

# Quantum Simulation of Abelian and non-Abelian Gauge Theories

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UNIVERSITÄT  
BERN

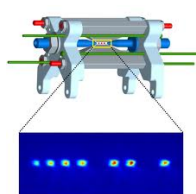
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January 16, 2015

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# Outline

A Brief History of Classical Computing

Pioneers of Quantum Computers and Quantum Simulators

Classical and Quantum Simulations of Quantum Spin Systems

From Wilson's Lattice QCD to Quantum Link Models

Quantum Simulators for Abelian Lattice Gauge Theories

Quantum Simulators for non-Abelian Lattice Gauge Theories

Conclusions

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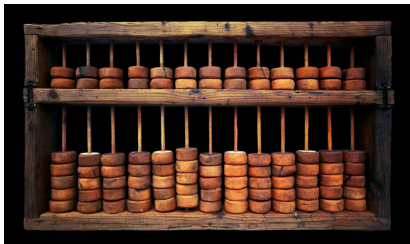
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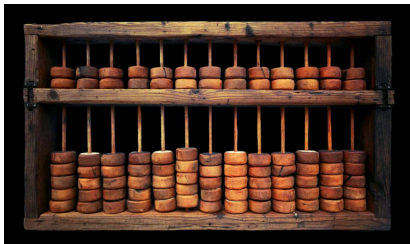
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The first “digital computer” in Babylonia about 2400 b.c.





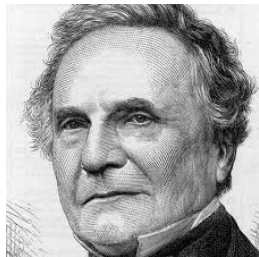
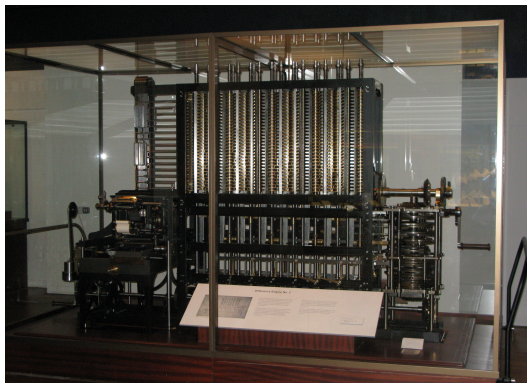
The first “digital computer” in Babylonia about 2400 b.c.



The first “analog computer”: Antikythera for determining the position of celestial bodies, Crete, about 100 b.c.

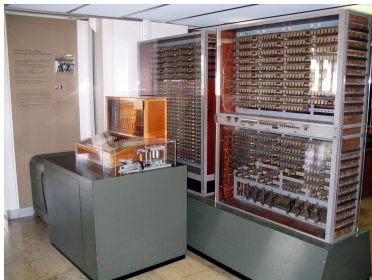


The first programmable computer:  
mechanical “difference engine”  
Charles Babbage (1791-1871)

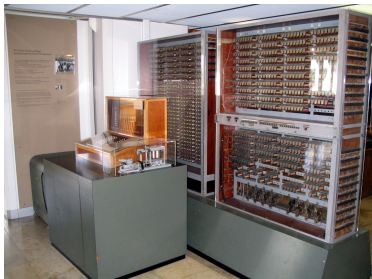


was realized by his son after Babbage's death.

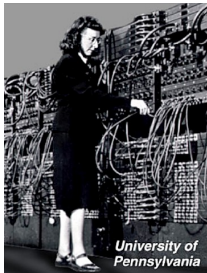
## Konrad Zuse's (1910-1992) relay-driven computer Z3



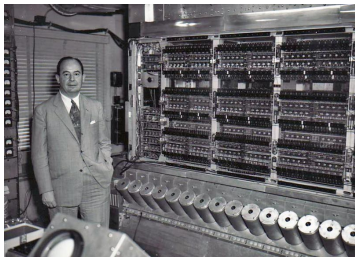
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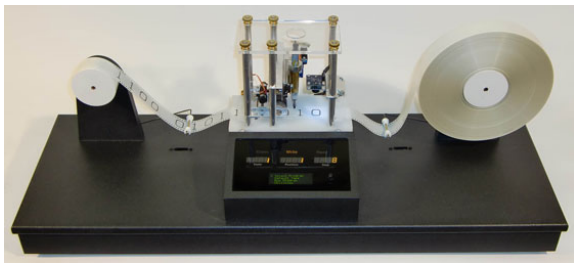
## From the vacuum-tube ENIAC to the IBM Blue Gene



## Pioneers of theoretical computer science: John von Neumann (1903-1992) and Alan Turing (1912-1954)



## Model of a universal Turing machine



RSA encryption: multiplication is easy, factorization is hard.

RSA decryption challenge in 1991:

factorize the following 174-digit number with 576 bits

*RSA576* = 18819881292060796383869723946165043980716356  
33794173827007633564229888597152346654853190  
60606504743045317388011303396716199692321205  
734031879550656996221305168759307650257059

RSA encryption: multiplication is easy, factorization is hard.

RSA decryption challenge in 1991:

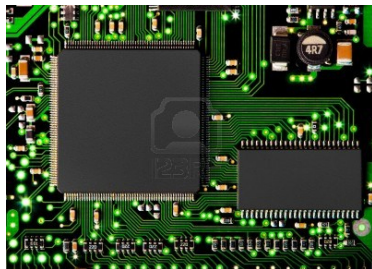
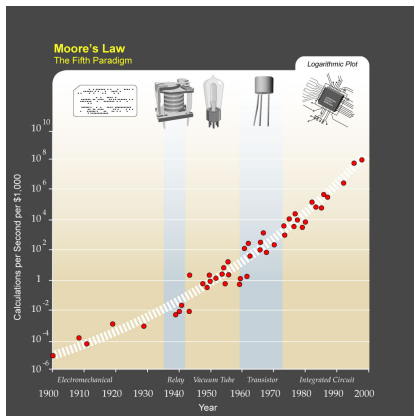
factorize the following 174-digit number with 576 bits

$$\begin{aligned} \text{RSA576} &= 18819881292060796383869723946165043980716356 \\ &\quad 33794173827007633564229888597152346654853190 \\ &\quad 60606504743045317388011303396716199692321205 \\ &\quad 734031879550656996221305168759307650257059 \\ &= 39807508642406493739712550055038649119906436 \\ &\quad 2342526708406385189575946388957261768583317 \\ &* 47277214610743530253622307197304822463291469 \\ &\quad 5302097116459852171130520711256363590397527 \end{aligned}$$

This problem was solved only in 2003 by two mathematicians in Bonn using very large computer resources.

Only in 2009, when the challenge was no longer active, the 232-digit number RSA768 with 768 bits has finally been factorized.

Moore's law: "Every two years the number of transistors per area increases by a factor of 2."



Modern micro chips consist of several billions of transistors, each about  $10^{-8}$  m in size. This is already close to the quantum mechanical limit set by the size of individual atoms.



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## Richard Feynman's vision of 1982



“I’m not happy with all the analyses that go with just the classical theory, because nature isn’t classical, dammit, and if you want to make a simulation of nature, you’d better make it quantum mechanical, and by golly it’s a wonderful problem, because it doesn’t look so easy.”

Deutsch's universal quantum computer could use Shor's algorithm to solve the factorization problem.



David Deutsch



Peter Shor

Deutsch's universal quantum computer could use Shor's algorithm to solve the factorization problem.

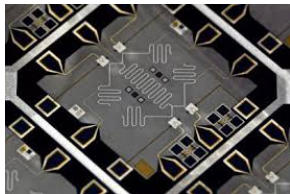


David Deutsch

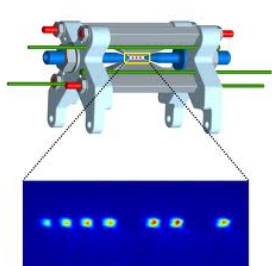
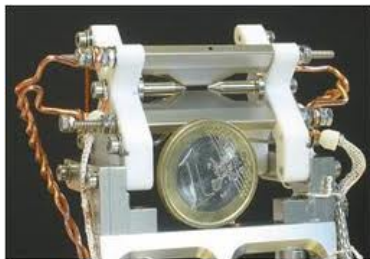


Peter Shor

Until today, only  $15 = 3 \cdot 5$  has been correctly factorized by a quantum computer, at least in about 50 % of all trials.



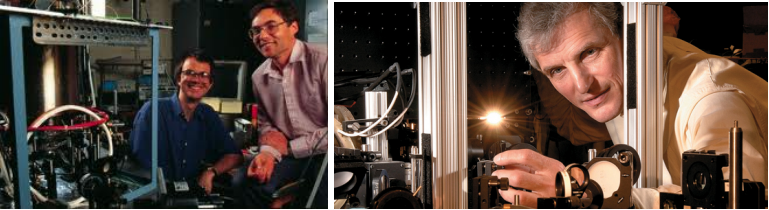
## Ion traps as a digital quantum computer?



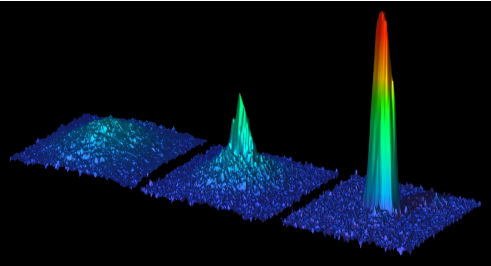
Franklin Medal 2010: I. Cirac, D. Wineland, P. Zoller



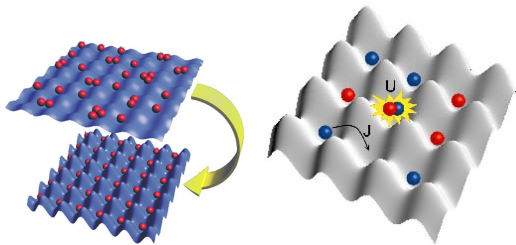
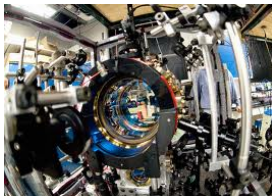
# Bose-Einstein condensation in ultra-cold atomic gases



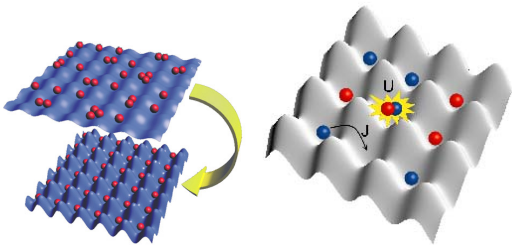
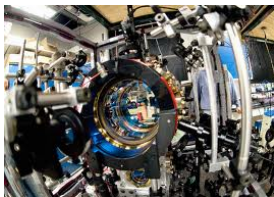
Eric Cornell, Carl Wieman, Wolfgang Ketterle, 1995



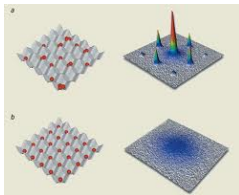
# Ultra-cold atoms in optical lattices as an analog quantum simulator for the bosonic Hubbard model



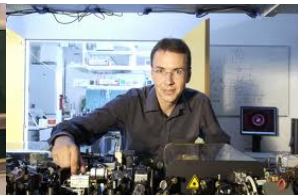
# Ultra-cold atoms in optical lattices as an analog quantum simulator for the bosonic Hubbard model



## Transition from a superfluid to a Mott insulator



Theodor Hänsch



Immanuel Bloch

Can one understand high- $T_c$  superconductivity in this way?



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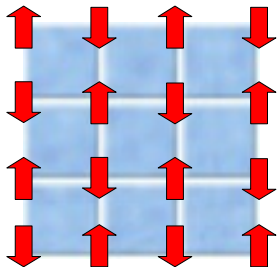
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## The spin $\frac{1}{2}$ quantum Heisenberg model



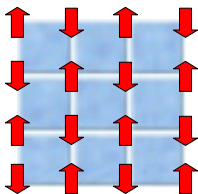
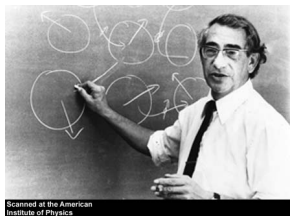
Quantum spins  $[S_x^a, S_y^b] = i\delta_{xy}\epsilon_{abc}S_x^c$  and their Hamiltonian

$$H = J \sum_{\langle xy \rangle} \vec{S}_x \cdot \vec{S}_y$$

Partition function at inverse temperature  $\beta = 1/T$

$$Z = \text{Tr} \exp(-\beta H)$$

## The Hubbard Model for doped antiferromagnets



$$H = -t \sum_{\langle xy \rangle} (c_x^\dagger c_y + c_y^\dagger c_x) + U \sum_x (c_x^\dagger c_x - 1)^2, \quad c_x = \begin{pmatrix} c_{x\uparrow} \\ c_{x\downarrow} \end{pmatrix}$$

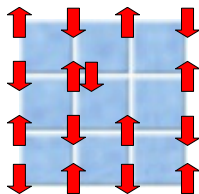
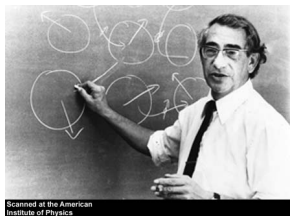
### Sign problem of fermionic path integrals

$$Z_f = \text{Tr} \exp(-\beta H) = \sum_{[n]} \text{Sign}[n] \exp(-S[n]), \quad \text{Sign}[n] = \pm 1$$

### Average sign is exponentially small

$$\langle \text{Sign} \rangle = \frac{\sum_{[n]} \text{Sign}[n] \exp(-S[n])}{\sum_{[n]} \exp(-S[n])} = \frac{Z_f}{Z_b} = \exp(-\beta V \Delta f)$$

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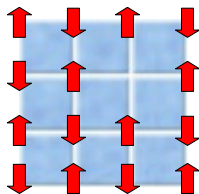
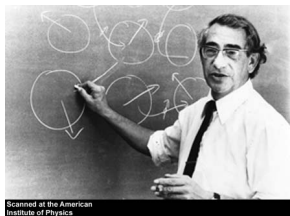
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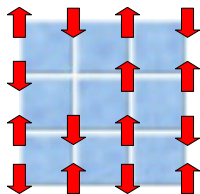
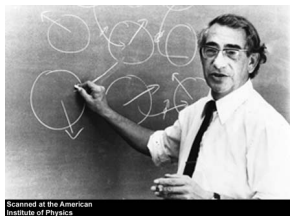
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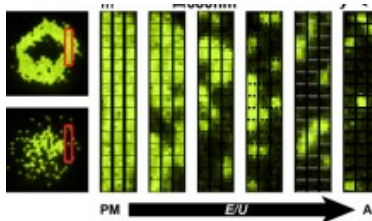
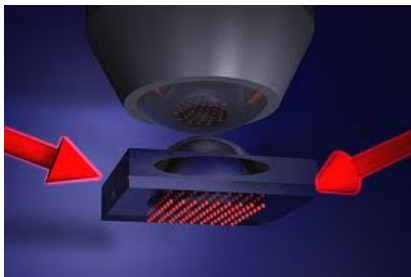
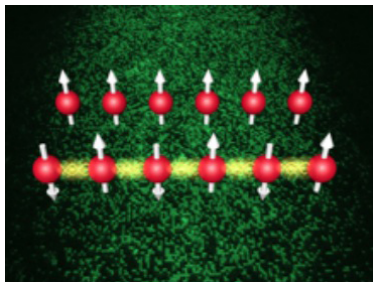
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## Optical lattice quantum simulation of quantum spin systems



J. Simon, W. S. Bakir, R. Ma, M. E. Tal, P. M. Preis, M. Greiner,  
Nature 472 (2011) 307.

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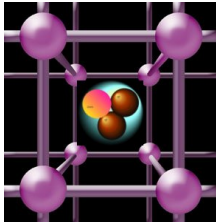
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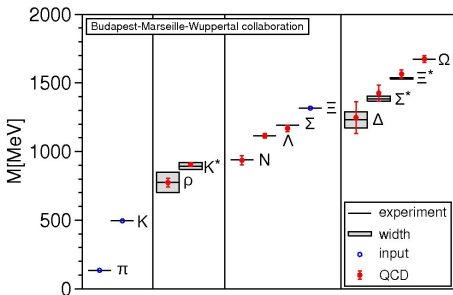
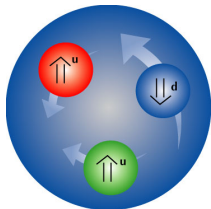
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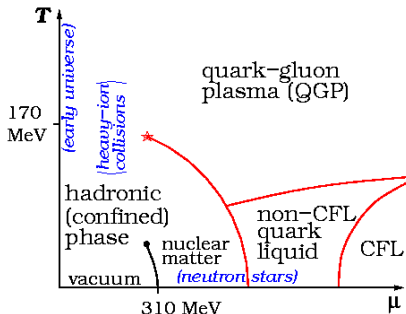
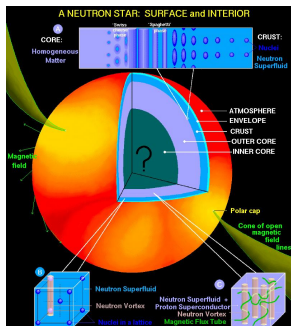
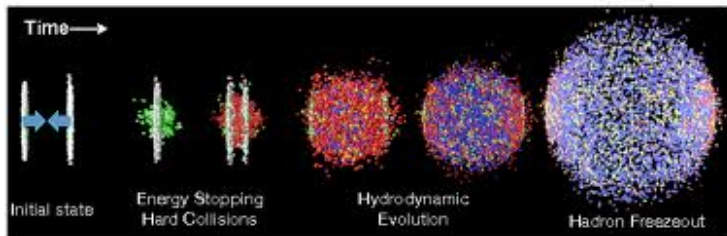
# Kenneth Wilson's lattice QCD describes confinement of quarks and gluons inside protons and neutrons



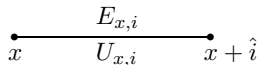
and confirms the experimentally measured mass spectrum



# Can heavy-ion collision physics or nuclear astrophysics benefit from quantum simulations in the long run?



## Hamiltonian formulation of Wilson's $U(1)$ lattice gauge theory


$$\begin{array}{ccc} & E_{x,i} & \\ \bullet & \text{---} & \bullet \\ x & U_{x,i} & x + \hat{i} \end{array}$$

$$U = \exp(iA), \quad U^\dagger = \exp(-iA) \in U(1)$$

### Electric field operator $E$

$$E = -i\partial_A, \quad [E, U] = U, \quad [E, U^\dagger] = -U^\dagger, \quad [U, U^\dagger] = 0$$

### Generator of $U(1)$ gauge transformations

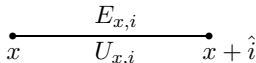
$$G_x = \sum_i (E_{x-\hat{i},i} - E_{x,i}), \quad [H, G_x] = 0$$

### $U(1)$ gauge invariant Hamiltonian

$$H = \frac{g^2}{2} \sum_{x,i} E_{x,i}^2 - \frac{1}{2g^2} \sum_{x,i \neq j} (U_{x,i} U_{x+\hat{i},j} U_{x+\hat{j},i}^\dagger U_{x,j}^\dagger + \text{h.c.})$$

operates in an infinite-dimensional Hilbert space per link

$U(1)$  quantum links from spins  $\frac{1}{2}$

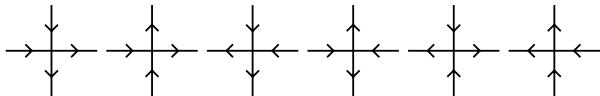


$$U = S_1 + iS_2 = S_+, \quad U^\dagger = S_1 - iS_2 = S_-$$

Electric flux operator  $E$

$$E = S_3, \quad [E, U] = U, \quad [E, U^\dagger] = -U^\dagger, \quad [U, U^\dagger] = 2E$$

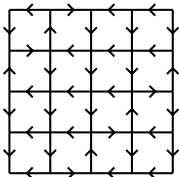
Gauss law



Ring-exchange plaquette Hamiltonian

$$H \begin{array}{c} \leftarrow \\ \downarrow \leftarrow \rightarrow \uparrow \\ \rightarrow \\ \leftarrow \end{array} = J \begin{array}{c} \rightarrow \\ \uparrow \leftarrow \rightarrow \downarrow \\ \leftarrow \\ \rightarrow \end{array}$$

$$H \begin{array}{c} \rightarrow \\ \downarrow \rightarrow \leftarrow \uparrow \\ \leftarrow \\ \rightarrow \end{array} = 0$$



D. Horn, Phys. Lett. B100 (1981) 149

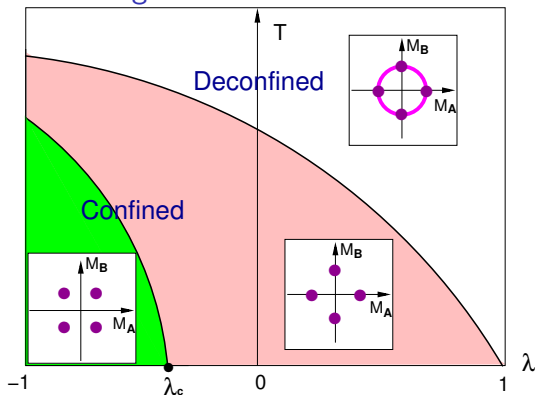
P. Orland, D. Rohrlich, Nucl. Phys. B338 (1990) 647

S. Chandrasekharan, UJW, Nucl. Phys. B492 (1997) 455

## Hamiltonian with Rokhsar-Kivelson term

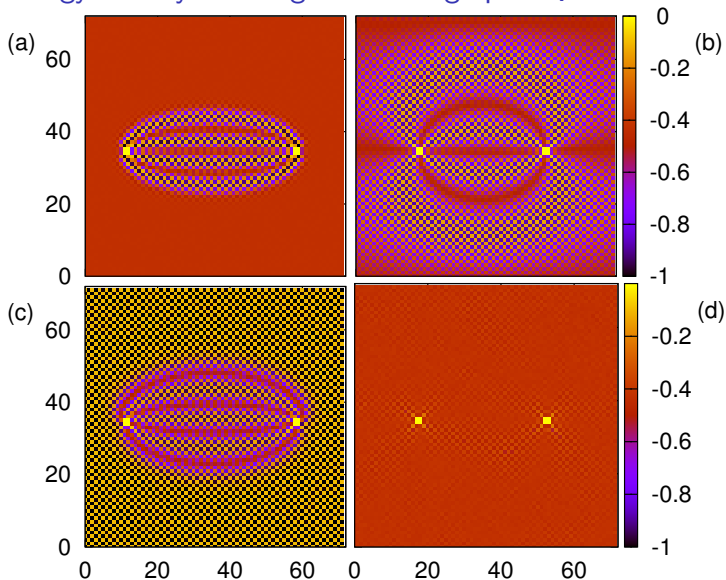
$$H = -J \left[ \sum_{\square} (U_{\square} + U_{\square}^{\dagger}) - \lambda \sum_{\square} (U_{\square} + U_{\square}^{\dagger})^2 \right]$$

## Phase diagram



D. Banerjee, F.-J. Jiang, P. Widmer, UJW, JSTAT (2013) P12010.

## Energy density of charge-anti-charge pair $Q = \pm 2$



D. Banerjee, F.-J. Jiang, P. Widmer, UJW, JSTAT (2013) P12010.

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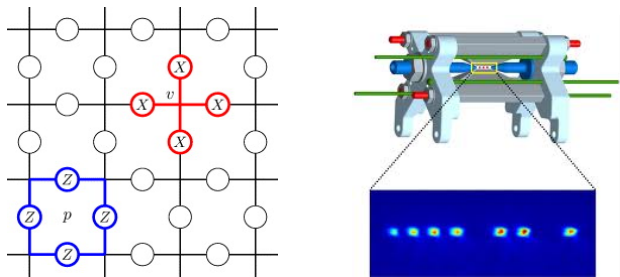
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## Digital quantum simulation of Kitaev's toric code (a $\mathbb{Z}(2)$ quantum link model) with trapped ions

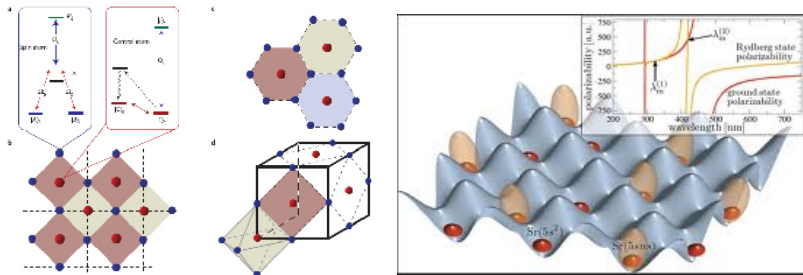


- Precisely controllable many-body quantum device, which can execute a prescribed sequence of quantum gate operations.
- State of simulated system encoded as quantum information.
- Dynamics is represented by a sequence of quantum gates, following a stroboscopic Trotter decomposition.

A. Y. Kitaev, *Ann. Phys.* 303 (2003) 2.

B. P. Lanyon, C. Hempel, D. Nigg, M. Müller, R. Gerritsma, F. Zähringer, P. Schindler, J. T. Barreiro, M. Rambach, G. Kirchmair, M. Hennrich, P. Zoller, R. Blatt, C. F. Roos, *Science* 334 (2011) 6052.

## $U(1)$ quantum link models can also be simulated with Rydberg atoms in an optical lattice



- Lasers can excite atoms to high-lying Rydberg states.
- Rydberg atoms are large and have collective interactions.
- Ensemble Rydberg atoms represent qubits at link centers.
- Control atoms at lattice sites ensure the Gauss' law.

M. Müller, I. Lesanovsky, H. Weimer, H. P. Büchler, P. Zoller, *Phys. Rev. Lett.* 102 (2009) 170502.

H. Weimer, M. Müller, I. Lesanovsky, P. Zoller, H. P. Büchler, *Nat. Phys.* 6 (2010) 382.

L. Tagliacozzo, A. Celi, P. Orland, M. Lewenstein, *Nature Communications* 4 (2013) 2615.

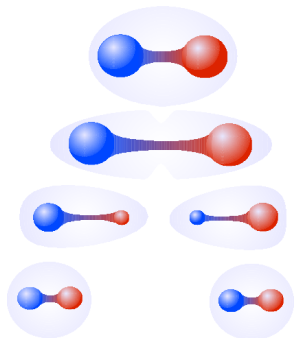
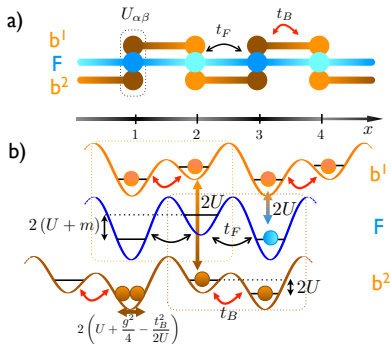
L. Tagliacozzo, A. Celi, A. Zamora, M. Lewenstein, *Ann. Phys.* 330 (2013) 160.

## Hamiltonian for staggered fermions and $U(1)$ quantum links

$$H = -t \sum_x \left[ \psi_x^\dagger U_{x,x+1} \psi_{x+1} + \text{h.c.} \right] + m \sum_x (-1)^x \psi_x^\dagger \psi_x + \frac{g^2}{2} \sum_x E_{x,x+1}^2$$

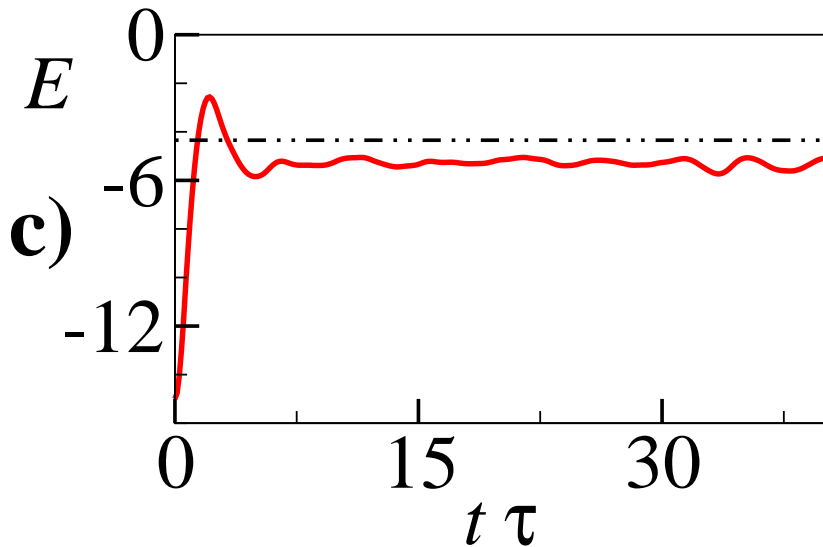
$$U_{x,x+1} = b_x b_{x+1}^\dagger, \quad E_{x,x+1} = \frac{1}{2} \left( b_{x+1}^\dagger b_{x+1} - b_x^\dagger b_x \right)$$

## Optical lattice with Bose-Fermi mixture of ultra-cold atoms



D. Banerjee, M. Dalmonte, M. Müller, E. Rico, P. Stebler, UJW,  
P. Zoller, Phys. Rev. Lett. 109 (2012) 175302.

# Quantum simulation of the real-time evolution of string breaking



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Conclusions

$U(N)$  quantum link operators

$$U^{ij} = S_1^{ij} + iS_2^{ij}, \quad U^{ij\dagger} = S_1^{ij} - iS_2^{ij}, \quad i, j \in \{1, 2, \dots, N\}, \quad [U^{ij}, (U^\dagger)^{kl}] \neq 0$$

$SU(N)_L \times SU(N)_R$  gauge transformations of a quantum link

$$[L^a, L^b] = if_{abc}L^c, \quad [R^a, R^b] = if_{abc}R^c, \quad a, b, c \in \{1, 2, \dots, N^2 - 1\}$$

$$[L^a, R^b] = [L^a, E] = [R^a, E] = 0$$

Infinitesimal gauge transformations of a quantum link

$$[L^a, U] = -\lambda^a U, \quad [R^a, U] = U\lambda^a, \quad [E, U] = U$$

Algebraic structures of  $U(N)$  quantum link models

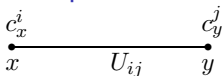
$$U^{ij}, L^a, R^a, E, \quad 2N^2 + 2(N^2 - 1) + 1 = 4N^2 - 1 \text{ } SU(2N) \text{ generators}$$

R. Brower, S. Chandrasekharan, UJW, Phys. Rev. D60 (1999) 094502

## Fermionic rishons at the two ends of a link

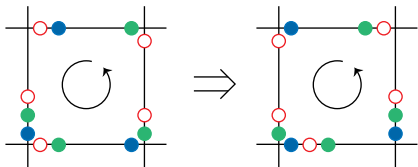
$$\{c_x^i, c_y^{j\dagger}\} = \delta_{xy}\delta_{ij}, \quad \{c_x^i, c_y^j\} = \{c_x^{i\dagger}, c_y^{j\dagger}\} = 0$$

## Rishon representation of link algebra

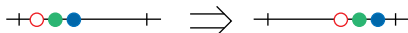


$$U_{xy}^{ij} = c_x^i c_y^{j\dagger}, \quad L_{xy}^a = c_x^{i\dagger} \lambda_{ij}^a c_x^j, \quad R_{xy}^a = c_y^{i\dagger} \lambda_{ij}^a c_y^j, \quad E_{xy} = \frac{1}{2}(c_y^{i\dagger} c_y^i - c_x^{i\dagger} c_x^i)$$

Can a “rishon abacus” implemented with ultra-cold atoms be used as a quantum simulator?

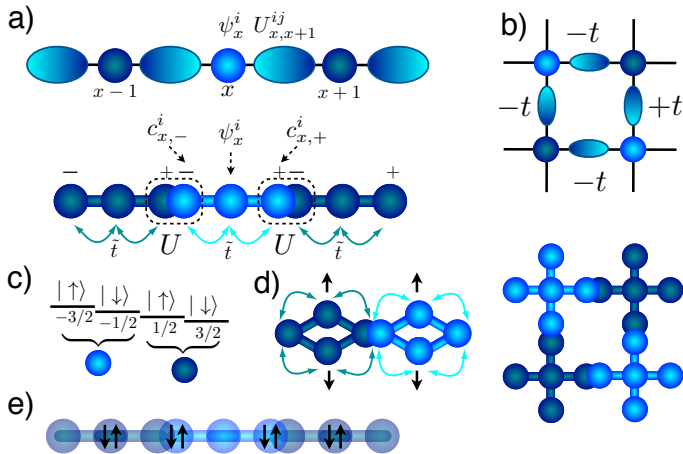


$\text{Tr Up}$



$\det U_{x,\mu}$

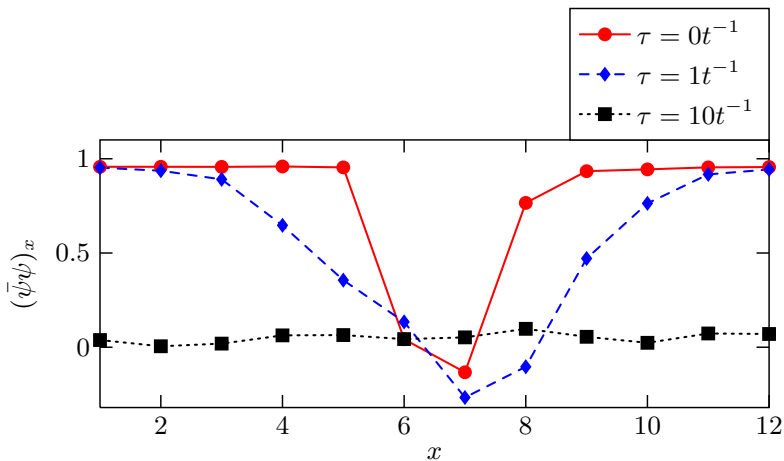
# Optical lattice with ultra-cold alkaline-earth atoms ( $^{87}\text{Sr}$ or $^{173}\text{Yb}$ ) with color encoded in nuclear spin



D. Banerjee, M. Bögli, M. Dalmonte, E. Rico, P. Stebler, UJW, P. Zoller, Phys. Rev. Lett. 110 (2013) 125303

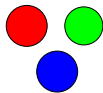


## Expansion of a “fireball” mimicking a hot quark-gluon plasma



# Nuclear Physics from $SU(3)$ QCD

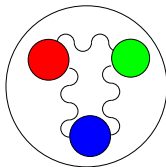
Quarks



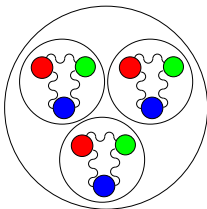
Gluon



Baryon

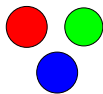


Nucleus



# Nuclear Physics from $SU(3)$ QCD

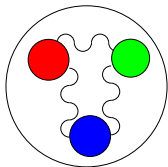
Quarks



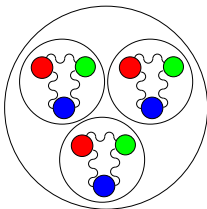
Gluon



Baryon



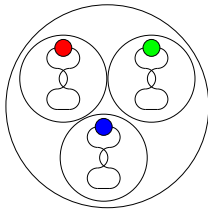
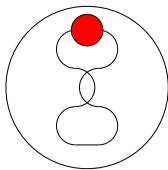
Nucleus



“Nuclear Physics” in an  $SO(3)$  lattice gauge theory?

$SO(3)$  “Nucleus”

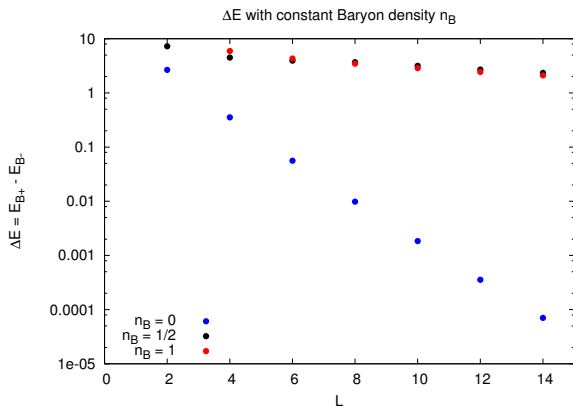
$SO(3)$  “Baryon”



## 1-d $SO(3)$ quantum link model with adjoint triplet-fermions

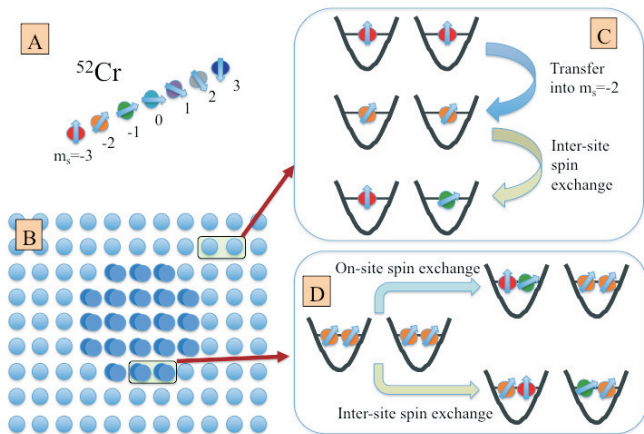
$$H = -t \sum_x \left[ \psi_x^{i\dagger} O_{x,x+1}^{ij} \psi_{x+1}^j + \text{h.c.} \right] + m \sum_x (-1)^x \psi_x^{i\dagger} \psi_x^i$$

Restoration of chiral symmetry at baryon density  $n_B \geq \frac{1}{2}$



M. Dalmonte, E. Rico, D. Banerjee, M. Bögli, P. Stebler, UJW, P. Zoller, in preparation

# Implementation with magnetic atoms (e.g. Cr), whose dipolar interactions allow spin-spin interactions without superexchange



A. de Paz, A. Sharma, A. Chotia, E. Marechal, J. H. Huckans, P. Pedri, L. Santos, O. Gorceix, L. Vernac, and B. Laburthe-Tolra, Phys. Rev. Lett. 111 (2013) 185305.

## Analog quantum simulator proposals

H. P. Büchler, M. Hermele, S. D. Huber, M. P. A. Fisher, P. Zoller, Phys. Rev. Lett. 95 (2005) 040402.

E. Zohar, B. Reznik, Phys. Rev. Lett. 107 (2011) 275301.

E. Zohar, J. I. Cirac, B. Reznik, Phys. Rev. Lett. 109 (2012) 125302;

Phys. Rev. Lett. 110 (2013) 055302; Phys. Rev. Lett. 110 (2013) 125304.

D. Banerjee, M. Dalmonte, M. Müller, E. Rico, P. Stebler, UJW, P. Zoller, Phys. Rev. Lett. 109 (2012) 175302.

D. Banerjee, M. Bögli, M. Dalmonte, E. Rico, P. Stebler, UJW, P. Zoller, Phys. Rev. Lett. 110 (2013) 125303

## Digital quantum simulator proposals

M. Müller, I. Lesanovsky, H. Weimer, H. P. Büchler, P. Zoller, Phys. Rev. Lett. 102 (2009) 170502; Nat. Phys. 6 (2010) 382.

L. Tagliacozzo, A. Celi, P. Orland, M. Lewenstein, Nature Communications 4 (2013) 2615.

L. Tagliacozzo, A. Celi, A. Zamora, M. Lewenstein, Ann. Phys. 330 (2013) 160.

## Review on quantum simulators for lattice gauge theories

UJW, Annalen der Physik 525 (2013) 777, arXiv:1305.1602.

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- The path towards quantum simulation of QCD will be a long one. **However, with a lot of interesting physics along the way,**