrapolated to large masses (\sim ton)



Favour low background rate experimental techniques

The problem of backgrounds

$$T_{1/2}^{0\nu} \propto \varepsilon \sqrt{\frac{Mt}{c \Delta E}} \rightarrow \text{energy resolution}$$

↓
background rate (per unit energy, mass and time)

Current-generation detector performance:

Experiment	c (counts / (keV · kg · yr))	ΔE (keV)	Notes
GERDA	$(0.7-3.5) \times 10^{-3}$	3.5	Measured
CUORE	6×10^{-2}	5	CUORE-0
KamLAND-Zen	1.6×10^{-4}	250	Measured
EXO	1.7×10^{-3}	75	Measured
NEXT	4×10^{-4}	17.5	Prediction

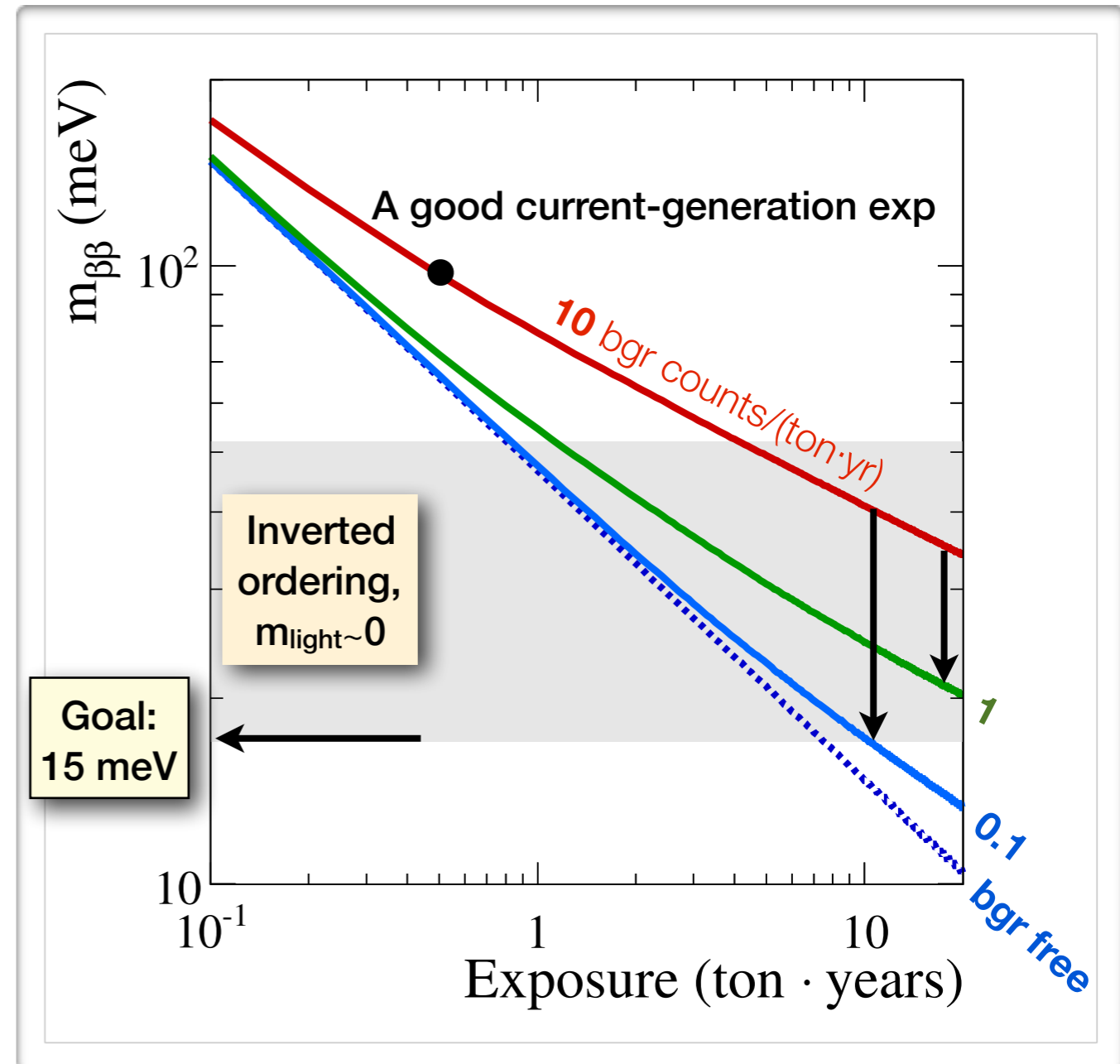
Ton-scale extrapolation:

Background (counts / (ton · yr))
2-12
300
40
130
7

- Today, no technique extrapolates to background-free regime at ton-scale

Ton-scale detector and need for background R&D

- Ton-scale detector **necessary but not sufficient** requirement
- Need at least 1-2 orders of magnitude background reduction with respect to current-generation
- **R&D** on active background reduction techniques



Background reduction R&D



Technology	R&D	Current experiments	Ton-scale proposal
Ge detectors	Larger Ge detectors, improved LAr scint. detection	GERDA, MAJORANA	LEGEND (200 kg)
Bolometers	Scintillating bolometers, isotopic enrichment	CUORE	CUPID
Liquid scintillators	High yield LS, light concentrators, high QE PMTs, enrichment	KamLAND-Zen, SNO+	KamLAND2-Zen
LXe-TPCs	Xe scint. readout with SiPMs, cold electronics, Ba tagging	EXO-200	nEXO
HPXe-TPCs	Low diffusion gas mixtures, finer tracking readout, Ba tagging	NEXT-NEW	NEXT-ton

Appendices:

*T*hinking outside the box

The standard story I just told you

Small neutrino mass

Baryogenesis

Majorana neutrinos and
lepton number violation

Neutrinoless double beta decay search

Possible variations

Small neutrino mass

Baryogenesis

No connection?

No connection?

Majorana neutrinos and
lepton number violation

$\beta\beta 0\nu$ non-optimal probe?

Neutrinoless double beta decay search

Other lepton number violating processes?

- $|\Delta L| = 2$ process mediated by: $W^- W^- \rightarrow l_\alpha^- l_\beta^-$, $\alpha, \beta = e, \mu, \tau$

→ can probe different neutrino mass matrix elements $m_{\alpha\beta}$

Other lepton number violating processes?

- $|\Delta L| = 2$ process mediated by: $W^- W^- \rightarrow l_\alpha^- l_\beta^-$, $\alpha, \beta = e, \mu, \tau$

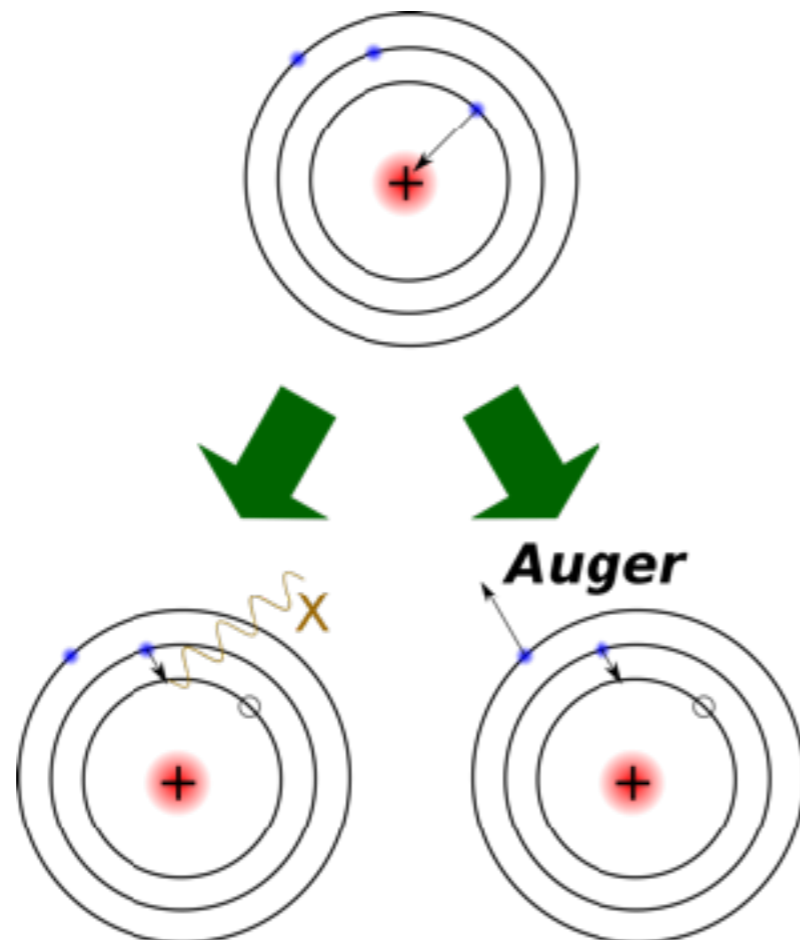
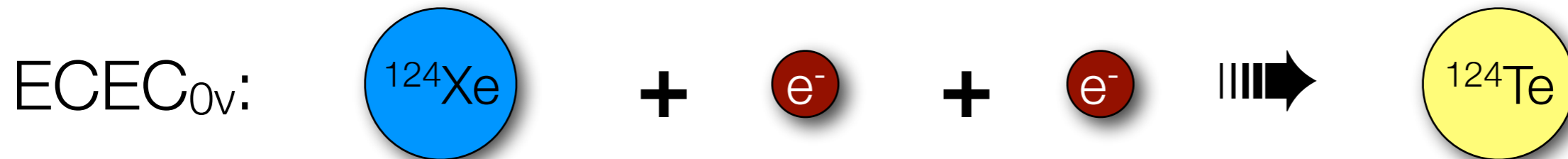
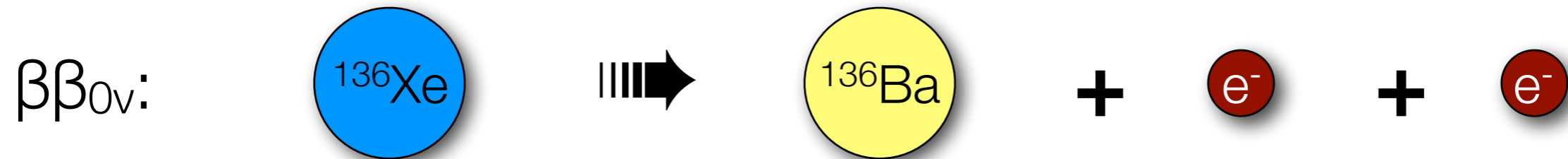
→ can probe different neutrino mass matrix elements $m_{\alpha\beta}$

Flavors	Exp. technique	Mass bound (eV)
(e,e)	$\beta\beta 0\nu$	$m_{ee} < 1 \times 10^{-1}$
(e, μ)	$\mu^- \rightarrow e^+$ conversion	$m_{e\mu} < 2 \times 10^7$
(e, τ)	Rare τ^- decays	$m_{e\tau} < 3 \times 10^{12}$
(μ , μ)	Rare K^+ decays	$m_{\mu\mu} < 3 \times 10^8$
(μ , τ)	Rare τ^- decays	$m_{\mu\tau} < 2 \times 10^{12}$
(τ , τ)	None	None

→ best constraint by far on **(e,e)** element and from **$\beta\beta 0\nu$** , currently

Example: neutrinoless double **E**lectron **C**apture

Inverse of neutrinoless double beta decay!



Observables:

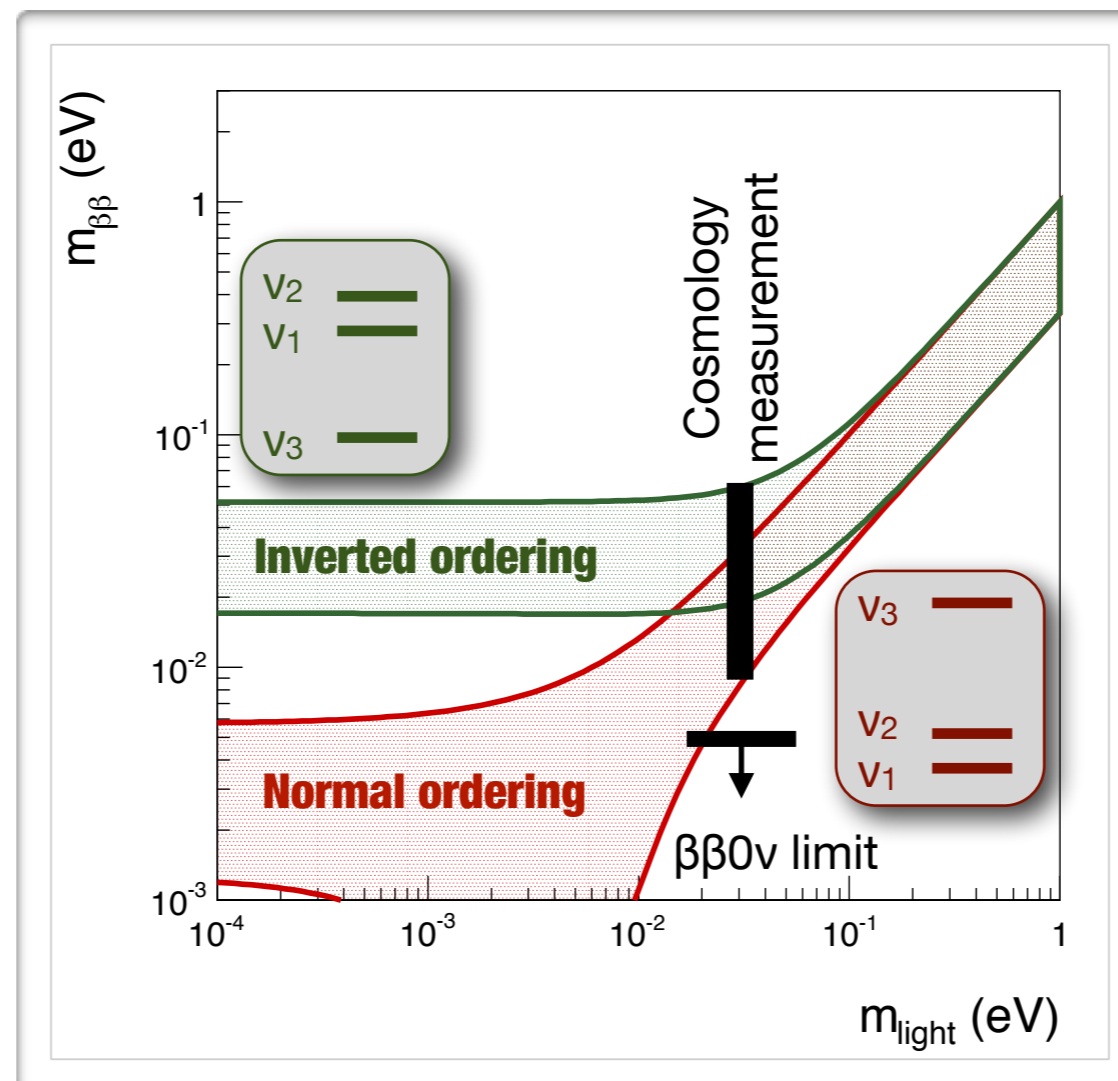
- Two X-rays
- Gamma-rays

What if ν mass is not connected to Majorana neutrinos?

- **Possible:** Dirac neutrinos!

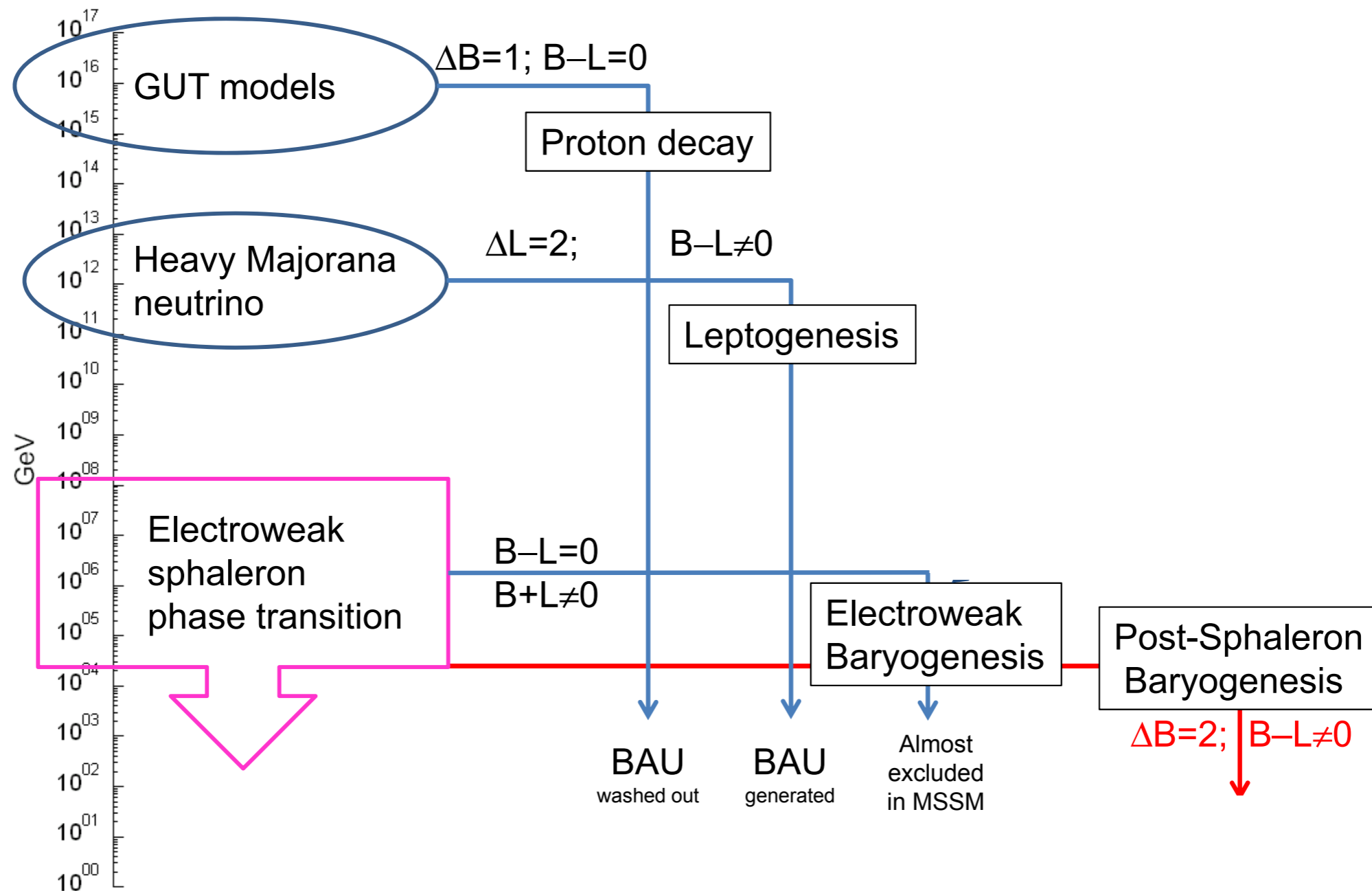
What if ν mass is not connected to Majorana neutrinos?

- **Possible:** Dirac neutrinos!
- How can we **prove** that neutrinos are Dirac particles? Difficult!
- **Best bet:** neutrino mass measured with cosmology, not in $\beta\beta 0\nu$



What if baryogenesis is not connected to Majorana neutrinos?

- Possible:** alternatives to leptogenesis-induced baryogenesis!

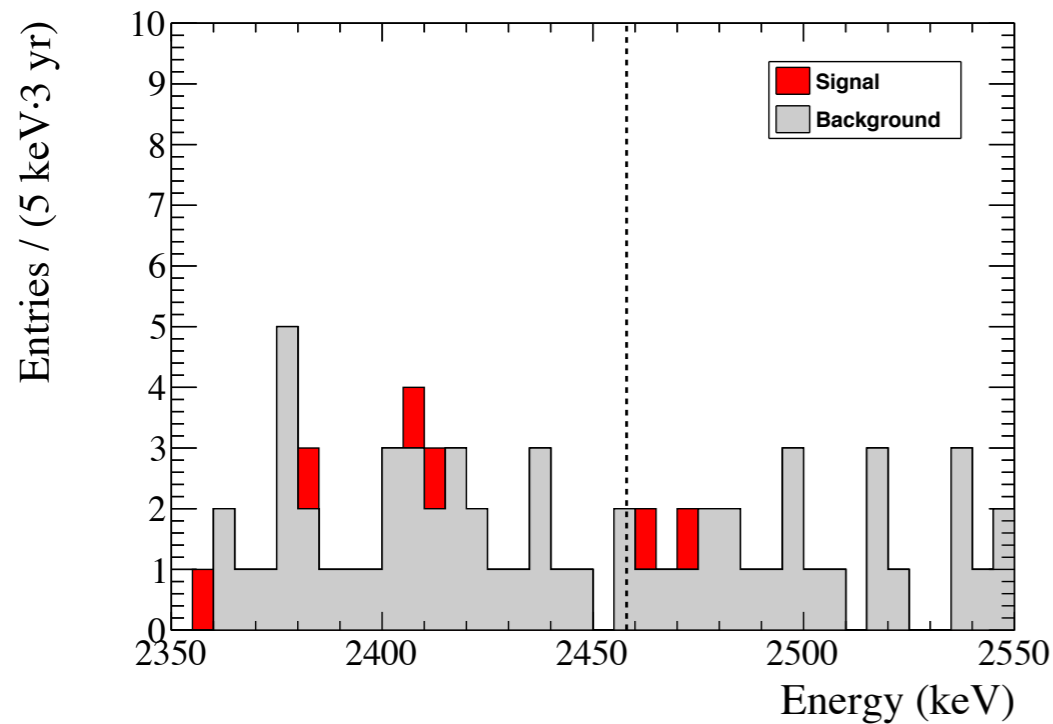


Epilogue:

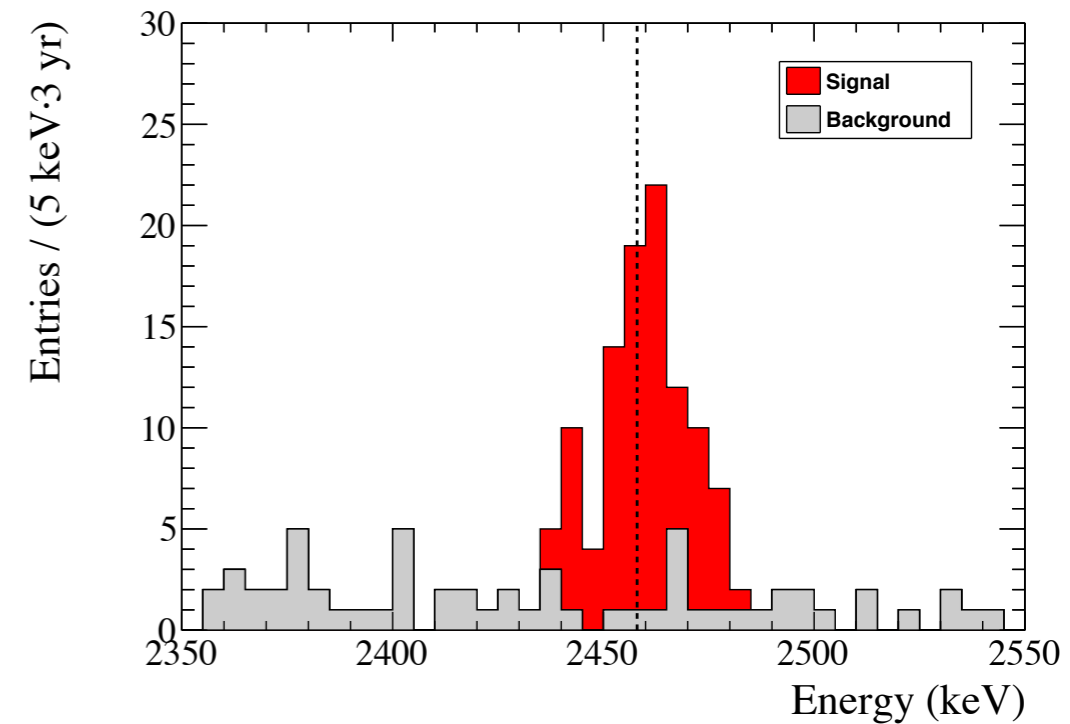
*F*inal thoughts

Double beta decay experiments are challenging

$$T_{1/2}^{0\nu} \propto \varepsilon \sqrt{\frac{Mt}{c \Delta E}}$$

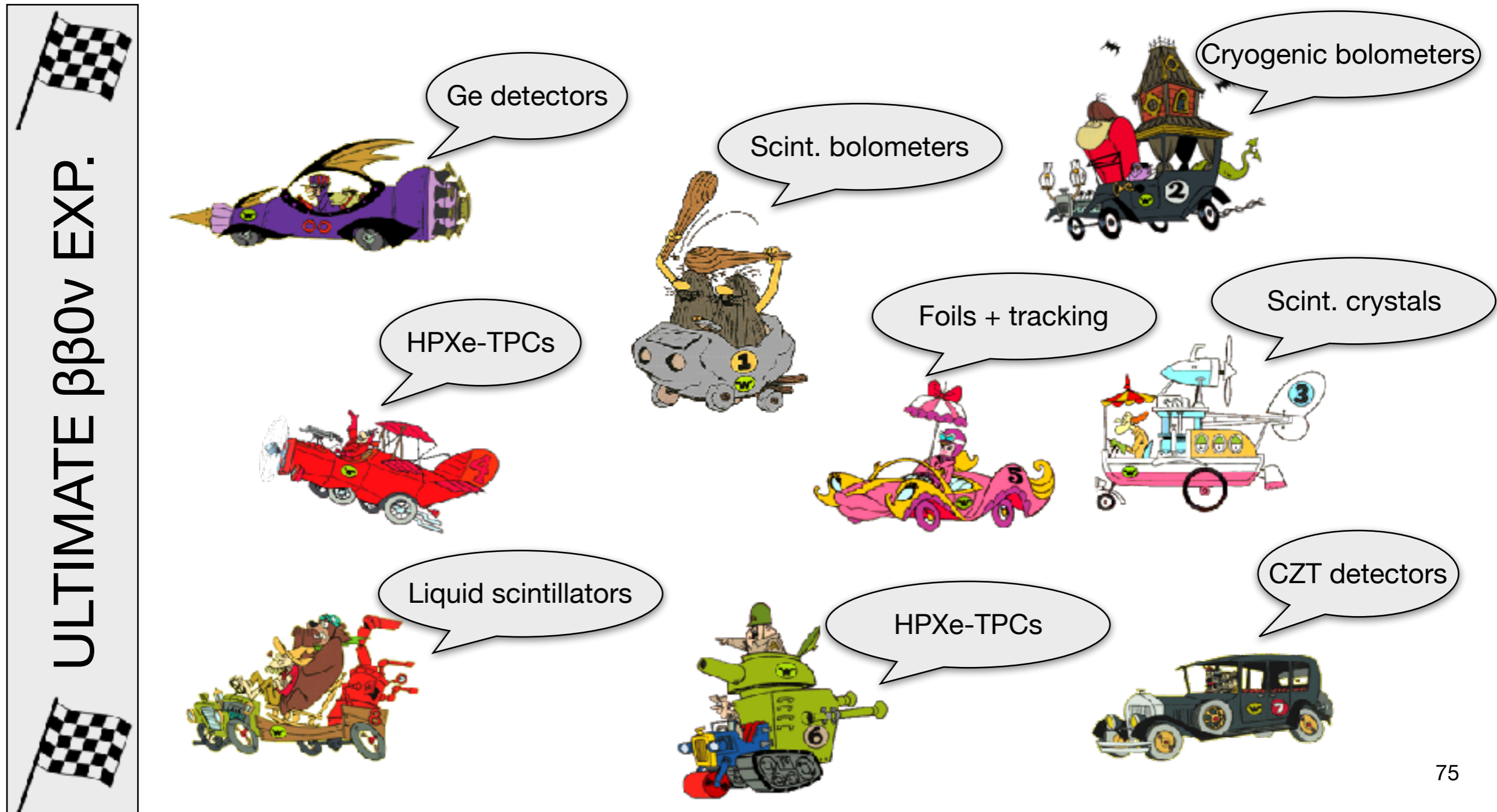


DIFFICULT



A variety of double beta decay experiments

- Non-trivial detector optimisation process, relatively inexpensive → diversity



The theory-experiment connection is essential

Small neutrino mass

Baryogenesis

Majorana neutrinos and
lepton number violation

Neutrinoless double beta decay search

Think outside the box!

Small neutrino mass

Baryogenesis

No connection?

No connection?

Majorana neutrinos and
lepton number violation

$\beta\beta 0\nu$ non-optimal probe?

Neutrinoless double beta decay search

The End.

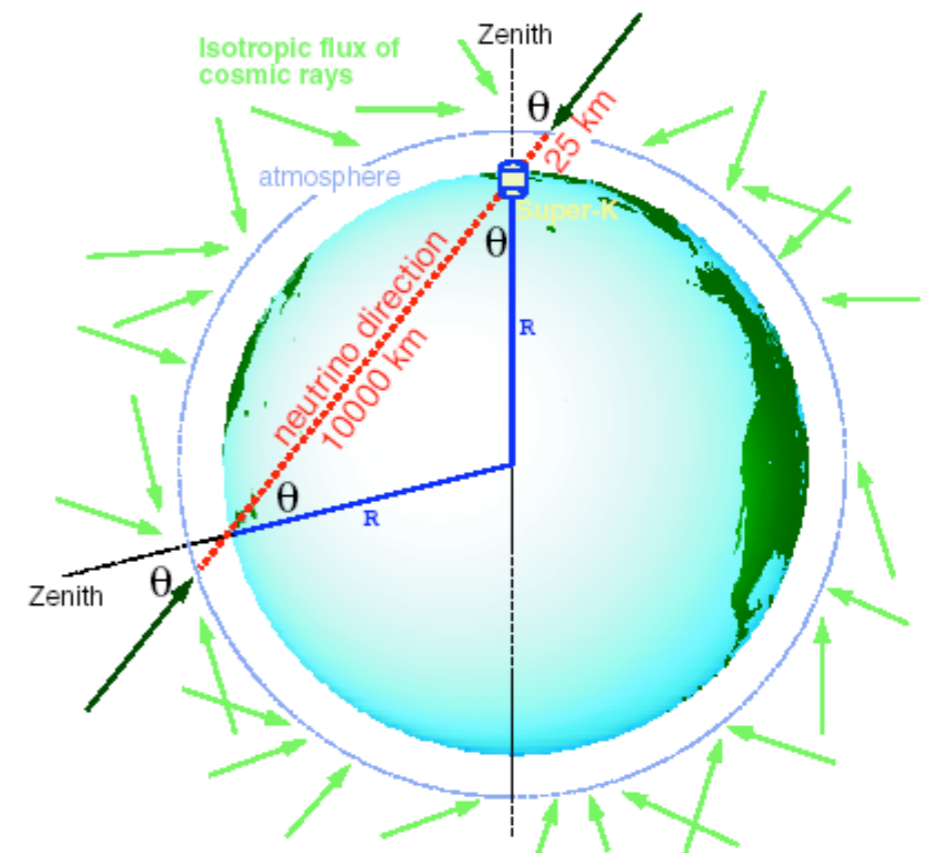
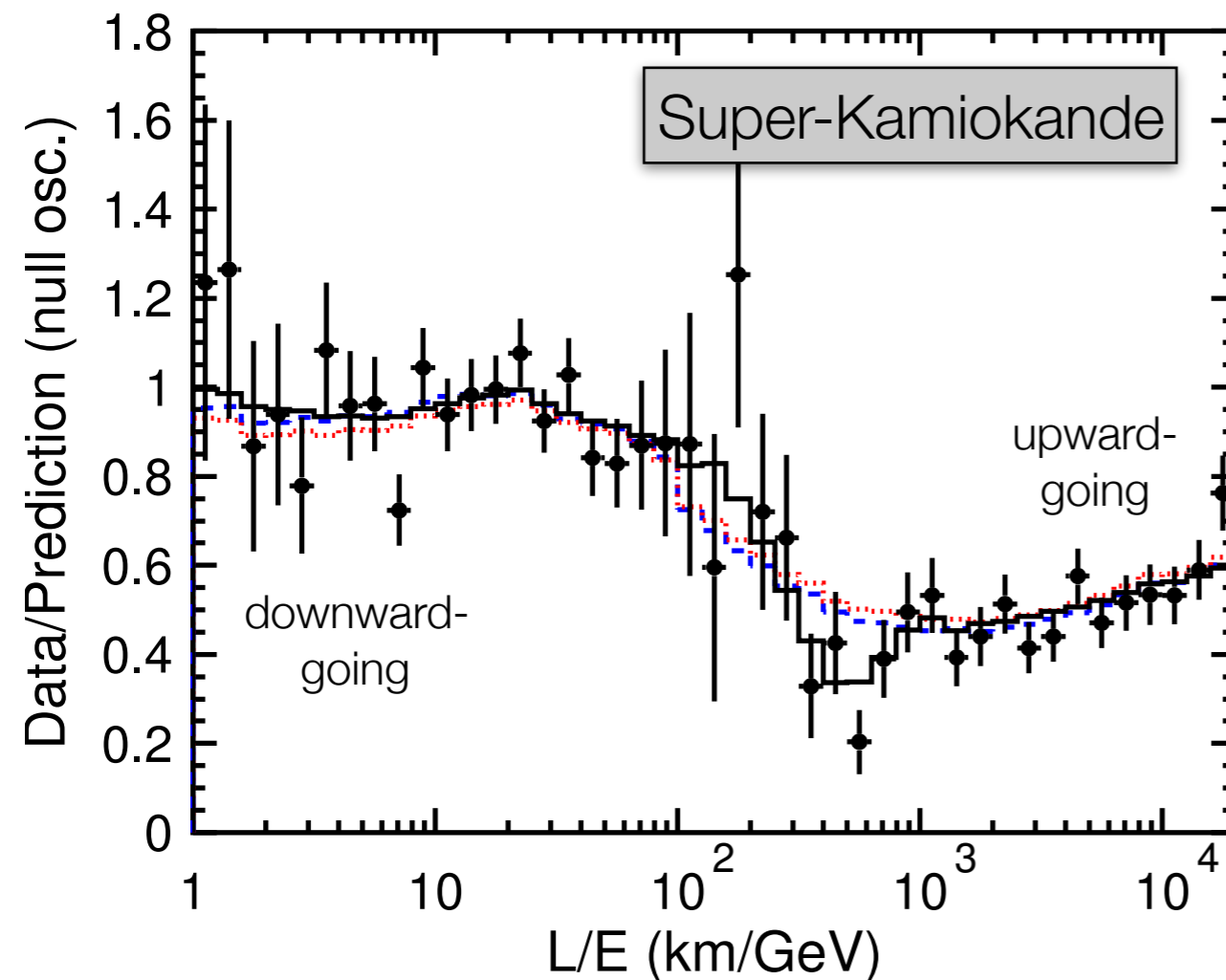
Backups

The discovery of neutrino oscillations

With atmospheric neutrinos

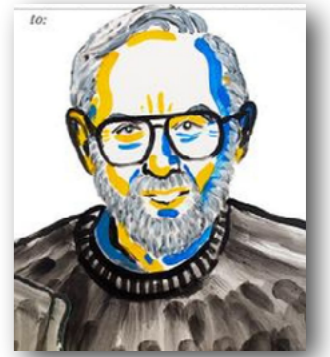


- Disappearance of atmospheric ν_μ 's and $\bar{\nu}_\mu$'s
- First conclusive evidence for oscillations: zenith angle-dependent deficit

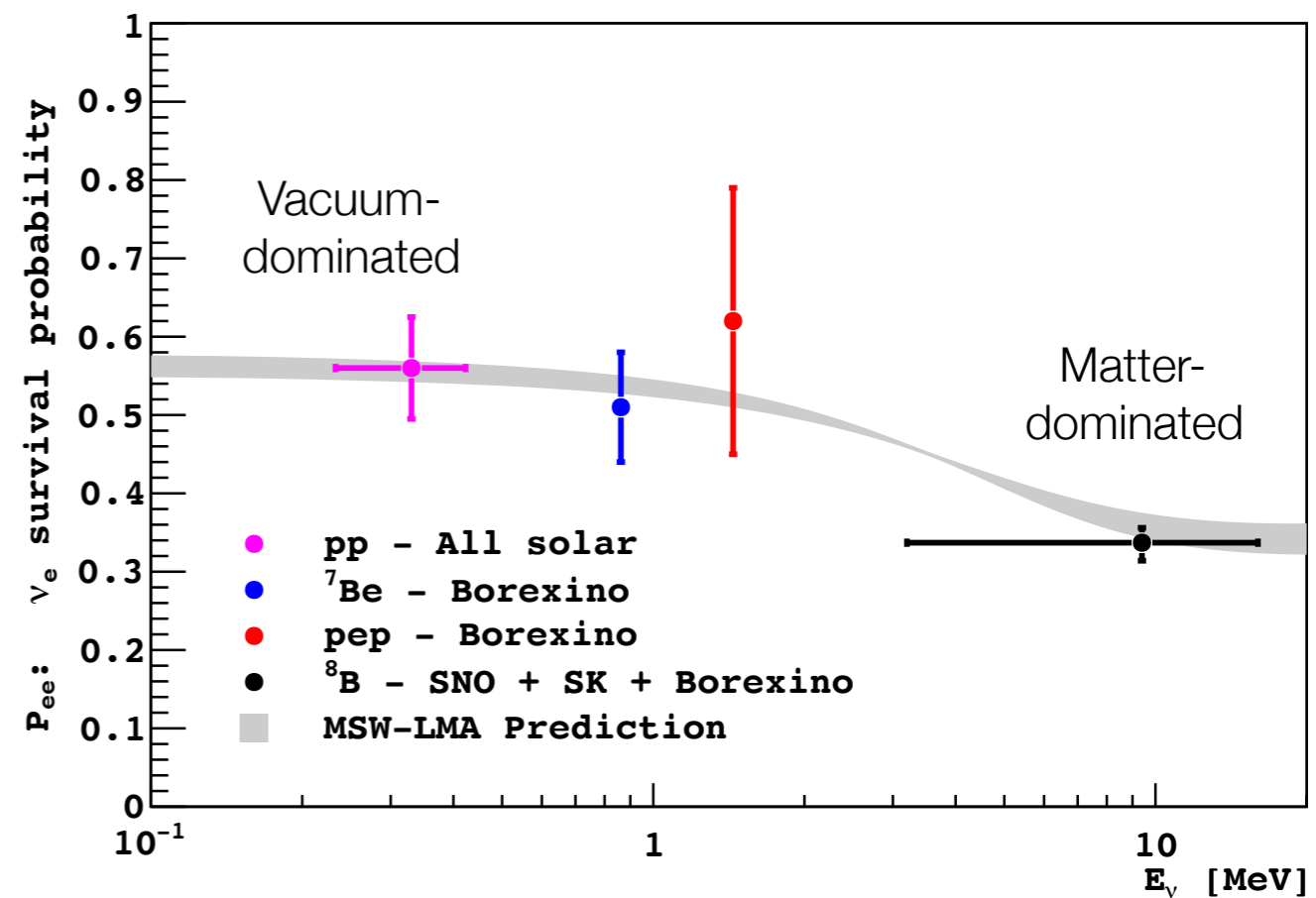
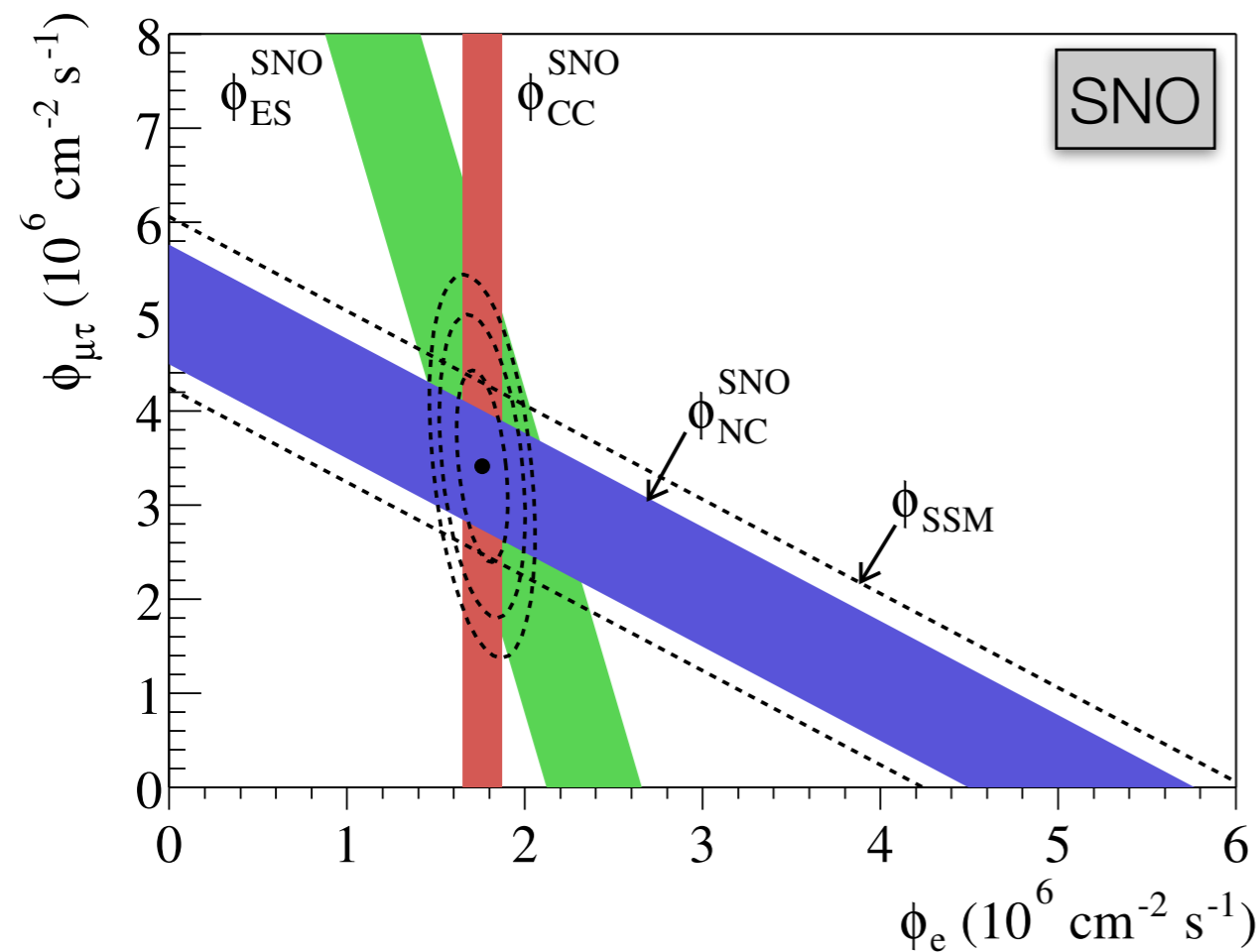


The discovery of neutrino oscillations

With solar neutrinos



- Disappearance of solar ν_e 's, appearance into other "active" flavours (μ , τ)
- Energy dependence of ν_e suppression also measured



3-Neutrino mixing parametrisation

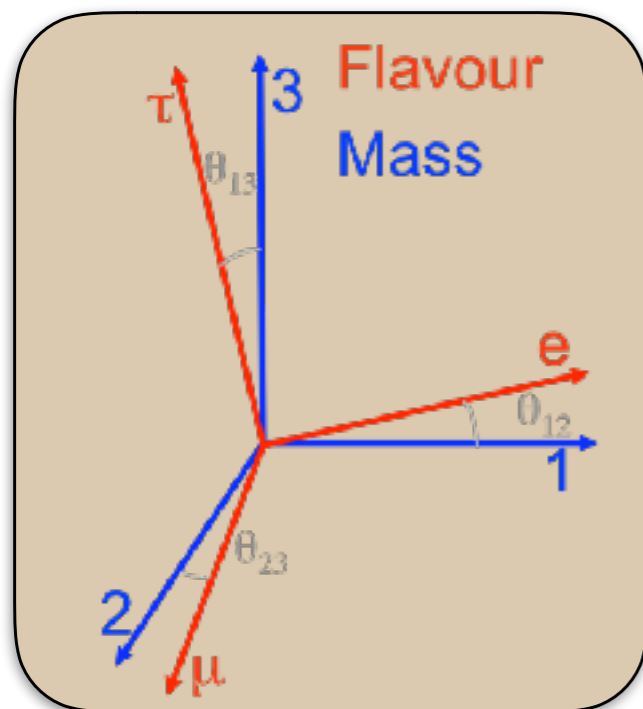
Atmospheric Oscillations

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{Atmospheric Oscillations}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}}_{\text{Interference}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar Oscillations}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

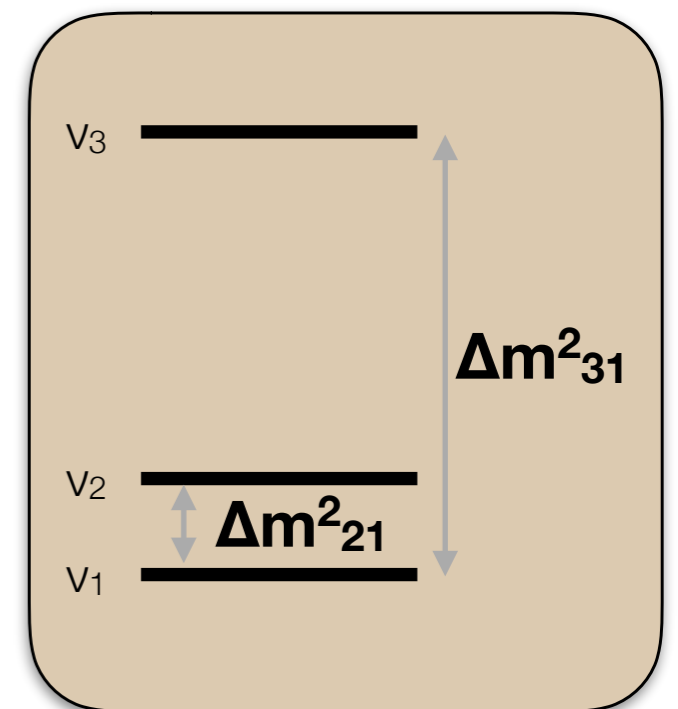
$c_{23} = \cos \theta_{23}$ etc...

Interference

Solar Oscillations



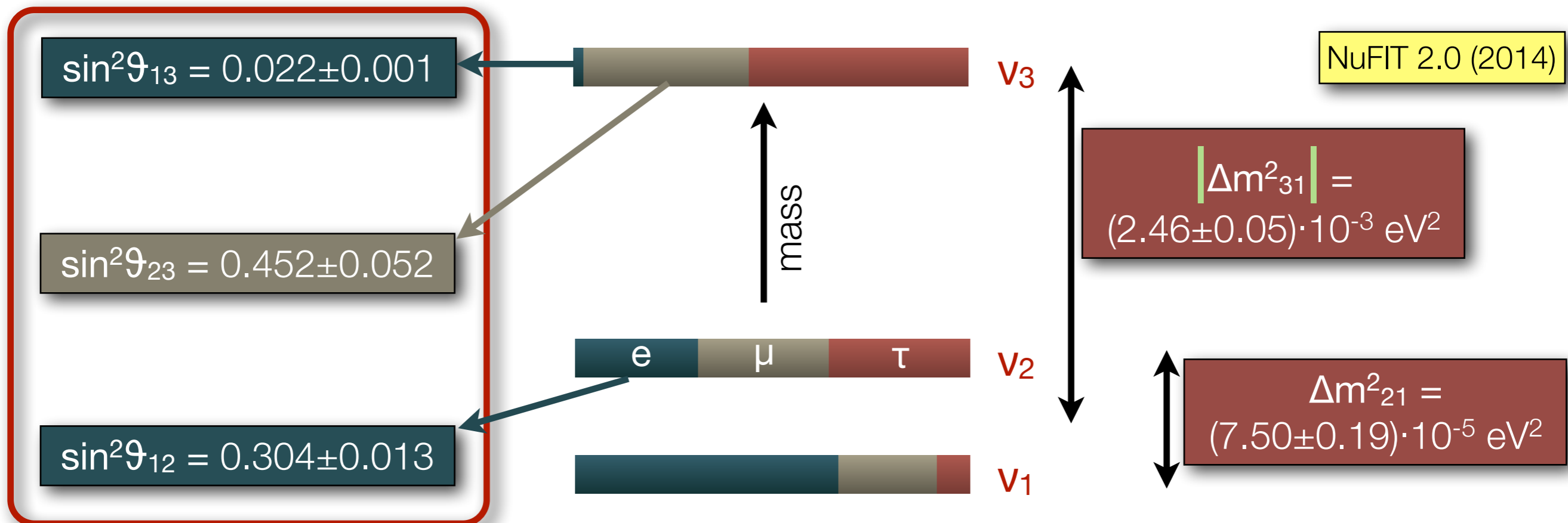
- 2 mass splittings, 3 mixing angles, 1 CPV phase
- Describe all convincing evidence for neutrino oscillations



Neutrino oscillation experimental status

3-neutrino mixing parametrisation

Mass splittings and mixing angles measured with 10% precision or better



- θ_{12} and Δm^2_{21} : solar and reactor experiments
- θ_{23} and $|\Delta m^2_{31}|$: atmospheric and accelerator experiments
- θ_{13} : reactor and accelerator experiments
- δ_{CP} phase compatible with any value at 3σ , $\text{sgn}(\Delta m^2_{31})$ unknown

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Current-generation

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- Larger phase space and less backgrounds

>2 MeV



Typical

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● Large $\beta\beta$ isotope mass

● Only way to probe 10^{26} yr half-lives!

100 kg

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● Low backgrounds in energy region of interest

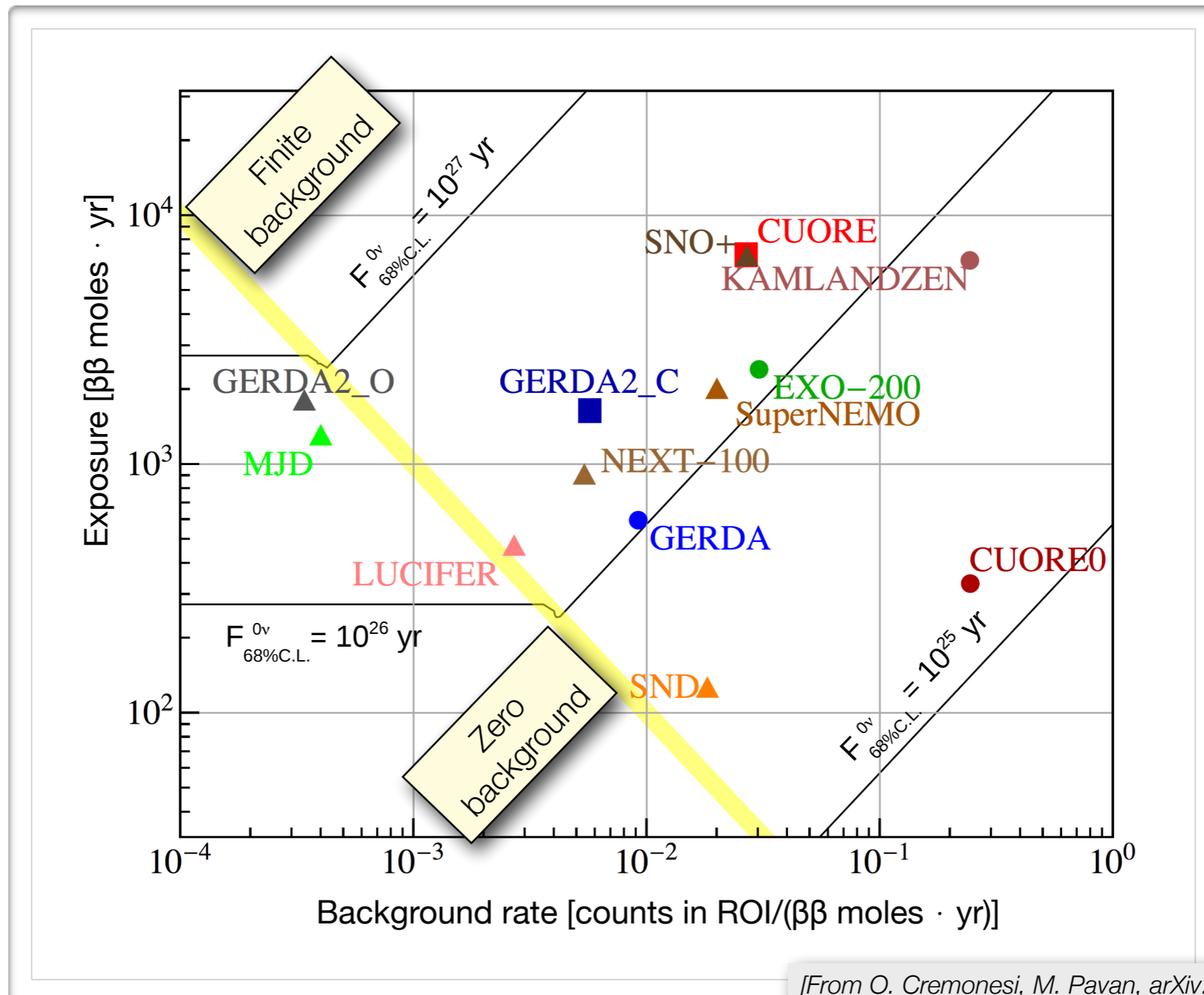
● Excellent energy resolution and/or very low background rates
(per unit energy) near $Q_{\beta\beta}$

1 cts/(100 kg·yr)

Typical

Current goal

$\beta\beta 0\nu$ experiments comparison: mass, background



[From O. Cremonesi, M. Pavan, arXiv:1310.4692]